

1 **A discriminatory diagram of massive vs stratified deposits based on the sedimentation and**  
2 **bedload transportation rates. Experimental investigation and application to pyroclastic density**  
3 **currents.**

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10

11 **Abstract**

12 Dense gas-particle jets similar to collapsing eruption columns were generated by large-scale  
13 experiments. The column collapse resulted in a ground-hugging current forming layers with bedding  
14 similar to natural pyroclastic density current deposits. At the impact of the column on the ground, a  
15 thick, massive bed was formed due to a high sedimentation rate. The flow expanded laterally into a  
16 fully turbulent, dilute gas-particle current that deposited thin layers showing ripple-like structures.  
17 The ripples were interpreted as tractional structures formed by the current ability to move bedload  
18 material settled from suspension.

19 Experiments fed with fine ash formed deposits without tractional structures, because fine particles  
20 require a large tangential stress to be sheared.

21 Experimental outcomes show that massive beds are formed where the sedimentation rate per unit  
22 width  $S_{rw}$  exceeds the bedload transportation rate  $Q_b$  by 100 times. A lower ratio generates traction at  
23 the base of the flow and formation of shear structures that increase in wavelength and height with a  
24 decreasing flux.

25 We present a diagram that provides a useful addition for facies analysis of pyroclastic density  
26 currents, provided that deposits representing sustained sedimentation can be identified in the field. In  
27 the diagram a decrease in the  $S_{rw}/Q_b$  ratio corresponds to an increase of bedform scale. Application  
28 of the diagram for hazard assessment purposes allow the reconstruction of the mass eruption rate of  
29 the Agnano-Monte Spina eruption at Campi Flegrei and of the Bingham rheology of the massive  
30 underflow of the Mercato pyroclastic density current at Vesuvius.

31

## 32 1. Introduction

### 33 1.1 General context of pyroclastic density currents

34 Pyroclastic density currents often form from the collapse of a dense gas-particle jet, called eruption  
 35 column. They are fast moving highly hazardous flows associated with medium to large magnitude  
 36 explosive volcanic eruptions. The likelihood of their formation and the nature of formative flows is  
 37 strongly controlled by the flow rate at the volcanic vent,

$$38 \rho_{mix} V_0 A \quad (1)$$

39 which is called the Mass Eruption Rate, *MER*, where

$$40 \rho_{mix} = \rho_s C + \rho_g (1-C) \quad (2)$$

41 is the gas-particle mixture density, with  $\rho_s$  particle density,  $\rho_g$  gas density and  $C$  particle volumetric  
 42 concentration.  $V_0$  is vent exit velocity and  $A$  is the vent cross-sectional area of radius  $r$  (for symbols  
 43 refer to the Symbol List). Typical values of *MER* for the formation of substantial pyroclastic density  
 44 currents are in the range  $10^7$ -  $10^9$   $\text{kg s}^{-1}$  (Cioni et al., 2000) with  $V_0$  in the range of 100 m/s,  $r$  in the  
 45 range of 100 m,  $C$  a few percent and a specific mass flow rate per unit area at conduit exit in the range  
 46  $10^3$ - $10^4$   $\text{kg m}^{-2} \text{s}^{-1}$ . It has been demonstrated (Dellino et al., 2014) that the condition for the formation  
 47 of a completely collapsing column is met when the densimetric Froude number at the vent,  $Fr'_0$ , is  
 48 between 1 and 3, where

$$49 Fr'_0 = \frac{V_0}{\sqrt{g' r_0}} \quad (3)$$

50 with  $g' = g \left( \frac{\rho_{mix} - \rho_{atm}}{\rho_{atm}} \right)$  and  $\rho_{atm}$  atmosphere density. When a complete collapse happens (i.e. no  
 51 contemporaneous ash convective dispersal in the surrounding atmosphere occurs) the column is  
 52 characterized by a very low, or null entrainment of atmosphere air (Dellino et al. 2014) and the  
 53 eruptive mixture, along the ascending-descending trajectory, conserves its bulk density and respects  
 54 Bernoulli's equation of the conservation of total pressure, with maximum height  $H = V_0^2 / 2g$ . It means  
 55 that at collapse, the flow hits the ground at the same speed and particle volumetric concentration of  
 56 conduit exit. The flow, after the impact, moves down the volcano flank due to the higher density of

57 the gas-particle mixture with respect to the surrounding atmosphere. The resulting deposits show a  
 58 wide range of structures, including both massive beds and stratified layers with tractional structures  
 59 like dunes and ripples (e.g., Pritchard and Gladstone, 2009). Since the last half of the last century a  
 60 debate goes on whether there is a rigid demarcation between dense and dilute currents (e.g., Dade,  
 61 Fisher, 1979; 2003; Sparks, R.S.J., 1976; Sulpizio et al., 2014) forming, separately, massive or  
 62 stratified deposits, or there is the possibility that a flow could form both deposit types as a function  
 63 of changes in the regime of sedimentation (e.g., Branney and Kokelaar, 2002; Sumner et al., 2008;  
 64 Sulpizio et al., 2007, 2014). The subject is relevant for hazard assessment, because dense and dilute  
 65 currents have a different impact potential. The impact of pyroclastic density currents on  
 66 infrastructures and life depends on flow characteristics such as particle volumetric concentration,  
 67 grain-size distribution and flow speed, in addition to temperature (Valentine, 1998; Baxter et al.,  
 68 2005; Dellino et al. 2008; Manville et al., 2009; Caricchi et al., 2014; Mele et al. 2015; Esposti Ongaro  
 69 et al., 2016). Metric thick, dense, massive flows are valley ponded, tend to move slowly on steep  
 70 slopes, maintain high temperature and are destructive at the valley apron at the base of the volcanic  
 71 edifice. Hundreds of meters thick, turbulent currents, move on the plain areas surrounding the  
 72 volcano. They are able to surmount substantial topographies and are lethal also in distal reaches,  
 73 because of their content of harmful fine ash taken into suspension by fluid turbulence (Druitt and  
 74 Kokelaar, 2002).

75

## 76 ***1.2 Physical properties and sedimentation regimes of turbulent suspension currents***

77 Sediment-laden flows governed by fluid turbulence, i.e. fluvial, aeolian, turbiditic and pyroclastic  
 78 density currents, can transport solid particles in three main modes: suspension load, bedload and wash  
 79 load. Suspension refers to particles that are sustained by turbulence and settle when their terminal  
 80 velocity  $w_t$  is higher than the current shear velocity,  $u_*$ .

$$81 \quad w_t = \sqrt{\frac{4Dg(\rho_s - \rho_{mix})}{3C_d\rho_{mix}}} \quad (4)$$

82 where  $D$  is particle size,  $g$  is gravitational acceleration,  $C_d$  is particle drag coefficient.

83 Bedload includes particles moved on the ground by shear stress from the overlying current. Wash  
84 load is made of particles so fine that are intimately coupled to fluid turbulence, are carried by the  
85 current and settle during the waning stage of flowage (e.g., Gladstone et al., 1998; Duller et al., 2008;  
86 Chanson, 2009; Meiburg and Kneller, 2010; Kok et al., 2012; Dellino et al., 2019).

87 In the time-space evolution of a current, when the settling velocity approaches flow shear velocity,  
88 particles, after being transported in suspension, eventually settle to the ground and form the bedload  
89 material (Branney and Kokelaar, 2002; Sulpizio and Dellino, 2008; Roche et al., 2013; Sulpizio et  
90 al., 2014; Lube et al., 2015; Dufek, 2016).

