

THE FIBERING METHOD APPROACH FOR A NON-LINEAR SCHRÖDINGER EQUATION COUPLED WITH THE ELECTROMAGNETIC FIELD

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ABSTRACT. We study, with respect to the parameter $q \neq 0$, the following Schrödinger-Bopp-Podolsky system in \mathbb{R}^3

$$\begin{cases} -\Delta u + \omega u + q^2 \phi u = |u|^{p-2} u, \\ -\Delta \phi + a^2 \Delta^2 \phi = 4\pi u^2, \end{cases}$$

where $p \in (2, 3]$, $\omega > 0$, $a \geq 0$ are fixed. We prove, by means of the fibering approach, that the system has no solutions *at all* for large values of q 's, and has two *radial* solutions for small q 's. We give also qualitative properties about the energy level of the solutions and a variational characterization of these extremals values of q . Our results recover and improve some results in [2, 5].

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1. INTRODUCTION

In the recent paper [2] for the first time in the mathematical literature the following system in \mathbb{R}^3 has been studied

$$(1.1) \quad \begin{cases} -\Delta u + \omega u + q^2 \phi u = |u|^{p-2} u, \\ -\Delta \phi + a^2 \Delta^2 \phi = 4\pi u^2, \end{cases}$$

where $a, \omega > 0$, $q \neq 0$ and $p \in (2, 6)$. The system appears when one looks for stationary solutions $u(x)e^{i\omega t}$ of the Schrödinger equation coupled with the Bopp-Podolski Lagrangian of the electromagnetic field, in the purely electrostatic situation. Here u represents the modulus of the wave function and ϕ the electrostatic potential. From a physical point of view, the parameter q has the meaning of the electric charge and a is the parameter of the Bopp-Podolski term.

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In the cited paper, it has been shown that the problem can be addressed variationally. Indeed introducing the Hilbert space

$$\mathcal{D} := \{\phi \in D^{1,2}(\mathbb{R}^3) : \Delta\phi \in L^2(\mathbb{R}^3)\}$$

normed by

$$\|\phi\|_{\mathcal{D}}^2 = a^2\|\Delta\phi\|_2^2 + \|\nabla\phi\|_2^2,$$

it can be proved that for every $u \in H^1(\mathbb{R}^3)$ there is a unique solution $\phi_u \in \mathcal{D}$ of the second equation in the system, that is satisfying

$$(1.2) \quad -\Delta\phi_u + a^2\Delta^2\phi_u = 4\pi u^2.$$

Moreover it turns out that

$$\phi_u = \frac{1 - e^{-|\cdot|/a}}{|\cdot|} * u^2.$$

Observe from (1.2) that, for every $u \in H^1(\mathbb{R}^3)$,

$$4\pi \int \phi_u u^2 = \|\phi_u\|_{\mathcal{D}}^2,$$

which will be used throughout the paper.

By using the classical by now *reduction argument* one is led to study, equivalently, the single equation

$$(1.3) \quad -\Delta u + \omega u + q^2\phi_u u = |u|^{p-2}u \quad \text{in } \mathbb{R}^3$$

containing the nonlocal term $\phi_u u$. Then whenever from now on we speak of solution of the system (1.1) we mean just the solution u of the above equation since $\phi = \phi_u$ is univocally determined. To the equation (1.3) is related the energy functional

$$\mathcal{J}_q(u) = \frac{1}{2}\|u\|^2 + \frac{q^2}{4} \int \phi_u u^2 - \frac{1}{p}\|u\|_p^p, \quad u \in H^1(\mathbb{R}^3)$$

which is well defined and C^1 . In this way we are simply reduced to find critical points of \mathcal{J}_q . We are denoting (here and throughout the paper) by $\|u\|_p$ the L^p -norm and by

$$\|u\|^2 = \|\nabla u\|_2^2 + \omega\|u\|_2^2$$

the (squared) norm in $H^1(\mathbb{R}^3)$, being ω a fixed positive constant.

In [2, Theorem 1.1 and Theorem 1.2] it is proved that if $p \in (3, 6)$, then problem (1.1) admits a solution for every $q \neq 0$. On the other hand, if $p \in (2, 3]$ the existence of a solution is proved just for $q \neq 0$ sufficiently small. As we can see, there is a difference in the result depending on the range where p varies. Indeed in the case of p 's "small" the value of q may prevent the existence of critical points for the functional \mathcal{J}_q .

Of course, if $a = 0$ system (1.1) reduces to the so called Schrödinger-Poisson system in \mathbb{R}^3

$$(1.4) \quad \begin{cases} -\Delta u + \omega u + q^2 \phi u = |u|^{p-2}u, \\ -\Delta \phi = 4\pi u^2 \end{cases}$$

or, equivalently, to

$$-\Delta u + \omega u + q^2 \phi_u^{SP} u = |u|^{p-2}u,$$

where now

$$\phi_u^{SP} = \frac{1}{|\cdot|} * u^2.$$

In the mathematical literature there is a huge number of papers concerning this last problem. However we cite here just [1] where the authors for the first time introduced the *reduction method* which allows to study a single equation instead of a system, and [5] where the author studies the problem depending on the parameter q^2 . In particular Ruiz in [5], among other results, shows that, in the case $p \in (2, 3)$ the system (1.4) has two radial solutions for small q^2 and has no solutions at all, that is radial or not, for $q^2 \geq 1/4$. See also [6] for similar results related to the problem in bounded domains.

Motivated by the cited papers [2, 5], our aim here is to understand in a more satisfactory way the existence of solutions for (1.1), or (1.3), namely the behaviour of \mathcal{J}_q for what concerns its critical points in the case $p \in (2, 3]$ and how they are influenced by the value of q .

We prove two type of results. The first one concerns with the smallness of q^2 as a necessary condition in order to have a nontrivial solution of the problem (the sufficiency being proved in [2]). Indeed we show that for q^2 suitably large the problem has no solutions at all. See Theorem 1.1 below.

The second result concerns the existence of solutions for q^2 small. However, due to the technique used (we borrow some ideas from [5]), we are able to state such a result just in the radial case: in this case obtaining two solutions (in spite of the result in [2] which states the existence of one solution in the nonradial case). See Theorem 1.2 below.

Before stating rigorously the results, we observe that in the problem is appearing the positive parameter q^2 . In view of this, the results are stated and proved for simplicity just for $q > 0$, being understood that they are valid by changing q with $|q|$. More specifically, under the assumption $p \in (2, 3]$ our main results in this work are the following.

Theorem 1.1. *There exists $q^* > 0$ such that, for every $q > q^*$ the problem admits only the trivial solution.*

Theorem 1.2. *There exist $\varepsilon > 0, q_0^* > 0$ satisfying $q_0^* + \varepsilon < q^*$ (with q^* given in Theorem 1.1) such that, for every $q \in (0, q_0^* + \varepsilon)$ the problem has two radial solutions.*

Few comments on these results are in order.