91 The bedload transportation rate is expressed, in volumetric form (Wilcock and Crowe, 2003), in terms  
92 of

$$93 \quad W_i^* = \frac{(\rho_s/\rho_{fmix} - 1)gq_{bi}}{\phi_i u_*^3} \quad (5)$$

94 where  $W_i^*$  is a non-dimensional factor, which is a function of the shear stress at the base of the  
95 current,  $q_{bi}$  is the volumetric bedload transport rate of the  $i_{th}$  size fraction per unit width of the flow,  
96  $\phi_i$  the proportion of the  $i_{th}$  size fraction of the bedload and  $u_*$  is the flow shear velocity.  $Q_b = \sum_i^n q_{bi}$   
97 is the bedload transportation rate of all size fractions.

98 Sedimentation can be sustained as long as the flow is fed with particles (Brown and Branney, 2013;  
99 Rowley et al., 2014; Roche, 2015). The sedimentation rate  $S_r = \sum_{i=1}^n S_{ri}$  expresses the rate of  
100 particles sedimented over a unit surface of the ground per unit time, where

$$101 \quad S_{ri} = w_{ti} \rho_{si} \gamma_i \quad (6)$$

102 with  $\gamma_i$  the proportion of particles of the  $i_{th}$  size fraction (out of the total concentration) that are  
103 sedimented from the current (Dellino et al. 2019).

$$104 \quad A_r = \frac{S_r}{\rho_{dep}} \quad (7)$$

105 is the aggradation rate of deposit thickness  $d$ , with a bulk deposit density  $\rho_{dep}$ .

106 The time of deposit aggradation

$$107 \quad t = \frac{d}{A_r} \quad (8)$$

108 represents the time spent by the current to form the whole deposit.

109 In some cases, it has been demonstrated that thick, massive, deposits are formed directly from  
110 suspension, because of a high sedimentation rate, also called suspended-load fallout rate, which  
111 inhibits traction at the bedload (Lowe, 1982 and 1988; Fisher, 1990; Druitt, 1992; Kneller and  
112 Branney, 1995; Woods et al., 2002; Postma et al., 2009).

113 The hypothesis of this paper is that the limit between the formation of massive and stratified deposits  
114 is governed by a threshold ratio between the sedimentation and the bedload transportation rates, and  
115 is also a function of the amount of cohesive material and particle volumetric concentration (Sumner  
116 et al, 2008; Baas et al., 2011). In order to verify this hypothesis, the sedimentation and bedload  
117 transportation rates need to be measured in actual currents. Unfortunately, it is difficult to monitor  
118 the formation of deposit structures directly during an eruption. In fact, pyroclastic density currents  
119 are hard to predict and/or to instrument, and the measurement of their internal characteristics, as of  
120 today, is impossible from remote sensing. In the present paper, the sedimentation and bedload  
121 transportation rates are measured in currents formed by means of large-scale experiments that were  
122 designed as to simulate pyroclastic density currents formed by the collapse of an eruption column.  
123 As an outcome of experiments, a diagram for discriminating between massive and stratified deposits  
124 is proposed. In the diagram, the scale of bedforms of stratified deposits increases where the ratio  
125 between sedimentation and bedload sedimentation rates decreases. The diagram is applied to deposits  
126 of Vesuvius and Campi Flegrei volcanoes in Southern Italy, as to verify its utility in the reconstruction  
127 of the physical characteristics of past eruptions of the two hazardous Italian active volcanoes.

128

## 129 **2. Experimental setup, materials and methods**

130 Experiments were carried out with the aim of generating density currents forming measurable  
131 massive and stratified deposits, starting from a dense gas-particle jet similar to a collapsing eruption

132 column. The experiments consider the case of a completely collapsing eruption column and do not  
133 include all the possible variations as, for example, the occurrence of contemporaneous fallout, re-  
134 fragmentation of particles after conduit exit, or the effect of valleys and/or topographic barriers.  
135 On previous papers (Dellino et al., 2011; Dellino et al., 2014) experiments carried out both at ambient  
136 temperature and up to 300° C demonstrated that, in the case of completely collapsing columns (in  
137 contrast with convective plumes), temperature doesn't play a major role neither in determining the  
138 jet dynamics nor in the sedimentation from the base of the ground-moving density current. In the  
139 present paper, experiments carried out at ambient temperatures are considered.

140 The experiments were performed with the apparatus described in detail by Dellino et al., 2007;  
141 2010a; 2010b; and 2014. The design (Fig. 1) consists of pressurized-gas bottles (Fig. 1a) connected,  
142 via valves and manometers, to long hoses that are charged with a known amount of gas (nitrogen) at  
143 known pressure. In this way, a known quantity of potential energy is stored, and is released by gas  
144 expansion upon opening the valves. The hoses are connected to a hub of fast solenoid valves (Fig.  
145 1b) that are attached, via short hoses and blow nozzles, to the base plate of a 3.2 m long, vertical  
146 cylindrical conduit with 0.6 m internal diameter (Fig. 1c). The blow nozzles allowed the injection of  
147 gas into the conduit and its coupling to the particulate material resting on the base plate.

148 The particulate material used in the experiments comes from pyroclastic deposits of the Mercato  
149 eruption at Vesuvius and of the Agnano Monte Spina eruption at Campi Flegrei in Southern Italy.  
150 The grain size of two compositions, representing the coarse (from Vesuvius) and fine (from Campi  
151 Flegrei) end members, are plot on Fig. 2 as to show the range used. The coarser composition ranges  
152 from lapilli to fine ash (a few micrometres to a few centimetres, Fig. 2a), with a median size  $D_{sm} =$   
153 0.75 mm. It is made of dense lithic, vesicular glass and crystal components. The fine composition  
154 (Fig. 2b) has median size  $D_{sm} = 0.066$  mm and is made almost exclusively of glass fragments.

155 In this paper, five runs are considered. For each run data were collected at multiple locations along  
156 flow runout. The particle load (up to 350 kg), gas pressure and grain size were varied among runs as  
157 to obtain density currents of different speed, particle concentration and granulometry (see Tab.1).

158 Blue plastic sheets were mounted on the ground, over an ample area around the conduit (about 2000  
159 m<sup>2</sup>), so to allow easy sampling of the deposits left after the experiments. The deposits were  
160 photographed and sampled over the whole dispersal area. The mass of sediment per unit area was  
161 calculated by sampling the sediment over a rectangular area, at regular distances (Fig. 3), and  
162 weighting it. The deposits were analysed in the laboratory for grain-size characterization and  
163 componentry analysis. Grain size analysis was carried on by dry sieving for the fractions > 0.125  
164 mm and by means of a Coulter Multisizer IV for the fractions  $\leq 0.125$  mm. Data were combined by  
165 software as histograms at half-phi interval.

166 The experiments start with the computer-controlled opening of the solenoid valves, allowing a fast  
167 coupling of the pressurized gas with the pyroclastic material filling the conduit that, on mixing with  
168 the expanding gas, accelerates along the conduit (Dioguardi et al., 2013). The two-phase mixture is  
169 finally expelled out of the conduit in the form of a dense gas-particle jet. The jet, with conduit exit  
170 velocities up to 20 m/s, reaches a maximum height (typically more than 10 m over conduit exit, Fig.  
171 4a), which follows the prescription of Bernoulli's equation. Along the ascending-descending  
172 trajectory no significant air entrainment occurs (Dellino et al., 2014). The density of the gas-particle  
173 mixture, hence particle volumetric concentration and mass flow rate, do not change with respect to  
174 the conduit exit.

175 The experiments were monitored with a network of multiple high definition videocameras (recording  
176 at 50 frames per second) and pressure sensors, which allowed measurements of the jet speed and  
177 geometry (including heights, lengths and volume) and speed and dynamic pressure of the density  
178 currents formed after the jet impact on the ground.

179 A frame of the close-up view of the passage of an experimental current is shown on Fig. 5, with insets  
180 of average flow velocity, shear velocity, and dynamic pressure values. In the figure, the current  
181 thickness is also shown and the position of a pressure sensor is indicated.

182 The shear velocity of the density currents was calculated following the method of Dellino et al., 2008,  
183 by means of the law of the wall of a turbulent shear flow moving on a rough substrate:  $u_{avg}(y) = u_*$   
184  $\left(2.5 \ln \frac{H_T}{k_s} + 6.5\right)$  (9)

185 where  $k_s = 1$  cm is the substrate roughness as measured on the ground,  $H_T$  is current thickness and  
186  $u_{avg}(y)$  is the height-averaged current velocity measured at flow front by video analysis.

187 The mass flow rate at the impact of the jet on the ground was calculated by means of Equation (1),  
188 by assuming that it was conserved along the ascending descending jet trajectory (Dellino et al., 2014).

189 The particle volumetric concentration of the current was calculated by equation (2), by means of the  
190 known gas and particle densities. The current density  $\rho_{mix}$  was calculated by dividing the mass flow  
191 rate by the average velocity. The mass flow rate in the current was calculated, at each location, by  
192 subtracting the mass of sediment deposited at previous locations from the mass flow rate at the impact,  
193 and using as the cross-sectional area of the flow, the value measured at each location by multiplying  
194  $H_T$  for  $2\pi L$ , where  $L$  is the distance, from the impact, reached by the flow front.

195 The particle volumetric concentration in the jet,  $C_1$  ranged between 0.019 and 0.063, as it was  
196 measured by the ratio of total particle load and jet volume (see Tab. 1). At the passage from the  
197 sedimentation of the descending jet to the sedimentation of the density current, the particle volumetric  
198 concentration in the gas-particle mixture,  $C_2$  decreases (see Tab. 1). It decreases also along the current  
199 runout as a consequence both of sedimentation and of increase of the cross-sectional area of the flow  
200 at increasing distance.

201

### 202 **3. Results**

#### 203 ***3.1 Generation, evolution and deposition of experimentally-generated PDCs***

204 The collapsing gas-particle jet, on hitting the ground (Fig. 4b), resembles the collapse of a natural  
205 eruption column leading to the formation of pyroclastic density currents. Upon impact, the jet is more  
206 than one meter wide and leaves on the ground a tens of centimetre thick deposit (Fig. 4c). The collapse

207 occurs on the right side because of the slight inclination of the conduit. This condition allowed the  
208 formation of a column impacting on the ground and not falling back in to the conduit. The thickness  
209 of the deposit decreases from the centre of the impact to the margin, with a shape approximating a  
210 cone. At the centre of the impact the deposit is massive, and does not show any visible sedimentary  
211 structure related to shear movement (Fig. 4d). The deposit at the margin of the impact, while  
212 decreasing in thickness, starts to show some undulations (Fig. 4c), similar to ripples, likely to be  
213 attributed to an incipient shear movement. In the case of fine-grained experimental runs, no tractive  
214 structures are formed, neither at the impact (Fig. 4e), nor at distal locations (Fig. 4f). Similar  
215 characteristics are reported in natural fine-ash deposits, which show a monotonous, structureless  
216 facies, with no grain-size variation along runout (e.g., Branney and Kokelaar, 1997; Darteville et al.,  
217 2002; Kokelaar et al., 2007).

218 After the impact of the fountain on the ground, the normal stress of the fluid is transformed into a  
219 tangential stress, leading to a shear flow that evolves laterally into a few meters thick, fully turbulent,  
220 gas-particle current, moving at several m/s (Fig. 6a). It spreads for more than 20 m over the zone of  
221 the impact and leaves on the ground a thin bed (a few cm thick) formed by the sustained alimentation  
222 of the sediment flux from the collapse. When the run is fed with the pyroclastic material of the coarser  
223 composition, the thin layer shows ripple-like sedimentary structures (Fig. 6b), which demonstrate that  
224 the bedload was moved at a rate sufficient to form tractional bedforms. The structures represent the  
225 lateral evolution of the undulations visible at the lateral margin of the impact.

226 A close-up view of two frames at the current front (Fig. 6c and d) shows the progressive shearing  
227 effect that the tangential stress exerts on the particles of the bedload, which favours the formation of  
228 ripples.

229 The ripple structures are characterized by an asymmetrical distribution of particles: the finer sizes  
230 occur at the foreset and the coarser ones at the backset, implying that the transportation process, in  
231 the shear current, is selective in terms of grain size. This suggests that the finer particles rested in

232 the shadow zone within the ripple troughs. These features are similar to those shown by natural  
233 pyroclastic density currents (Sulpizio et al., 2014).

234 Some variants of the ripple structures were observed in the experimental deposits. Centimetre-thick  
235 structures that show an average wavelength about 30cm (Fig. 6b) are found in locations a few metres  
236 outside the impact zone. They are associated to fast currents ( $u_{avg}=13$  m/s;  $u_*=0.73$  m/s,  $0.0035 C$  in  
237 the example of Fig. 6b) and are made up of coarse ash and lapilli (median size  $D_{sm} = 0.79$  mm in the  
238 example of Fig 7b). Other cm-thick narrowly spaced ripple-like structures, showing an average  
239 wavelength of about 8 cm, are also found a few metres outside the impact zone. The ripples are made  
240 of coarse ash and lapilli, (median size  $D_{sm} = 1.69$  mm in the example of Fig. 6e) and are associated  
241 to a lower current speed ( $u_{avg} = 6.3$  m/s,  $u_*=0.38$  m/s in the example of Fig. 6e). An average  
242 wavelength of 20 cm is shown by mm-thick ripples found in distal locations, more than 10 m outside  
243 the impact zone. They are made of medium ash (median size  $D_{sm} = 0.26$  mm in the example of Fig.  
244 6f), and are formed by slow currents ( $u_{avg} = 3.9$  m/s,  $u_*=0.19$  m/s, in the example of Fig. 6f). Such  
245 data suggest that the spacing and thickness of tractional features are somehow linked to the flow  
246 regime and grain size of the particulate material, as it is postulated for sedimentary currents (Lowe,  
247 1988; Postma et al., 2009; Cartigny et al., 2014). It is, however, quite difficult to find a consistent  
248 relationship between speed, spacing and grain size in the distal deposits, especially in the case of gas-  
249 particle currents (Middleton, 1984).

250 In further locations that were reached by the flow at the time the fountain had stopped feeding the  
251 current (Fig. 7a), only fine-ash material, lacking sedimentary structures, was deposited. This material  
252 is spread well over the blue plastic sheets and formed a submillimetric veil of ash (Fig. 7b) that  
253 represents the waning-flow washload.

254 A general facies model of the experimental deposits, associated to the flow evolution along runout,  
255 is shown on Fig. 8. The proximal facies is massive and lays at the center of the impact. The proximal-  
256 transitional facies shows some incipient shear structures at the surface and is formed around the lateral  
257 zone of impact, where the flow rate developed basal shear stress. The medial facies is associated to

258 the fully turbulent, fast, shear current and is characterized by coarse-ash ripple-like bedforms. The  
259 distal facies is associated to the down-current evolution of the flow, where velocity decreases and the  
260 particle load is less concentrated and finer grained, resulting in ripple-like structures of medium ash.  
261 The ultra-distal zone is associated to the waning stage of the flow where the wash load is slowly  
262 settled and forms a thin, structureless, fine-ash bed.