As we already said, we are reduced to find critical points of \mathcal{J}_q . We remark explicitly that no Pohozaev identity is involved in proving the nonexistence result in Theorem 1.1: it just follows due to the properties of the fibering map.

To prove Theorem 1.2 we will use the Mountain Pass Theorem on the space of radial functions, that is in $H_r^1(\mathbb{R}^3)$. We take advantage of the smallness of q to prove that the energy functional has the Mountain Pass Geometry. However, in contrast to [2] where the condition of q small was used in order to find a function where \mathcal{J}_q is negative (and then apply the Mountain Pass Theorem in a standard way), the value q_0^* we find here is a threshold: for $q < q_0^*$ there is a function in $H_r^1(\mathbb{R}^3)$ where the functional is negative, while for $q \geq q_0^*$ the functional is non-negative. Hence the argument employed in both papers [2, 5] do not work for $q > q_0^*$, nevertheless also in this case we exhibit here a Mountain Pass structure.

As shown in [2], the solutions we find in Theorem 1.2 are classical and positive by the Maximum Principle.

Our results holds for any fixed $a \geq 0$. We also notice that, in case $a = 0$ we do not need Lemma 2.1 and the inequality in Proposition 2.2 just follows (up to some positive constant) by multiplying the second equation in (1.4) by $|u|$ and integrating, which is the relation used in [5]. In this sense our result recovers the one in [5] and even a better understanding of the fiber maps is given here, since the Mountain Pass structure for the functional related to (1.4) holds although the functional is non-negative for $q \geq q_0^*$.

Actually we deduce Theorem 1.1 and Theorem 1.2 as consequences of the following result which gives additional information on the solutions.

Theorem 1.3. *Let $p \in (2, 3], a \geq 0$ be fixed. There exist positive numbers ε, q_0, q_0^* satisfying $q_0^* + \varepsilon < q^*$ such that:*

1. *for each $q > q^*$ the functional \mathcal{J}_q has no critical points in $H^1(\mathbb{R}^3)$ other than the zero function;*
2. *for each $q \in (0, q_0^* + \varepsilon)$ the functional \mathcal{J}_q has two nontrivial critical points $u_q, w_q \in H_r^1(\mathbb{R}^3)$ where w_q is a Mountain Pass critical point with*

$$\mathcal{J}_q(w_q) > \max\{0, \mathcal{J}_q(u_q)\}.$$

More specifically,

(i) if $q \in (0, q_0^*]$, then u_q is a global minimum, with

$$\mathcal{J}_q(u_q) < 0 \quad \text{if } q \in (0, q_0^*), \quad \mathcal{J}_{q_0^*}(u_{q_0^*}) = 0;$$

(ii) if $q \in (q_0^*, q_0^* + \varepsilon)$, then u_q is a local minimum with

$$\mathcal{J}_q(u_q) > 0.$$

As we can see, whenever $q = q_0^*$, then $\mathcal{J}_{q_0^*}$ is non-negative and we find a global minimizer at zero energy; then an additional work is needed in order to show that this is not the zero function. This result is new also in the case $a = 0$.

Concerning the extremal values q_0^* and q^* we say here they have a variational characterisation (see Section 2). Furthermore, although we were not able to prove it, it seems plausible that for all $q \in (0, q^*)$, the system (1.1) has two positive solutions satisfying the properties in Theorem 1.3 and q^* is in fact a bifurcation parameter where the two solutions (the local minimum and the Mountain Pass type solution) collapses.

We point out finally that similar results have been obtained in some nonlinear problems depending on a parameter in some recent papers: see Il'yasov and Silva [4], Silva and Macedo [7].

This paper is organised as follows.

In Section 2 some preliminaries and technical results (true in the general nonradial setting) are given. As a byproduct of these results, the proof of item 1. in Theorem 1.3 follows, see Corollary 2.8.

In Section 3 the important Proposition 3.1 is proved. It concerns with radial functions and is fundamental in proving our result.

Finally, the proof of Theorem 1.3 is completed in Section 4.

As a matter of notations, we use the generic letters C, C', \dots to denote a positive constant, usually related to Sobolev embedding, whose value may also change from line to line: no confusion should arise.

2. PRELIMINARIES AND TECHNICAL RESULTS

In [2] some properties of the solution ϕ_u are found. However to deal with the case $p \in (2, 3]$ we need also the next ones. Of course this applies just for $a \neq 0$.

Lemma 2.1. *For each $u \in H^1(\mathbb{R}^3)$ we have*

$$(i) \quad \Delta \phi_u \in \mathcal{D},$$

$$(ii) \quad a^2 \Delta \phi_u \leq \phi_u.$$

Proof. Let us fix $u \in H^1(\mathbb{R}^3)$ and let $\psi_u := -\Delta\phi_u$. Then

$$-a^2\Delta\psi_u + \psi_u = 4\pi u^2.$$

Since $u^2 \in L^2(\mathbb{R}^3)$, by standard results we have $\psi_u \in H^2(\mathbb{R}^3)$ and in particular,

$$\psi_u \in L^6(\mathbb{R}^3), \quad \nabla\psi_u \in L^2(\mathbb{R}^3) \quad \text{and} \quad \Delta\psi_u \in L^2(\mathbb{R}^3).$$

This gives $\Delta\phi_u \in D^{1,2}(\mathbb{R}^3)$ and $\Delta^2\phi_u \in L^2(\mathbb{R}^3)$, namely $\Delta\phi_u \in \mathcal{D}$ proving (i).

On the other hand, if we set $v = -a^2\Delta\phi_u + \phi_u$, then

$$-\Delta v = 4\pi u^2 \geq 0$$

and $v \in D^{1,2}(\mathbb{R}^3)$ and is continuous. Define $\Omega^- = \{x \in \mathbb{R}^3 : v(x) < 0\}$ and suppose that $\Omega^- \neq \emptyset$. Once v is continuous, the set Ω^- is open. Let $v^- = \max\{-v, 0\}$. It follows that

$$-\int_{\Omega^-} |\nabla v|^2 = \int \nabla v \nabla v^- \geq 0,$$

which is a contradiction, therefore, $\Omega^- = \emptyset$ and $a^2\Delta\phi_u \leq \phi_u$ in \mathbb{R}^3 , proving (ii). \square

The next result will be useful to get a generalisation of [5, Formula (19)] to the case $a \neq 0$.