263

### 264 ***3.2 Scaling between experimental and natural currents***

265 The obvious difference between the experiments and natural pyroclastic density currents is the much  
266 larger thickness of natural layers, which results from the vertical aggradation of multiple laminae,  
267 formed during the time-integrated passage of thick (up to hundreds of meters), expanded currents,  
268 which in nature last much longer (several minutes up to hours) than the experiments, the latter lasting  
269 up to 30 s. In the natural case, the passage of a long-lasting pyroclastic density current is likely to  
270 generate a vertical facies association (Mele et al., 2015), also depending on the evolution of the  
271 eruption feeding the collapsing column, which is not considered here.

272 Notwithstanding these differences, the values of the fluid-flow parameters of the experiments are in  
273 the same order of magnitude of natural pyroclastic density currents, suggesting that a scaling between  
274 experiments and nature can be proposed. For example: i) the specific mass flow rate at conduit exit

275 was up to 4000 kg/m<sup>2</sup>s; ii) the current velocity was up to 13 m/s; iii) the Reynolds number  $Re_* = \frac{\rho_f u H_T}{\mu}$

276 , with  $\mu$  fluid viscosity, was always in excess of 10<sup>6</sup>, meaning the flow was fully turbulent.

277 Furthermore, shear velocity (up to 0.73 m/s) and dynamic pressure (up to about 500 Pa) are almost

278 comparable to those calculated for some natural pyroclastic density currents (Mele et al., 2011 and

279 2015). The experimental deposits share some features with the natural counterparts, for example: i)

280 the deposit at the centre of the impact looks similar to the layer deposited by massive pyroclastic

281 flows (Bursik and Woods, 1996; Branney and Kokelaar, 2002); ii) dilute pyroclastic density current

282 deposits have bedforms that are visually similar to those found in the experiments (Sulpizio et al.,

283 2014); iii) the facies model of the experimental deposits of Fig. 8 is similar to what is observed in

284 some ignimbrites (e.g., Vazquez and Ort, 2006; Brown and Branney, 2013; Wright et al., 2016); iv)  
285 the contemporaneous occurrence of massive and stratified deposits related to both concentrated and  
286 dilute currents has been already noted in volcanology (Druitt et al., 2002; Gernon et al., 2013; Breard  
287 et al., 2016; Breard and Lube, 2017).

#### 288 289 **4. Modelling of the sedimentation and bedload transportation rates**

290 In order to identify the threshold ratio that marks the transition between massive and stratified  
291 deposits, the sedimentation and bedload transportation rates were calculated at multiple locations  
292 along runout. They include the impact point, as for recording the formation of the thick massive  
293 layer, and further locations, where the turbulent suspension current lead to the sedimentation of the  
294 stratified beds. A total of 13 locations are considered in this paper, 5 out of them represent as many  
295 thick massive layers formed at the impact, and 8 represent beds formed by the turbulent suspension  
296 currents. The experimental deposits were sampled as for obtaining the particle parameters to be  
297 used in the calculation of  $S_r$  and  $Q_b$ . They include grain size, density and componentry.

#### 298 299 **4.1 Methods and measurements used in the calculation of $S_{rw}$ and $Q_b$**

300 The sedimentation rate  $S_r$  (see Tab. 1) was measured as the mass sedimented per unit area (the latter  
301 being obtained by an area-integrated sample of sediment), over the time of sedimentation, which was  
302 measured by using the time length of the video-frames between the arrival of the current and the end  
303 of the jet feeding the current.

304 For the thick, massive deposits, the sedimentation rate  $S_r$  ranged between 39 and 220  $\text{kgm}^{-2}\text{s}^{-1}$ ; while  
305 it ranged from 0.19 to 1.17  $\text{kgm}^{-2}\text{s}^{-1}$  for the deposits with ripple-like tractional structures. It was about  
306 0.022  $\text{kgm}^{-2}\text{s}^{-1}$  for the thin structureless fine-ash beds.

307 The bedload transportation rate was calculated by modifying the model of Wilcock and Crowe 2003,  
308 which is based on flume experiments. In the model, a correlation between  $W_i^*$  (Equation 5) and the  
309 shear stress  $\tau$  at the base of the current is introduced

$$310 \quad W_i^* = \begin{cases} 0.002\xi^{7.5} & \text{for } \xi < 1.35 \\ 14\left(1 - \frac{0.894}{\xi^{0.5}}\right) & \text{for } \xi \geq 1.35 \end{cases} \quad (10)$$

311 with

$$312 \quad \xi = \tau/\tau_{ri} \quad (11)$$

313 where the shear stress at the base of the current  $\tau$ , was calculated by the relation  $\tau = \rho_{mix}u^*$ .

314  $\tau_{ri}$  is the minimum shear stress needed to move by traction the  $i_{th}$  size fraction;  $\tau_{ri}/\tau_{rsm} = (D_i/D_{sm})^b$ ,

315 where  $\tau_{rsm}$  is the minimum shear stress needed to move the bedload mixture by traction, with  $b =$

316  $0.67/1 + \exp(1.5 - D_i/D_{sm})$ ,  $D_i$  the diameter of the  $i_{th}$  size fraction and  $D_{sm}$  the diameter of the

317 median size of the grain-size distribution of the bedload material;  $\tau_{rsm}^* =$

318  $\tau_{rsm} / \left( \frac{\rho_s}{\rho_f - 1} g \rho_f D_{sm} \right)$  with  $\tau_{rsm}^*$  the critical Shields number normalized to the mixture of the

319 bedload material.  $\tau_{rsm}^* = 0.75\theta_{tmedian}$  where  $\theta_{tmedian}$  is the critical Shields number of the median

320 size. The critical Shields number  $\theta_t = \frac{\rho_f u_*^2}{Dg(\rho_s - \rho_f)}$  is the ratio of surface to volume forces and represents

321 the threshold of initiation of motion of particles resting on the substrate. It is a function of the

322 Reynolds' number of shear  $Re_* = \frac{\rho_f u_* D_{sm}}{\mu}$ . The critical values, for various fluids, including liquids

323 and gas, are found in the review of Miller et al. (1977).

324 In the original Wilcock and Crowe 2003 model the relation  $\tau_{rsm}^* = 0.021 + 0.015\exp[-20F_s]$

325 was used, with  $F_s$  representing the percent of sand in the sediment. In this paper, as resulting from

326 our tests on volcanic material, a value of 0.75 of the critical Shields number of the median size

327  $\theta_{tmedian}$  was used, as representing the critical Shields number normalized to the mixture of the

328 bedload material  $\tau_{rsm}^*$ . The value of the Shields number of the median size was derived from the

329 Miller et al., (1977) curve for gas. It ranged from 0.010 for the fastest, coarse grained, currents to

330 0.019 for the slowest fine-grained current that formed thin structureless fine-ash beds.

331 The procedure for the calculation of  $Q_b$  starts with the analysis of the multicomponent grain-size  
 332 distribution of the deposit. For each size fraction and component (glass, lithics and crystals) the  
 333 critical shear stress for the initiation of motion is calculated and the ratio between the shear stress at  
 334 the base of the current and the critical shear stress is calculated for each size class. Equations (10)  
 335 and (5) are used to calculate the bedload transportation rate of each size class,  $q_{bi}$ . The total bedload  
 336 transportation rate per unit width  $Q_b$  is obtained by summation of the bedload transportation rate of all  
 337 size classes.