Proposition 2.2. *There holds*

$$\int |u|^3 \leq \frac{1}{\pi} \|\phi_u\|_{\mathcal{D}} \|\nabla u\|_2, \quad \forall u \in H^1(\mathbb{R}^3).$$

Proof. For $u \in H^1(\mathbb{R}^3)$ fixed, let us consider equation (1.2). Since by Lemma 2.1 it holds in particular that $\nabla\Delta\phi_u \in L^2(\mathbb{R}^3)$, by multiplying the equation (1.2) by $|u| \in H^1(\mathbb{R}^3)$ and integrating we get

$$\begin{aligned} 4\pi \int |u|^3 &= a^2 \int \nabla(-\Delta\phi_u) \nabla|u| + \int \nabla\phi_u \nabla|u| \\ &\leq a^2 \|\nabla(-\Delta\phi_u)\|_2 \|\nabla u\|_2 + \|\nabla\phi_u\|_2 \|\nabla u\|_2 \\ &= (a^2 \|\nabla(-\Delta\phi_u)\|_2 + \|\nabla\phi_u\|_2) \|\nabla u\|_2 \\ (2.1) \quad &\leq (a^2 \|\nabla(\Delta\phi_u)\|_2 + \|\phi_u\|_{\mathcal{D}}) \|\nabla u\|_2. \end{aligned}$$

On the other hand by multiplying (1.2) by $\Delta\phi_u \in \mathcal{D}$ and making use of (ii) of Lemma 2.1 we get

$$\begin{aligned} a^2 \|\nabla(\Delta\phi_u)\|_2^2 &= 4\pi \int \Delta\phi_u u^2 - \int \nabla\phi_u \nabla(\Delta\phi_u) \\ &\leq \frac{1}{a^2} \|\phi_u\|_{\mathcal{D}}^2 + \frac{\varepsilon^2}{2} \|\nabla(\Delta\phi_u)\|_2^2 + \frac{1}{2\varepsilon^2} \|\nabla\phi_u\|_2^2 \\ &\leq \frac{1}{a^2} \|\phi_u\|_{\mathcal{D}}^2 + \frac{\varepsilon^2}{2} \|\nabla(\Delta\phi_u)\|_2^2 + \frac{1}{2\varepsilon^2} \|\phi_u\|_{\mathcal{D}}^2. \end{aligned}$$

By choosing $\varepsilon = a$ above we conclude that, for all $u \in H^1(\mathbb{R}^3)$,

$$(2.2) \quad a^2 \|\nabla(\Delta\phi_u)\|_2^2 \leq \frac{2}{a^2} \|\phi_u\|_{\mathcal{D}}^2 + \frac{1}{a^2} \|\phi_u\|_{\mathcal{D}}^2 = \frac{3}{a^2} \|\phi_u\|_{\mathcal{D}}^2.$$

From (2.1) and (2.2) we conclude that

$$\int |u|^3 \leq \frac{1}{\pi} \|\phi_u\|_{\mathcal{D}} \|\nabla u\|_2,$$

completing the proof. \square

We conclude this section by showing a first simple property of the energy functional. The next result says that the functional \mathcal{J}_q has a strict local minimum in 0, uniformly in q . However to have the complete Mountain Pass structure q has to be small, as it will be shown in Corollary 2.9.

Proposition 2.3. *There exist $\rho > 0$ and $M > 0$ such that*

$$\forall q \in \mathbb{R}, u \in H^1(\mathbb{R}^3) \text{ with } \|u\| = \rho : \mathcal{J}_q(u) \geq M.$$

Proof. Since

$$\mathcal{J}_q(u) \geq \frac{1}{2} \|u\|^2 - \frac{1}{p} \|u\|_p^p \geq \frac{1}{2} \|u\|^2 - C \|u\|^p,$$

the conclusion easily follows. \square

In this Section we establish some notations and study the geometry of the functional \mathcal{J}_q . We observe that $\phi_{tu} = t^2 \phi_u$ and therefore, if $\psi_{q,u} : [0, \infty) \rightarrow \mathbb{R}$ is defined by $\psi_{q,u}(t) = \mathcal{J}_q(tu)$, we have that

$$\psi_{q,u}(t) = \frac{t^2}{2} \|u\|^2 + \frac{q^2 t^4}{4} \int \phi_u u^2 - \frac{t^p}{p} \|u\|_p^p.$$

Whenever q and u are fixed, we will use for brevity also the notation $\psi := \psi_{q,u}$. A simple analysis shows that

Proposition 2.4. *For each $q \in \mathbb{R} \setminus \{0\}$ and $u \in H^1(\mathbb{R}^3) \setminus \{0\}$, there are only three possibilities for the graph of ψ :*

- (i) *the function ψ has only two critical points when $t > 0$, to wit, $0 < t_q^-(u) < t_q^+(u)$. Moreover, $t_q^-(u)$ is a local maximum with $\psi''(t_q^-(u)) < 0$ and $t_q^+(u)$ is a local minimum with $\psi''(t_q^+(u)) > 0$;*
- (ii) *the function ψ has only one critical point when $t > 0$ at the value $t_q(u)$. Moreover, $\psi''(t_q(u)) = 0$ and ψ is increasing;*
- (iii) *the function ψ is increasing and has no critical points.*

It is important to notice that (i) happens for q small, and (iii) for q large.

Let us consider the Nehari manifold associated with the functional \mathcal{J}_q , that is

$$\mathcal{N}_q = \{u \in H^1(\mathbb{R}^3) \setminus \{0\} : \psi'_{q,u}(1) = 0\}.$$

Note that, for $u \in \mathcal{N}_q$:

$$\|u\|^2 \leq \|u\|^2 + q^2 \int \phi_u u^2 \leq C \|u\|^p$$

and then all the Nehari manifolds are bounded away from zero uniformly in q , in the sense that

$$(2.3) \quad \exists \tilde{C} > 0 \text{ such that for all } q \in \mathbb{R}, u \in \mathcal{N}_q : \|u\| \geq \tilde{C}.$$

Moreover since

$$\mathcal{N}_q = \mathcal{N}_q^+ \cup \mathcal{N}_q^0 \cup \mathcal{N}_q^-,$$

where

$$\begin{aligned} \mathcal{N}_q^+ &= \{u \in \mathcal{N}_q : \psi''(1) > 0\}, \\ \mathcal{N}_q^0 &= \{u \in \mathcal{N}_q : \psi''(1) = 0\}, \\ \mathcal{N}_q^- &= \{u \in \mathcal{N}_q : \psi''(1) < 0\}. \end{aligned}$$

as an application of the Implicit Function Theorem one has the following:

Proposition 2.5. *If $\mathcal{N}_q^+, \mathcal{N}_q^- \neq \emptyset$, then $\mathcal{N}_q^+, \mathcal{N}_q^- \neq \emptyset$ are C^1 manifolds of codimension 1 in $H^1(\mathbb{R}^3)$. Moreover, $u \in \mathcal{N}_q^+ \cup \mathcal{N}_q^-$ is a critical point for the functional \mathcal{J}_q if and only if u is a critical point of the constrained functional $(\mathcal{J}_q)|_{\mathcal{N}_q^+ \cup \mathcal{N}_q^-} : \mathcal{N}_q^+ \cup \mathcal{N}_q^- \rightarrow \mathbb{R}$.*

Note that, fixed $u \in H^1(\mathbb{R}^3) \setminus \{0\}$ we have $tu \in \mathcal{N}_q^0$ if, and only if $\psi'_{q,tu}(1) = \psi''_{q,tu}(1) = 0$, i.e. the following system of equations is satisfied:

$$(2.4) \quad \begin{cases} t\|u\|^2 + q^2 t^3 \int \phi_u u^2 - t^{p-1} \|u\|_p^p = 0, \\ \|u\|^2 + 3q^2 t^2 \int \phi_u u^2 - (p-1)t^{p-2} \|u\|_p^p = 0. \end{cases}$$

We can solve the system (2.4) with respect to the variables q and t to obtain a unique solution given by

$$t(u) = \left(\frac{2\|u\|^2}{(4-p)\|u\|_p^p} \right)^{1/(p-2)}$$

and

$$(2.5) \quad q(u) = C_p \frac{\|u\|_p^{p/(p-2)}}{\|u\|^{(4-p)/(p-2)} \|\phi_u\|_{\mathcal{D}}}, \quad C_p = \frac{2(p-2)^{1/2} \pi^{1/2} (4-p)^{(4-p)/2(p-2)}}{2^{1/(p-2)}}.$$

In addition the solutions $q(u)$ and $t(u)$ are related by

$$t(u) = \left(\frac{2q^2(u)}{4\pi(p-2)} \frac{\|\phi_u\|_{\mathcal{D}}^2}{\|u\|_p^p} \right)^{1/(p-4)}.$$

Define the extremal value (see Il'yasov [3])

$$q^* = \sup \{q(u) : u \in H^1(\mathbb{R}^3) \setminus \{0\}\}.$$

Lemma 2.6. *The function $H^1(\mathbb{R}^3) \setminus \{0\} \ni u \mapsto q(u)$ defined in (2.5) is 0-homogeneous. Moreover $q^* < \infty$.*

Proof. That $q(u)$ is zero homogeneous is obvious. Let us prove that $q^* < \infty$. Indeed, once $p \in (2, 3]$ we have from the interpolation inequality that, for all $u \in H^1(\mathbb{R}^3)$ we have

$$(2.6) \quad \|u\|_p^p \leq \|u\|^{6-2p} \|u\|_3^{3p-6}.$$

Combining the inequality (2.6) with the Proposition 2.2 we conclude that

$$(2.7) \quad \|u\|_p^p \leq C \|u\|^{6-2p} \|u\|^{(3p-6)/3} \|\phi_u\|_{\mathcal{D}}^{(3p-6)/3} = C \|u\|^{4-p} \|\phi_u\|_{\mathcal{D}}^{p-2}.$$

for some constant $C > 0$. It follows from (2.7) that

$$q(u) \leq C \frac{\|u\|^{(4-p)/(p-2)} \|\phi_u\|_{\mathcal{D}}}{\|u\|^{(4-p)/(p-2)} \|\phi_u\|_{\mathcal{D}}} \leq C$$

completing the proof. \square

Another extremal value which will be important for us is the one such that, for larger values of the parameter, the functional is always non-negative. Let us start by fixing $u \in H^1(\mathbb{R}^3) \setminus \{0\}$ and considering the system

$$\begin{cases} \psi_{q_0, u}(t_0) = \frac{t_0^2}{2} \|u\|^2 + q_0^2 \frac{t_0^4}{4} \int \phi_u u^2 - \frac{t_0^p}{p} \|u\|_p^p = 0, \\ \psi'_{q_0, u}(t_0) = t_0 \|u\|^2 + q_0^2 t_0^3 \int \phi_u u^2 - t_0^{p-1} \|u\|_p^p = 0. \end{cases}$$

One can solve this system with respect to the variables t_0 and q_0 to obtain the unique solution given by

$$(2.8) \quad q_0(u) = C_{0,p} \frac{\|u\|_p^{p/(p-2)}}{\|u\|^{(4-p)/(p-2)} \|\phi_u\|_{\mathcal{D}}}, \quad C_{0,p} = \frac{2^{3/2}(p-2)^{1/2} \pi^{1/2} (4-p)^{(4-p)/2(p-2)}}{p^{1/(p-2)}}$$

and $t_0(u)$ is given by

$$t_0(u) = \left(\frac{pq_0^2(u)}{2(p-2)} \frac{\|\phi_u\|_{\mathcal{D}}^2}{\|u\|_p^p} \right)^{1/(p-4)}.$$

Observe that $C_{0,p} < C_p$, where C_p is the one appearing in (2.5). Then $q_0(u) < q(u)$. Define the extremal value as

$$q_0^* = \sup \{q_0(u) : u \in H^1(\mathbb{R}^3) \setminus \{0\}\}.$$

Remark 1. *Once $q_0(u)$ is a multiple of $q(u)$, Lemma 2.6 also holds true for the function q_0 .*

The solutions $q(u)$ and $q_0(u)$ given in (2.5) and (2.8) have the following geometrical interpretation which can be proved starting from the Proposition 2.4.

Proposition 2.7. *For each $u \in H^1(\mathbb{R}^3) \setminus \{0\}$ there holds:*

- (i) $q(u)$ is the unique parameter $q > 0$ for which the fiber map $\psi_{q,u}$ has a critical point with second derivative zero at $t(u)$. Moreover, if $0 < q < q(u)$, then $\psi_{q,u}$ satisfies (i) of the Proposition 2.4 while if $q > q(u)$, then $\psi_{q,u}$ satisfies (iii) of the Proposition 2.4.
- (ii) $q_0(u)$ is the unique parameter $q > 0$ for which the fiber map $\psi_{q,u}$ has a critical point with zero energy at $t_0(u)$. Moreover, if $0 < q < q_0(u)$, then $\inf_{t>0} \psi_{q,u}(t) < 0$ while if $q > q_0(u)$, then $\inf_{t>0} \psi_{q,u}(t) = 0$.

Moreover the parameter q_0^* has the geometrical interpretation that for each $q \in (0, q_0^*)$, there exists at least one $u \in H^1(\mathbb{R}^3) \setminus \{0\}$ for which $\mathcal{J}_q(u) < 0$, while if $q \geq q_0^*$, then $\mathcal{J}_q(u) \geq 0$ for all $u \in H^1(\mathbb{R}^3)$. In both works [2, 5] the necessity of small values of q was imposed in order to show that there exists a function where the functional is negative, in such a way that \mathcal{J}_q possesses a Mountain Pass Geometry. Therefore, the argument employed in both papers do not work for $q > q_0^*$.

The above proposition has the following important consequences.

Corollary 2.8. *If $q > q^*$ the functional \mathcal{J}_q has no critical points other than the zero function. Moreover if $q < q^*$, then $\mathcal{N}_q^- \neq \emptyset$ and $\mathcal{N}_q^+ \neq \emptyset$.*

In particular item 1. of Theorem 1.3 is proved.

Proof. It is sufficient to show that, for each $u \in H^1(\mathbb{R}^3) \setminus \{0\}$, the function $\psi_{q,u}$ has no critical points for $q > q^*$. Actually this is a consequence of the inequalities $q(u) \leq q^* < q$ and (i) of Proposition 2.7.