338 With the aim of illustrating step by step the calculation of  $Q_b$ , in the supporting files A, B and C,  
 339 examples of calculation of a massive layer at the impact; a layer showing ripple-like tractional  
 340 structures and a thin structureless fine-ash bed are reported, respectively. Grain size, density and  
 341 componentry data are included.

342 The bedload transportation rate data are reported in Tab. 1. It ranged from  $9.3 \times 10^{-5} \text{ m}^2\text{s}^{-1}$  to  $2 \times 10^{-3}$   
 343  $\text{m}^2\text{s}^{-1}$  for the massive layers, while it ranged from  $7.4 \times 10^{-6} \text{ m}^2\text{s}^{-1}$  to  $1.15 \times 10^{-3} \text{ m}^2\text{s}^{-1}$  for the layers  
 344 with ripple-like tractional structures. It is  $1.77 \times 10^{-7} \text{ m}^2\text{s}^{-1}$  for the thin structureless fine-ash beds.

345

#### 346 ***4.2 The threshold ratio of sedimentation and bedload transportation rates and the discriminatory*** 347 ***diagram between massive and stratified deposits***

348 In order to compare sedimentation and bedload transportation rates easily, they need to be  
 349 expressed in the same physical dimension ( $\text{m}^2\text{s}^{-1}$ ). For this purpose  $S_r$  has been transformed into the  
 350 volumetric sedimentation rate per unit width,  $S_{rw}$  by the formula of transformation

$$351 \quad S_{rw} = A_r \text{Ref}_a \text{Ref}_w \quad (12)$$

352 The transformation of Equation (12) starts by first calculating the growth rate of sediment thickness  
 353  $A_r$  ( $\text{ms}^{-1}$ ) by means of equation (7). A square of one meter width is considered as the reference area  
 354 on the ground,  $\text{Ref}_w$ , which is the same area on which  $Q_b$  is calculated. By means of the  
 355 multiplication of  $A_r$  for the reference area (which is one square meter), the volumetric sedimentation

356 rate is obtained ( $\text{m}^3\text{s}^{-1}$ ). By dividing the volumetric sedimentation rate by the width of the reference  
 357 area,  $Ref_w$  (which is one metre), we obtain  $S_{rw}$  ( $\text{m}^2\text{s}^{-1}$ ), which has the same dimension of  $Q_b$ .  
 358 Upon calculation of  $S_{rw}$  and  $Q_b$ , we verified that: i) the ratio of particle volumetric concentration  $C_1$   
 359 in the impacting fountain that fed the massive layer, and ii) the particle volumetric concentration  
 360 transported in the suspension current  $C_2$  at the first location after the impact are a function of the  
 361 ratio of  $Q_b^*/S_r$ , which is defined by:

$$362 \quad \frac{C_1}{C_2} = 0.5581 \left( \frac{Q_b^*}{S_r} \right)^{-0.578} \quad (13)$$

363 where  $Q_b^* = Q_b \frac{D_{sm}}{1\text{mm}}$  is the bedload transportation rate normalised to a grain size of one millimetre.  
 364 Data of  $C_1$ ,  $C_2$ ,  $Q_b^*$ ,  $Q_b$ ,  $D_{sm}$ ,  $S_{rw}$  and  $S_r$  are found in Tab. 1. The trend of equation (13) is illustrated  
 365 by Fig. 9. Equation (13), being based on a single data set, would need further investigation for proving  
 366 its general validity. It can be useful for back calculating the concentration in the dense eruption  
 367 column, starting from the concentration in the current, once the sedimentation and bedload  
 368 transportation rates are known, which could help in the reconstruction of the *MER* of past eruptions.  
 369 Data of the sedimentation rate per unit width of the experiments are listed in Tab. 1. They range  
 370 from  $0.030 \text{ m}^2\text{s}^{-1}$  to  $0.2 \text{ m}^2\text{s}^{-1}$  for the massive layers and from  $8.2 \times 10^{-5} \text{ m}^2\text{s}^{-1}$  to  $9.75 \times 10^{-3} \text{ m}^2\text{s}^{-1}$  for the  
 371 layers with ripple-like tractional structures. It is  $1.78 \times 10^{-5} \text{ m}^2\text{s}^{-1}$  for the thin structureless fine-ash  
 372 beds.

373 Experiments suggest that massive and stratified layers form under different conditions of bedload  
 374 transport and sedimentation, and that the abundance of very fine material, and particle volumetric  
 375 concentration, have an influence in the formation of deposit structures. This is similar to what has  
 376 been already suggested by Lowe (1988) for subaqueous density flows. Our results indicate that in the  
 377 thick massive beds formed at impact, where no shear structures are visible,  $S_{rw}$  is at least 100 times  
 378 larger than  $Q_b$  (see Tab. 1). These data suggest that the threshold limit between the conditions leading  
 379 to the formation of massive vs stratified deposits can be represented, in a  $Q_b$  vs  $S_{rw}$  diagram, as the  
 380 line of  $S_{rw}=100Q_b$ . Such a diagram is shown on Fig. 10, with the massive beds formed at impact

381 plotting below the threshold limit on the right-hand side, where both the sedimentation rate per unit  
382 width and the bedload transportation rates are high. Also the thin structureless fine-ash beds formed  
383 by dilute currents plot below the threshold limit, on the left-hand side of the diagram, where both the  
384 sedimentation rate and bedload transportation rates are very low. This means that both with high and  
385 low sediment flux a value of  $S_{rw}$  lower than  $100 Q_b$  inhibits the formation of tractional features in the  
386 deposit. In the thin structureless fine-ash beds, the absence of tractional features indicates that,  
387 although the sedimentation rate is very low (see Tab. 1), the shear stress was not effective in moving  
388 the bedload material. This is due to the very high value that the critical Shields number  $\theta_t$  reaches  
389 where the shear Reynolds number is lower than 1 for gas, as it is shown in the diagram of Miller et  
390 al., (1977). In the case of the structureless fine-ash beds of the experiments, the shear Reynolds  
391 number is about 0.9 and the critical Shield number is 0.19, much higher than that of the coarse-ash  
392 ripple layers, which is 0.10. This means that very fine volcanic ash, as for other fine sedimentary  
393 materials, is difficult to be moved by traction because it is cohesive and exposes a limited surface to  
394 the shear stress (Zanke, 2001; Manning and Schoellhamer, 2013; Xu et al., 2015). This outcome can  
395 help interpreting the origin of structureless finely-grained pyroclastic deposits. In fact, structureless  
396 fine-ash beds are not necessarily and univocally the result of fallout processes from a convective  
397 eruptive column (as it is frequently assumed in volcanology), but likely the far facies left from  
398 pyroclastic density currents, also known as co-ignimbrite ash plumes (Bursik and Woods, 1996). It  
399 has implications in the assessment of the eruptive history and, hence, hazard of active volcanoes.

400  $S_{rw}$  of the layers with ripple-like structures is much lower than that of the massive deposits (and of  
401 the fine-ash structureless beds). For these layers, the ratio  $S_{rw}/Q_b$  is always lower than 100, and  
402 deposits plot above the line of  $S_{rw}=100Q_b$  of Fig. 10. In the diagram of Fig. 10, deposits with ripple-  
403 like structures of different wavelength and height are also discriminated, and a line separating the  
404 field of widely spaced vs narrowly spaced ripples is tentatively set at  $S_{rw}=4Q_b$ . This suggests that,  
405 while it is difficult to rationalize the dependence of bedform scale in terms of grain size and speed  
406 (as noted in the introduction section), it seems that the wavelength and height of sedimentary

407 structures can be better followed by the ratio between sedimentation and bedload transportation  
408 rates.