Now assume that $q < q^*$. From the definition of q^* , there exists $u \in H^1(\mathbb{R}^3) \setminus \{0\}$ such that $q < q(u) < q^*$. Therefore, from (i) of Proposition 2.7 we conclude that $\mathcal{N}_q^- \neq \emptyset$ and $\mathcal{N}_q^+ \neq \emptyset$. \square

Corollary 2.9. *For each $q \geq q_0^*$, there holds $\mathcal{J}_q(u) \geq 0$ for all $u \in H^1(\mathbb{R}^3)$. Moreover, if $q < q_0^*$, then there exists $u \in H^1(\mathbb{R}^3)$ such that $\mathcal{J}_q(u) < 0$.*

Proof. Indeed, assume that $q \geq q_0^*$. It follows that $q > q_0(u)$ for each $u \in H^1(\mathbb{R}^3) \setminus \{0\}$ and from item (ii) of Proposition 2.7 there holds $\inf_{t>0} \psi_{q,u}(t) = 0$. Therefore $\mathcal{J}_q(u) \geq 0$.

Now assume that $q < q_0^*$. From the definition of q_0^* , there exists $w \in H^1(\mathbb{R}^3) \setminus \{0\}$ such that $q < q_0(w) < q_0^*$. Therefore, from (ii) of Proposition 2.7 we conclude that $\inf_{t>0} \psi_{q,w}(t) < 0$ and hence there exists $t > 0$ such that if $u := tw$, it holds $\mathcal{J}_q(u) < 0$. \square

Let us conclude this section with the following important result.

Proposition 2.10. *There exists a positive constant m such that*

$$\forall q \in \mathbb{R}, u \in \mathcal{N}_q^0 : \mathcal{J}_q(u) \geq m.$$

Proof. From the equations (2.4) with $t = 1$ we have that

$$\begin{cases} \|u\|^2 + q^2 \int \phi_u u^2 - \|u\|_p^p = 0, \\ \|u\|^2 + 3q^2 \int \phi_u u^2 - (p-1)\|u\|_p^p = 0. \end{cases}$$

It follows that

$$q^2 \int \phi_u u^2 = \|u\|_p^p - \|u\|^2 \quad \text{and} \quad \|u\|_p^p = \frac{2}{4-p} \|u\|^2$$

so that $\mathcal{J}_q(u) = \frac{p-2}{4p} \|u\|^2$ for each $u \in \mathcal{N}_q^0$. From (2.3) the proof is completed. \square

It is worth to point out that all that we have done in this section does not use the radial setting, and clearly these results also holds in $H_r^1(\mathbb{R}^3)$.

3. GLOBAL MINIMA AND (PS) SEQUENCES FOR \mathcal{J}_q

In this section we prove the following result. Here and in the next section is fundamental to work with radial functions.

Proposition 3.1. *There holds:*

- (i) for each $q \in (0, q_0^*)$ we have that $-\infty < \inf_{u \in H_r^1(\mathbb{R}^3)} \mathcal{J}_q(u) < 0$;
- (ii) for each $q > 0$, if $\{u_n\} \subset H_r^1(\mathbb{R}^3)$ is a sequence such that $\mathcal{J}'_q(u_n) \rightarrow 0$ as $n \rightarrow \infty$, then $\{u_n\}$ is convergent, up to subsequences.

Proof. Let us show (i). Indeed, from the Corollary 2.9 we know that $\inf_{u \in H_r^1(\mathbb{R}^3)} \mathcal{J}_q(u) < 0$. We claim that $-\infty < \inf_{u \in H_r^1(\mathbb{R}^3)} \mathcal{J}_q(u)$. In fact, given $\varepsilon > 0$ such that $D := \frac{q^2}{16\pi} - \varepsilon^4 > 0$, by Proposition 2.2, we have

that for each $u \in H_r^1(\mathbb{R}^3)$ there holds

$$\begin{aligned}
\mathcal{J}_q(u) &= \frac{1}{4}\|\nabla u\|_2^2 + \frac{1}{4}\|\nabla u\|_2^2 + \frac{1}{2}\|u\|_2^2 + \frac{q^2}{4} \int \phi_u u^2 - \frac{1}{p}\|u\|_p^p \\
&\geq \frac{1}{4}\|\nabla u\|_2^2 + D\|\phi_u\|_{\mathcal{D}}^2 + \frac{1}{2}\|u\|_2^2 + \frac{\pi\varepsilon^2}{4}\|u\|_3^3 - \frac{1}{p}\|u\|_p^p \\
(3.1) \quad &= \frac{1}{4}\|u\|^2 + D\|\phi_u\|_{\mathcal{D}}^2 + \int f(u)
\end{aligned}$$

where

$$f(t) = \frac{1}{4}t^2 + \frac{\pi\varepsilon^2}{4}t^3 - \frac{1}{p}t^p, \quad \forall t > 0.$$

A simple analysis shows that $I := \inf_{t>0} f(t) > -\infty$ and if $f(t) < 0$ for some $t > 0$, then $f^{-1}((-\infty, 0)) = (\alpha, \beta)$, where $0 < \alpha < \beta < \infty$.

If $I \geq 0$, being $D_q > 0$, from (3.1) we conclude that $-\infty < \inf_{u \in H_r^1(\mathbb{R}^3)} \mathcal{J}_q(u)$.

If $I < 0$ then

$$(3.2) \quad \mathcal{J}_q(u) \geq \frac{1}{4}\|u\|^2 + D_q\|\phi_u\|_{\mathcal{D}}^2 + I \text{meas}(A)$$

where $A = \{x \in \mathbb{R}^3 : u(x) \in (\alpha, \beta)\}$. If there exists a sequence $\{u_n\} \subset H_r^1(\mathbb{R}^3)$ such that $\mathcal{J}_q(u_n) \rightarrow -\infty$ as $n \rightarrow \infty$, then $\|u_n\| \rightarrow +\infty$. Moreover by (3.2) we can assume without loss of generality that

$$(3.3) \quad \frac{1}{4}\|u_n\|^2 < |I| \text{meas}(A_n), \quad \forall n \in \mathbb{N}.$$

By the result of Strauss [8] we know that there exist a positive constant C such that

$$(3.4) \quad |u(x)| \leq C|x|^{-1}\|u\|, \quad \forall u \in H_r^1(\mathbb{R}^3).$$