409 Even though the numerical limits of Fig. 10 would need additional tests to verify their precision and  
410 general validity, the rationale on which the diagram is based is to be noted, i.e.: “that is the balance  
411 between sedimentation and transportation rates that controls deposit structures”. This outcome, which  
412 is confirmed by experiments, encourage application to natural cases.

413

#### 414 **5. Application to natural deposits**

415 Campi Flegrei and Vesuvius in Southern Italy (Fig. 11) represent the most hazardous explosive  
416 volcanoes in Europe. The deposits investigated in this paper come from the 4.2 ka Plinian eruption  
417 of Agnano Monte-Spina at Campi Flegrei (de Vita et al., 1999; Mele et al., 2015) and from the 8.2 ka  
418 Plinian eruption of Mercato at Vesuvius (Mele et al., 2011). The sedimentation and bedload  
419 transportation rates were calculated as to both test the applicability of the discriminatory diagram of  
420 Fig. 10 to natural deposits, and to verify its utility in the reconstruction of the physical parameters of  
421 the eruptions.

422 The sedimentation rate of deposits was reconstructed by means of the model of Dellino et al., (2019)

$$423 \quad S_r = \left( \sum_i^n \rho_{s_i} w_{t_i} \left( \frac{\frac{\phi_i / \rho_{s_i}}{\sum_{i=1}^n \phi_i / \rho_{s_i}} * C_{tot}}{\left( \left( 10.065 * \frac{P_{ni}}{P_n^*} \right) + 0.1579 \right)} * 0.7 + \frac{\frac{\phi_{i+1} / \rho_{s_{i+1}}}{\sum_{i=1}^n \phi_{i+1} / \rho_{s_{i+1}}} * C_{tot}}{\left( \left( 10.065 * \frac{P_{ni}}{P_n^*} \right) + 0.1579 \right)} * 0.3 \right) \right) - 0,01 \quad (14)$$

424 that was developed starting from experiments of which the ones used in this paper represent a subset.

425  $P_{ni} = w_{ti} / ku$  is the Rouse number of the  $i_{th}$  size fraction of the solid material suspended in the current  
426 with  $k$  the Von Karman constant (0.4). The Rouse number defines the particle volumetric  
427 concentration as a function of flow height, by means of

$$428 \quad C_{tot} = C_0 \frac{1}{H_T - y_0} \int_{y_0}^{H_T} \left( \frac{H_T - y_0}{y_0} \frac{y}{H_T - y} \right)^{P_{nsusp}} dy \quad (15)$$

429 where  $C_{tot}$  is the total concentration of the current averaged over total flow thickness  $H_T$ .  $C_0$  is a value  
430 of known concentration at a specific height  $y_0$ , which is generally assumed very close to the ground  
431 (Middleton, 1984).  $P_{nsusp}$  is the Rouse number associated to the maximum concentration that can be  
432 held in suspension.  $P_n^*$  is the normalized Rouse number, which is obtained by the ratio of  $P_{n_{avg}}/P_{nsusp}$ ,  
433 with  $P_{n_{avg}} = \sum_{i=1}^n P_{ni}C_i/C_{tot}$  being the average Rouse number of the particulate material.

434 The bedload transportation rate of deposits was calculated by means of Equations (5) and (10).

435 The flow input parameters used for the calculation of the sedimentation and of the bedload  
436 transportation rates of the Mercato and Agnano Monte-Spina density currents (i.e. shear velocity,  
437 shear stress and particle volumetric concentration), were derived from Mele et al., (2011) and Mele  
438 et al. (2015), respectively. These input parameters (see Tab. 2) were calculated by means of the  
439 PYFLOW code (Dioguardi et al. 2014, 2018). PYFLOW solves for the flow properties of turbulent  
440 pyroclastic density currents based on the approach first presented by Dellino et al. (2008). The  
441 calculation utilizes the concept of hydraulic equivalence, which states that if two components (for  
442 example vesicular glass and crystals) that have different median size, density and shape, coexist in a  
443 deposit, it means they were settling from suspension with the same settling velocity (Jutzeler et al.,  
444 2015). By the assumption that sedimentation occurs when the settling velocity approaches current  
445 shear velocity (Middleton, 1984), the software equates the settling velocity of the two components  
446 and solves for the current shear velocity  $u^*$ , total concentration  $C_{tot}$  and Rouse number at maximum  
447 suspension capacity  $P_{nsusp}$ , with  $y_0$  obtained numerically (Dellino et al., 2008). Input parameters of  
448 particles: multicomponent grain-size distribution, density and drag coefficient, are obtained by  
449 laboratory investigations on sediment samples. The settling velocity,  $w_t$  was calculated by means of  
450 the model of Dioguardi et al. (2018). PYFLOW gives on output the probability density functions of  
451 each flow parameter. In the present study, the values of the average solution were utilized (see Tab.2).

452

#### 453 **4.1 The B2 Agnano Monte-Spina pyroclastic density current**

454 The deposits of Agnano Monte-Spina refer to the pyroclastic density current of the eruption phase  
455 B2. This phase followed the fallout layer of the Plinian convective column of phase B1 (de Vita et  
456 al., 1999; Mele et al., 2015). B2 phase resulted from a dense eruptive column that on impact, which  
457 occurred on an almost flat surface, generated a massive deposit. As the result of the lateral evolution  
458 of the impact, a thick, dilute pyroclastic density current was formed, which generated stratified  
459 deposits with dune-like bedforms. The fine-ash component, which represents a substantial hazard at  
460 Campi Flegrei, spread well away the caldera border. The most proximal outcrops of B2 are found at  
461 Vallone del Corvo and at Sbirri (Fig. 12a and b), two localities beneath the Agnano Monte-Spina  
462 caldera rim. At Vallone del Corvo a 8 m thick, poorly sorted, massive deposit, composed of lapilli  
463 and scattered scoriaceous bombs supported in an ash matrix crops out (Fig. 12b). Alignments of  
464 coarse clasts occur at various levels, suggesting some incipient basal shear. At Sbirri, a 4 m thick  
465 (Fig. 12c), poorly sorted, massive bed composed of lapilli and scattered bombs, supported in an ash  
466 matrix, occurs, and is correlated with the deposit at Vallone del Corvo. The deposit is overlain,  
467 conformally, by a 1.5 m thick stratified deposit of ash and lapilli showing some normal grading, likely  
468 due to variations, over time, in the feeding of the eruption column. The stratified deposit has a low  
469 angle internal bedding, which is typical of volcanic bedforms (Crowe and Fisher, 1973), and shows  
470 small-dune bedforms (in the sense of Ashley, 1990) with wavelength about 2 m and 0.15 m height  
471 (Fig. 12c). The massive bed at Vallone del Corvo can be interpreted as formed by the impact of the  
472 dense gas-particle column that fed the eruption. In contrast, the massive layer at Sbirri can be  
473 interpreted as resulting from the lateral zone of the fountain impact, where the incipient lateral  
474 movement on the ground generated, afterward, the turbulent shear current that lead to stratification.  
475 Moving away from the caldera rim, the next locality where a stratified deposit with bedding structures  
476 of the same unit B2 occurs, is at Verdolino (Fig. 12d), about 2 km from the caldera rim. The deposit  
477 is much finer and thinner compared to the stratified one at Sbirri. The layer is composed of coarse  
478 and fine ash, has a low-angle internal bedding, and includes ripple structures of coarse ash with  
479 wavelength about 0.5 m and height about 0.05 m. The deposit can be interpreted as formed from the

480 down-flow evolution of the dilute suspension current. The most distal outcrop of B2 is at Camaldoli  
 481 (Fig. 11e), which is about 4 km from the caldera rim and shows a 0.2 m thick structureless fine-ash  
 482 bed.

483 The facies architecture of B2 is quite similar with that defined by the analogue experiments on figure  
 484 8. The deposit at Vallone del Corvo looks similar to the proximal facies of the impact zone; the  
 485 deposit at Sbirri represents the proximal-transitional facies; the stratified deposit at Verdolino is  
 486 similar to the distal facies and the fine-ash structureless bed at Camaldoli represents the ultra-distal  
 487 facies of the washload settling from the waning current.

488

489 ***5.1.1 Use of the discriminatory diagram with the proximal and transitional facies of deposits and***  
 490 ***reconstruction of the mass eruption rate of B2***

491 Judging from the vicinity of the two outcrops to the caldera rim, it is reasonable to hypothesize that  
 492 the sedimentation of the massive layer at Vallone del Corvo and at Sbirri resulted from the direct  
 493 feeding of the mass flow rate of the collapsing column. The bedload transportation rate  $Q_b$  of the  
 494 massive layer was calculated by means of Equations (5) and (10). As input parameters, the shear  
 495 velocity  $u_{*1}$ , flow density  $\rho_{mix1}$ , shear stress  $\tau$  and grain-size distribution of the deposit were needed.  
 496 As for the flow density of the gas-particle mixture forming the massive layer, which is  $\rho_{mix1} = \rho_s C_1$   
 497  $+ \rho_g(1 - C_1)$ , it is reasonable to assume that it must have been higher than that of the lateral dilute  
 498 current  $\rho_{mix2} = \rho_s C_2 + \rho_g(1 - C_2)$ , which formed the stratified layer with small-dune bedforms. The  
 499 difference between  $\rho_{mix1}$  and  $\rho_{mix2}$  is due to the lower particle volumetric concentration of  $C_2$  with  
 500 respect to  $C_1$ . The latter was calculated by using as  $C_2$ , in equation (13), the value obtained for the  
 501 dilute current at Sbirri by Mele et al. (2015), which is 0.001 (see Tab. 2) and using 0.01 as the value  
 502 of  $\frac{Q_b}{S_{rw}}$ , which is the minimum value for the formation of a massive layer. The resulting value of  $C_1$  is  
 503 0.009.

504 Since both the massive and the stratified deposits at Sbirri result from the incipient impact of a dense  
 505 eruptive column, the shear stress acting on the massive layer  $\tau = \rho_{mix1}u_{*1}^2$ , and that leading to the  
 506 formation of the lateral dilute current  $\tau = \rho_{mix2}u_{*2}^2$ , have to be similar. By means of the values of  $\tau$ ,  
 507  $\rho_{mix2}$  and  $u_{*2}^2$  that were calculated by Mele et al. (2015) for the lateral dilute current that formed the  
 508 stratified layer at Sbirri (see Tab. 2), and using the value of  $C_I$  previously obtained, the value of  $u_{*1}$   
 509 was calculated. With these input values, and by the use of the grain-size distribution of the massive  
 510 layer, a value of  $0.015 \text{ m}^2\text{s}^{-1}$  was obtained for  $Q_b$ . See the supporting file D for a step by step  
 511 calculation. By setting the value of  $S_{rw}=100Q_b$ , the value of  $S_{rw}$  was calculated as  $1.5 \text{ m}^2\text{s}^{-1}$ . With  
 512 these values, the thick massive layer at Sbirri plots, in the discriminatory diagram of Fig. 10, in the  
 513 same position of the thick massive layers from experiments.

514 For the stratified layer with small-dune bedforms at Sbirri, by using as input parameters in Equation  
 515 (14) those calculated by Mele et al. (2015) (see Tab. 2), a sedimentation rate  $S_r$  of  $0.59 \text{ kgm}^{-2}\text{s}^{-1}$   
 516 resulted. The step by step calculation of  $S_r$  is found in the supporting file E. By means of Equation  
 517 (12),  $S_r$  was converted into  $S_{rw}$ , with a value of  $0.00047 \text{ m}^2\text{s}^{-1}$  (see Tab.1). As for  $Q_b$  of the small-dune  
 518 bedform layer at Sbirri, it was calculated by means of Equations (5) and (10) and using as input  
 519 parameters those calculated by Mele et al. (2011), see Tab. 2. The resulting value of  $Q_b$  was  $0.022$   
 520  $\text{m}^2\text{s}^{-1}$ . With the calculated values of  $Q_b$  and  $S_{rw}$ , the layer with small-dune bedforms at Sbirri plots in  
 521 the discriminatory diagram of Fig. 10 in the field of stratified deposits, with a lower  $S_{rw}/Q_b$  ratio  
 522 compared to the experimental deposits. The structures of the stratified layer at Sbirri have a much  
 523 longer wavelength, about 2 m, and height, about 0.15 m, (see Fig. 12c) with respect to the ripple-like  
 524 structures of experimental deposits. The position of the small-dune bedforms at Sbirri confirms the  
 525 inverse relationship between bedform scale and  $S_{rw}/Q_b$  ratio, and extends the range of the diagram of  
 526 Fig. 10 to lower  $S_{rw}/Q_b$  values and larger bedforms.

527 With the data obtained on the massive and stratified deposits at Sbirri, an evaluation of the *MER* of  
 528 the B2 phase of Agnano Monte Spina can be obtained. An estimation of *MER* can be obtained by the  
 529 total volume of deposits divided by the duration  $t$  of deposit aggradation. An approximate volume of

530 0.15 km<sup>3</sup> has been proposed for the B2 deposit by Dell'Erba (2004). As for  $t$  (the time of deposit  
 531 aggradation), it is reasonable to hypothesize that the stratified layer at Sbirri represents deposition  
 532 from the continuous aggradation of the lateral current that was fed from the eruption column. The  
 533 time of deposit aggradation was calculated as  $t = d / S_r / \rho_{dep}$ , where  $d = 1.5$  m, is the thickness of the  
 534 stratified layer and  $\rho_{dep}$  is 1248 kgm<sup>-3</sup>, as measured in the laboratory after extracting a known volume  
 535 of deposit in the field and weighing it.  $S_r$ , as obtained previously, was set to 0.59 kgm<sup>-2</sup>s<sup>-1</sup>. The  
 536 resulting  $MER$  is about  $6 \times 10^7$  kg/s. In this estimation, the assumption is made that neither  
 537 contemporaneous ash fallout nor additional layers (unrecorded in the field) were formed during B2  
 538 phase. The value is therefore to be considered as a minimal  $MER$ .