Define $\rho_n = \sup\{|x| : x \in A_n\}$ and observe from the inequalities (3.3) and (3.4) that, for every $x \in \mathbb{R}^3$ with $|x| = \rho_n$ it holds,

$$0 < \alpha \leq u_n(x) \leq C\rho_n^{-1}\|u_n\| \leq 2C\rho_n^{-1}(|I|\text{meas}(A_n))^{1/2},$$

and hence, for some $C' > 0$, we deduce

$$(3.5) \quad C'\rho_n \leq \text{meas}(A_n)^{1/2}.$$

Similar to the deduction of (3.3), we can assume without loss of generality that

$$D\|\phi_{u_n}\|_{\mathcal{D}}^2 < |I|\text{meas}(A_n), \quad \forall n \in \mathbb{N},$$

and hence, once that the function $(0, \infty) \ni t \mapsto (1 - e^{-t/a})/t$ is decreasing we conclude that

$$\begin{aligned}
|I|meas(A_n) &> D\|\phi_{u_n}\|_{\mathcal{D}}^2 \\
&= \int \int \frac{1 - e^{-\frac{|x-y|}{a}}}{|x-y|} u_n^2(x)u_n^2(y) \\
&\geq \int_{A_n} \int_{A_n} \frac{1 - e^{-\frac{|x-y|}{a}}}{|x-y|} u_n^2(x)u_n^2(y) \\
&\geq \frac{1 - e^{-\frac{2\rho_n}{a}}}{2\rho_n} \alpha^4 meas(A_n)^2,
\end{aligned}$$

which implies that

$$(3.6) \quad \frac{|I|}{\alpha^4} \geq \frac{1 - e^{-\frac{2\rho_n}{a}}}{2\rho_n} meas(A_n), \quad \forall n \in \mathbb{N},$$

Observe from (3.5) that $|A_n| \rightarrow \infty$ as $n \rightarrow \infty$ and from (3.6) that $\rho_n \rightarrow \infty$ as $n \rightarrow \infty$. Combining (3.5) with (3.6) we obtain that

$$C'' \geq (1 - e^{-\frac{2\rho_n}{a}})\rho_n,$$

for some $C'' > 0$, which is clearly a contradiction and therefore (i) is proved.

Let us show (ii). From the convergence $\mathcal{J}'_q(u_n) \rightarrow 0$ as $n \rightarrow \infty$, we can assume without loss of generality that

$$(3.7) \quad \mathcal{J}'_q(u_n)[u_n] \leq \|u_n\|, \quad \forall n \in \mathbb{N}.$$

On the other hand, from Proposition 2.2 we have that

$$\|u_n\|_3^3 \leq \frac{1}{\pi} \left(\frac{1}{\varepsilon^2} \|\nabla u_n\|_2^2 + \varepsilon^2 \|\phi_{u_n}\|_{\mathcal{D}}^2 \right),$$

where $\varepsilon > 0$ is chosen now such that $\frac{q^2}{4\pi} - \frac{\varepsilon^4}{2} > 0$. It follows that, for all $n \in \mathbb{N}$:

$$\begin{aligned}
\mathcal{J}'_q(u_n)[u_n] &= \|\nabla u_n\|^2 + \|u_n\|_2^2 + \frac{q^2}{4\pi} \|\phi_{u_n}\|_{\mathcal{D}}^2 - \|u_n\|_p^p \\
&\geq \frac{1}{2} \|u_n\|^2 + \frac{1}{2} \|u_n\|_2^2 + \frac{\pi\varepsilon^2}{2} \|u_n\|_3^3 - \frac{\varepsilon^4}{2} \|\phi_{u_n}\|_{\mathcal{D}}^2 + \frac{q^2}{4\pi} \|\phi_{u_n}\|_{\mathcal{D}}^2 - \|u_n\|_p^p \\
&= \frac{1}{2} \|u_n\|^2 + \left(\frac{q^2}{4\pi} - \frac{\varepsilon^4}{2} \right) \|\phi_{u_n}\|_{\mathcal{D}}^2 + \frac{1}{2} \|u_n\|_2^2 + \frac{\pi\varepsilon^2}{2} \|u_n\|_3^3 - \|u_n\|_p^p \\
(3.8) \quad &\frac{1}{2} \|u_n\|^2 + \left(\frac{q^2}{4\pi} - \frac{\varepsilon^4}{2} \right) \|\phi_{u_n}\|_{\mathcal{D}}^2 + \int g(u_n),
\end{aligned}$$

where $g(t) = t^2/2 + \frac{\pi\varepsilon^2}{2}t^3 - t^p$ for $t > 0$. We combine (3.7) with (3.8) to conclude that

$$\|u_n\| \geq \frac{1}{2}\|u_n\|^2 + \left(\frac{q^2}{4\pi} - \frac{\varepsilon^4}{2}\right)\|\phi_{u_n}\|_{\mathcal{D}}^2 + \int g(u_n).$$

We conclude as in the proof of (i) that $\{u_n\}$ is bounded. Once we know that $\{u_n\}$ is bounded, standard arguments (observe that the analogous of [5, Lemma 2.1] is valid) produce a convergent subsequence. \square

Remark 2. Note that (ii) in the Proposition 3.1 can be extended in the following way: if $q_n \rightarrow q$ and if $\{u_n\} \subset H_r^1(\mathbb{R}^3)$ is a sequence such that $\mathcal{J}'_{q_n}(u_n) \rightarrow 0$ as $n \rightarrow \infty$, then $\{u_n\}$ is convergent, up to subsequences. This follows due to the smooth dependence of \mathcal{J}'_q on q .

4. EXISTENCE OF TWO RADIAL SOLUTIONS

In this section we prove item 2. of Theorem 1.3.

Proposition 4.1. For each $q \in (0, q_0^*)$ there exists a global minimum u_q such that $\mathcal{J}_q(u_q) < 0$.

Proof. It follows from the Proposition 3.1 and Ekeland's Variational Principle. \square

Now we prove the existence of a local minimizer for \mathcal{J}_q when q is near q_0^* . To do so, we first prove the existence of a global minimizer for the functional $\mathcal{J}_{q_0^*}$.

Corollary 4.2. There exists a global minimizer $u_{q_0^*} \neq 0$ of $\mathcal{J}_{q_0^*}$ such that $\mathcal{J}_{q_0^*}(u_{q_0^*}) = 0$.

Proof. Indeed, suppose that $q_n \uparrow q_0^*$ as $n \rightarrow \infty$. From the Proposition 4.1, for each n , there exists $u_n := u_{q_n}$ such that u_n is a global minimum for \mathcal{J}_{q_n} and $\mathcal{J}_{q_n}(u_n) < 0$. It follows that $\mathcal{J}'_{q_n}(u_n) = 0$ for each n and, being all the Nehari manifolds bounded away from zero uniformly in q , by (2.3) it is $\|u_n\| \geq \tilde{C}$ for each n . From (ii) in Proposition 3.1 (see Remark 2) we conclude that $u_n \rightarrow u \neq 0$. By $\mathcal{J}_{q_n}(u_n) < 0$ for each n , we conclude that $\mathcal{J}_{q_0^*}(u) \leq 0$ and from Corollary 2.9 it follows that $\mathcal{J}_{q_0^*}(u) = 0$. Then it is sufficient to set $u_{q_0^*} := u$ and the proof is completed. \square

Remark 3. From the definition of q_0^* and the Corollary 4.2 it follows that $q_0(u_{q_0^*}) = q_0^*$. Moreover $q_0^* < q(u_{q_0^*})$.

For $q > 0$, define

$$\widehat{\mathcal{J}}_q := \inf \{ \mathcal{J}_q(u) : u \in \mathcal{N}_q^+ \cup \mathcal{N}_q^0 \}.$$

Observe that

$$(4.1) \quad \widehat{\mathcal{J}}_q = \inf_{u \in H_r^1(\mathbb{R}^3)} \mathcal{J}_q(u) \quad \forall q \in (0, q_0^*]$$

and from Corollary 2.9 there holds $\widehat{\mathcal{J}}_q \geq 0$ for $q \geq q_0^*$.

Proposition 4.3. *Given $\delta > 0$, there exists $\varepsilon > 0$ such that for each $q \in (q_0^*, q_0^* + \varepsilon)$ there holds $\widehat{\mathcal{J}}_q < \delta$.*

Proof. Indeed, let $u_{q_0^*} \in \mathcal{N}_q^+$ be given as in Corollary 4.2. Observe that if $q \downarrow q_0^*$, then $\mathcal{J}_q(u_{q_0^*}) \rightarrow \mathcal{J}_{q_0^*}(u_{q_0^*}) = 0$. Moreover, once $q_0^* < q(u_{q_0^*})$, it follows that there exists $\varepsilon_1 > 0$ such that $q_0^* + \varepsilon_1 < q(u_{q_0^*})$. From Proposition 2.4 and (ii) in Proposition 2.7, for each $q \in (q_0^*, q_0^* + \varepsilon_1)$, there exists $t_q^+(u_{q_0^*})$ such that $t_q^+(u_{q_0^*})u_{q_0^*} \in \mathcal{N}_q^+$. Note that $t_q^+(u_{q_0^*}) \rightarrow 1$ as $q \downarrow q_0^*$ and therefore

$$\mathcal{J}_q(u_{q_0^*}) \leq \mathcal{J}_q(t_q^+(u_{q_0^*})u_{q_0^*}) \rightarrow \mathcal{J}_{q_0^*}(u_{q_0^*}) = 0, \quad q \downarrow q_0^*.$$

If $\varepsilon_2 > 0$ is chosen in such a way that $\mathcal{J}_q(t_q^+(u_{q_0^*})u_{q_0^*}) < \delta$ for each $q \in (q_0^*, q_0^* + \varepsilon_2)$, then we set $\varepsilon = \min\{\varepsilon_1, \varepsilon_2\}$ and the proof is completed. \square

Let us recall that by Proposition 2.3, there exist positive constants ρ, M such that $\mathcal{J}_q(u) \geq M$ for each $\|u\| = \rho$. We can assume without loss of generality that $\rho < \widetilde{C}$, where \widetilde{C} is such that

$$\text{for all } q \in \mathbb{R}, u \in \mathcal{N}_q : \|u\| \geq \widetilde{C},$$

(see (2.3)).

We choose $\delta > 0$ in the Proposition 4.3 in such a way that

$$(4.2) \quad \delta < \min\{M, m\},$$

where m is the positive constant such that, by Proposition 2.10,

$$\forall q \in \mathbb{R}, u \in \mathcal{N}_q^0 : \mathcal{J}_q(u) \geq m.$$

Let $\varepsilon > 0$ be as in Proposition 4.3 in correspondence of the above fixed $\delta > 0$.

Proposition 4.4. *There holds*

$$\inf \{ \mathcal{J}_q(u) : \|u\| \geq \rho \} = \widehat{\mathcal{J}}_q, \quad \forall q \in (q_0^*, q_0^* + \varepsilon).$$

Proof. First observe from the inequality $\rho < \widetilde{C}$ that $\inf \{ \mathcal{J}_q(u) : \|u\| \geq \rho \} \leq \widehat{\mathcal{J}}_q$. We claim that the equality holds. Indeed, by one hand, if the fiber map $\psi_{q,u}$ satisfies (ii) or (iii) of the Proposition 2.4, then $\inf_{t>\rho} \psi_{q,u}(t) = M$. On the other hand, if the fiber map $\psi_{q,u}$ satisfies (i) of the Proposition 2.4, then $\inf_{t>\rho} \psi_{q,u}(t) \geq \widehat{\mathcal{J}}_q$. Once $M > \delta > \widehat{\mathcal{J}}_q$ the proof is completed. \square

Corollary 4.5. *For each $q \in (q_0^*, q_0^* + \varepsilon)$ there holds: there exists $u_q \in \mathcal{N}_q^+$ such that $\mathcal{J}_q(u_q) = \widehat{\mathcal{J}}_q$. In particular $\mathcal{J}_q(u_q) > 0$ and $\|u_q\| \geq \widetilde{C} > \rho$.*

Proof. Fix $q \in (q_0^*, q_0^* + \varepsilon)$ and let $\{u_n\} \subset \mathcal{N}_q^+ \cup \mathcal{N}_q^0$ be a minimising sequence for $\widehat{\mathcal{J}}_q < \delta$ by Proposition 4.3. Since $m > \delta$ and, by Proposition 2.10, $\mathcal{J}_q(u) \geq m$ on \mathcal{N}_q^0 , we can assume that $\{u_n\} \subset \mathcal{N}_q^+$ and hence, by the Ekeland's Variational Principle, also that $\mathcal{J}'_q(u_n) \rightarrow 0$. We conclude from the Proposition 3.1 that $u_n \rightarrow u$ in $H_r^1(\mathbb{R}^3)$ with $\|u\| \geq \widetilde{C} > \rho$. Setting $u_q := u$ clearly we obtain that $u_q \in \mathcal{N}_q^+$ and $\mathcal{J}_q(u_q) = \widehat{\mathcal{J}}_q$. Due to the definition of q_0^* and the fact that $q > q_0^*$, we conclude that $\mathcal{J}_q(u_q) > 0$. \square

We observe two properties of the function $(0, q_0^* + \varepsilon) \ni q \mapsto \widehat{\mathcal{J}}_q$.

Lemma 4.6. *The function $(0, q_0^* + \varepsilon) \ni q \mapsto \widehat{\mathcal{J}}_q$ is increasing and continuous.*

Proof. Indeed, suppose that $q < q'$. From Proposition 4.1, Corollary 4.2 and the Corollary 4.5 there exists $u_{q'}$ such that $\widehat{\mathcal{J}}_{q'} = \mathcal{J}_{q'}(u_{q'})$.

If $q' \in (q_0^*, q_0^* + \varepsilon)$, from the Corollary 4.5 it is $\|u_{q'}\| \geq \widetilde{C} > \rho$ and hence from the Proposition 4.4 we obtain

$$\widehat{\mathcal{J}}_q \leq \mathcal{J}_q(u_{q'}) < \mathcal{J}_{q'}(u_{q'}) = \widehat{\mathcal{J}}_{q'}.$$

If $q' \in (0, q_0^*]$ the lemma follows by (4.1).

Now we prove that $(0, q_0^* + \varepsilon) \ni q \mapsto \widehat{\mathcal{J}}_q$ is continuous. In fact, suppose that $q_n \uparrow q \in (0, q_0^* + \varepsilon)$. From Proposition 4.1, Corollary 4.2 and Corollary 4.5, for each n , there exists $u_n := u_{q_n}$ such that $\widehat{\mathcal{J}}_{q_n} = \mathcal{J}_{q_n}(u_n)$. Similar to the proof of Corollary 4.2 we may assume that $u_n \rightarrow u \neq 0$.

As before, if $q \in (0, q_0^*]$ the lemma holds due to (4.1).

If $q \in (q_0^*, q_0^* + \varepsilon)$ observe from Corollary 4.5 that $\|u\| > \rho$. We claim that $\widehat{\mathcal{J}}_{q_n} \rightarrow \widehat{\mathcal{J}}_q$ as $n \rightarrow \infty$. Indeed, once $(0, q_0^* + \varepsilon) \ni q \mapsto \widehat{\mathcal{J}}_q$ is increasing, we can assume that $\widehat{\mathcal{J}}_{q_n} < \widehat{\mathcal{J}}_q$ for each n and $\widehat{\mathcal{J}}_{q_n} \rightarrow \mathcal{J}_q(u) \leq \widehat{\mathcal{J}}_q$ as $n \rightarrow \infty$, which implies that $\mathcal{J}_q(u) = \widehat{\mathcal{J}}_q$.

Now suppose that $q_n \downarrow q \in (0, q_0^* + \varepsilon)$. Once $(0, q_0^* + \varepsilon) \ni q \mapsto \widehat{\mathcal{J}}_q$ is increasing, we can assume that $\widehat{\mathcal{J}}_q < \widehat{\mathcal{J}}_{q_n}$ and $\widehat{\mathcal{J}}_q \leq \lim_{n \rightarrow \infty} \widehat{\mathcal{J}}_{q_n}$. Choose u_q such that $\widehat{\mathcal{J}}_q = \mathcal{J}_q(u_q)$ ($\|u\| > \rho$ in case $q \in (q_0^*, q_0^* + \varepsilon)$) and observe that $\widehat{\mathcal{J}}_q \leq \lim_{n \rightarrow \infty} \widehat{\mathcal{J}}_{q_n} \leq \lim_{n \rightarrow \infty} \mathcal{J}_{q_n}(u_q) = \widehat{\mathcal{J}}_q$. \square

Now we turn our attention to the second solution.

Let $q \in (0, q_0^*)$. As a consequence of the Corollary 2.9 we have that

$$\Gamma_q = \{\gamma \in C([0, 1], H_r^1(\mathbb{R}^3)) : \gamma(0) = 0, \mathcal{J}_q(\gamma(1)) < 0\}$$

is non-empty. Define the Mountain Pass level

$$c_q := \inf_{\gamma \in \Gamma_q} \max_{t \in [0,1]} \mathcal{J}_q(\gamma(t)) > 0.$$

By Proposition 2.3 and Proposition 3.1 we deduce the following

Proposition 4.7. *For each $q \in (0, q_0^*)$ there exists $w_q \in H_r^1(\mathbb{R}^3) \setminus \{0\}$ such that $\mathcal{J}_q(w_q) = c_q$ and $\mathcal{J}'_q(w_q) = 0$. In particular $\mathcal{J}_q(w_q) > \mathcal{J}_q(u_q) \in (-\infty, 0)$.*

Let now $q \in (q_0^*, q_0^* + \varepsilon)$, where $\varepsilon > 0$ is the one fixed in correspondence of δ in (4.2). Let $u_q \in \mathcal{N}_q^+$ (by Proposition 4.5) such that $\mathcal{J}_q(u_q) = \widehat{\mathcal{J}}_q$. Define

$$d_q = \inf_{\gamma \in \Gamma_q} \max_{t \in [0,1]} \mathcal{J}(\gamma(t)),$$

where $\Gamma_q = \{\gamma \in C([0, 1], H_r^1(\mathbb{R}^3)) : \gamma(0) = 0, \gamma(1) = u_q\}$.

Proposition 4.8. *For each $q \in (q_0^*, q_0^* + \varepsilon)$ there exists $w_q \in H_r^1(\mathbb{R}^3) \setminus \{0\}$ such that $\mathcal{J}_q(w_q) = d_q$ and $\mathcal{J}'_q(w_q) = 0$. In particular $\mathcal{J}_q(w_q) > \mathcal{J}_q(u_q)$.*

Proof. Indeed, we combine Proposition 2.3 with the inequality $M > \delta \geq \widehat{\mathcal{J}}_q = \mathcal{J}_q(u_q)$ to obtain a Mountain Pass Geometry for the functional \mathcal{J}_q . The proof follows from (ii) in Proposition 3.1. \square

Now we conclude the proof of item 2. of Theorem 1.3.

Let ε be given as in the Proposition 4.4. The existence of the minimum u_q follows from Proposition 4.1, Corollary 4.2, Corollary 4.5. The existence of a Mountain Pass critical point w_q satisfying $\mathcal{J}_q(w_q) > \max\{0, \mathcal{J}_q(u_q)\}$ follows by Proposition 4.7 and Proposition 4.8. That u_q and w_q are actually critical points of \mathcal{J}_q follows by Proposition 2.5.

REFERENCES

- [1] V. Benci and D. Fortunato, *A eigenvalue problem for the Schrödinger-Maxwell system*, Topological Meth. Nonlinear Anal **55** (1977), 149–162. [3](#)
- [2] P. d’Avenia and G. Siciliano, *Nonlinear Schrödinger equation in the Bopp-Podolsky electrodynamics: solutions in the electrostatic case*, arXiv:1802.03380. [1, 2, 3, 4, 5, 10](#)
- [3] Y. Ilyasov *On extreme values of Nehari manifold method via nonlinear Rayleigh’s quotient*, Topological Meth. Nonlinear Anal. **49** (2017) 683–714. [9](#)
- [4] Y. Ilyasov and K. Silva, *On branches of positive solutions for p -Laplacian problems at the extreme value of the Nehari manifold method*, Proc. Amer. Math. Soc. **146** (2018), 2925–2935. [5](#)
- [5] D. Ruiz, *The Schrödinger-Poisson equation under the effect of a nonlinear local term*, J. Funct. Anal. **237** (2006), 655–674. [1, 3, 4, 6, 10, 14](#)
- [6] D. Ruiz and G. Siciliano, *A Note on The Schrödinger-Poisson-Slater Equation on Bounded Domains*, Adv. Nonlinear Stud. **8** (2008), 179–190. [3](#)

- [7] K. Silva and A. Macedo, *Local minimizers over the Nehari manifold for a class of concave-convex problems with sign changing nonlinearity*, J.Diff. Equations, **265** 5, (2018), 1894–1921. [5](#)
- [8] W. A. Strauss, *Existence of solitary waves in higher dimensions*, Comm. Math. Phys. **55** (1977), 149–162. [12](#)

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