539 To constrain this figure of the  $MER$  of B2, a complementary method was applied. It is based on the  
 540 assumption that the massive layer at Sbirri is the result of the direct sedimentation from the lateral  
 541 zone of the impact, which was followed, in stratigraphic continuity, by the formation of the stratified  
 542 layer. Following the hypothesis that the eruption column collapsed completely and, as introduced  
 543 earlier, in its ascending-descending trajectory it did not entrain air and followed Bernoulli's equation,  
 544  $V_0$ ,  $\rho_{mix}$  and  $C_0$  at the vent are respectively equal to the velocity  $U_f$ , density  $\rho_{mix}$  and particle  
 545 volumetric concentration  $C_1$  of the eruption column at the impact on the ground.

$$546 \quad MER / \pi r^2 = V_0 \rho_s C_1 \quad (16)$$

547 is the  $MER$  per unit area at the vent, with  $\pi r^2$  the vent cross-sectional area. When air entrainment is  
 548 neglected in the eruption column, density doesn't change along the ascending descending trajectory,  
 549 and the mass flow rate per unit area at the vent is equivalent to the area specific mass flow rate  $Q_s$  of  
 550 the gas-particle mixture feeding the massive layer at the impact

$$551 \quad Q_s = \rho_s C_1 U_f \quad (17).$$

552 From the equality of (16) and (17)

$$553 \quad MER = \rho_s C_1 U_f \pi r_0^2 \quad (18).$$

554 By assuming that particles fall at the same descent velocity of the collapsing column (meaning a  
 555 mechanical coupling between particles and fluid in the dense fountain), it results that  $U_f = A_r/C_1$   
 556 where  $A_r$  is the aggradation rate of the massive deposit. By means of Equation (12), the aggradation  
 557 rate  $A_r$  is obtained by the previously calculated value of  $S_{rw}$  of the massive layer.  $A_r$  is about 1.53 m/s  
 558 and the falling velocity of the fountain  $U_f = A_r/C_1$  is about 170 m/s. As for  $C_1$ , the value of 0.009  
 559 previously calculated for the massive layer was used. From the equality of Equations (16) and (17),  
 560  $U_f$  represents also the exit velocity  $V_0$  of the gas-particle mixture at the vent.

561 Finally, using, in equation (3), the values of  $\rho_{f1}$  and  $V_0$ , and setting the value of  $Fr'_0$  between 1 and  
 562 3, which is the range proposed by Dellino et al. (2014) for completely collapsing eruption columns,  
 563 a range of  $MER$  was obtained, which is shown on Fig. 13 as a function of the Froude number. The  
 564  $MER$  value of  $6 \times 10^7 \text{ kgs}^{-1}$  previously obtained (by dividing the total mass of deposit by the time of  
 565 layer aggradation), falls inside the range of Fig. 13 and corresponds with a vent radius of 85 m. Both  
 566 the radius, the  $MER$  and the exit velocity are compatible with the values generally considered in the  
 567 literature (Dellino et al., 2010 and 2014). It is to note that the specific mass flow rate at the vent is in  
 568 the order of  $3 \times 10^3 \text{ kgm}^{-2}\text{s}^{-1}$ , which is in the range of values obtained in experiments, further justifying  
 569 the scaling between natural and experimental cases.

570

### 571 ***5.1.2 Use of the discriminatory diagram with the distal and ultra-distal facies of deposits for the*** 572 ***reconstruction of fine-ash deposition of B2***

573 The distal facies of B2 cropping out at Verdolino has a  $Q_b$  value of  $0.0029 \text{ m}^{-2}\text{s}^{-1}$  and  $S_{rw}$  of  $0.00015$   
 574  $\text{m}^{-2}\text{s}^{-1}$  (see Tab. 1). They were calculated, respectively, by means of Equations (10) and (5) and of  
 575 Equation (14). Data from Mele et al. (2011) were used as input (see Tab. 2). The deposit shows  
 576 coarse-ash ripple structures with a wavelength about 0.5 m and height about 0.05 m (Fig. 12d), which  
 577 are a little larger than the ripple beds of the experiments. The Verdolino deposit, by plotting, in the  
 578 diagram of Fig. 10, in a position characterized by a slightly lower value of  $S_{rw}/Q_b$  when compared to

579 the ripple beds formed by experiments, confirms the inverse relationship between  $S_{rw}/Q_b$  and bedform  
580 scale.

581 As for the current that formed the fine-ash structureless deposit of the ultra-distal facies of B2  
582 cropping out at Camaldoli, it is expected that a pyroclastic density current, due to the continuous air  
583 entrainment along runout, expands during its down-flow evolution. It can be hypothesized, therefore,  
584 that the thickness of the current, at Camaldoli, could not have been lower than that calculated by Mele  
585 et al. (2015) at Verdolino (see Tab. 2), which is about 300m. This value can be used as the minimum  
586 height  $H$  of the current at Camaldoli. Following the approach of Lajoie et al. (1989) and Dellino et  
587 al. (2008), who suggest that in steady conditions  $C = d/H$ , a particle volumetric concentration of  $C =$   
588 0.0006 results at Camaldoli, which can be considered as a maximum.  $P_{nsusp}$  was calculated  
589 numerically as 0.83 by using the values of  $d$ ,  $C$  and  $H$  as  $y_0$ ,  $C_{tot}$  and  $H_T$  in equation (15). By the use  
590 of  $P_{nsusp}$ ,  $\rho_{mix}$  and the grain-size distribution at Camaldoli, and setting the condition  $S_{rw} = 100Q_b$  as  
591 the minimum value for obtaining a structureless fine-ash bed (since the deposit does not show any  
592 shear structure), the combined use of Equations (10) and (14) resulted in values of  $u_*$ ,  $\tau$ ,  $S_{wr}$  and  $Q_b$   
593 at Camaldoli about  $0.43 \text{ ms}^{-1}$ ,  $0.34 \text{ Pa}$ ,  $6.7 \times 10^{-5} \text{ m}^2\text{s}^{-1}$  and  $6.7 \times 10^{-7} \text{ m}^2\text{s}^{-1}$ , respectively (see Tab. 1).

594 With the values of  $Q_b$  and  $S_{rw}$ , the fine-ash structureless bed of Camaldoli plots in the low-left corner  
595 of diagram of Fig. 10, in a similar position of the fine-ash structureless beds of the experiments.

596 When  $S_{rw}$  is transformed into  $S_r$  by means of Equations (7) and (12), a value of the sedimentation rate  
597 of  $9.6 \times 10^{-2} \text{ kgm}^{-2}\text{s}^{-1}$  results. The aggradation rate  $A_r$  is  $6.7 \times 10^{-5} \text{ ms}^{-1}$ , and the deposition time  $t$  is about  
598 3000 s. Step by step calculation can be found in the supporting file F. Such a value of deposition time  
599 must be considered as a minimum since, by hypothesizing a thicker current, the concentration would  
600 be lower, the sedimentation rate would further decrease, and the deposition time would increase.

601 This means that in ultra-distal localities, as at Camaldoli, the area is engulfed, for almost an hour,  
602 with a fine-ash cloud moving slowly on the ground. While the cloud loses much of its mechanical  
603 and thermal destructive power along runout, it is still rich of very harmful to breath fine material,  
604 which could compromise human survival also in zones well outside the Campi Flegrei caldera.

605

606 *5.2 Rheology and mobility of Mercato massive undercurrent*

607 The stratigraphy of Mercato eruption at Vesuvius (Mele et al., 2011), is made up of alternating Plinian  
608 fallout layers and pyroclastic density current deposits that are well exposed in the northern sector of  
609 the volcano, at the rim and down the slope of Monte-Somma (Fig. 14a). The deposits considered in  
610 this paper are from the first phase of the eruption and consist of a m-thick poorly-sorted massive layer  
611 of lapilli and scattered bombs and blocks supported in an ashy matrix (Fig. 14b), and a distal deposit  
612 made of a stratified lapilli deposit exhibiting small-dune bedforms (Fig. 14c). Both result from a  
613 dense eruption column that impacted, upon collapse, on the uneven slope of the volcanic cone. The  
614 massive layer is interpreted as fed and formed directly from the sedimentation of the collapsing  
615 column, whereas the stratified layer is interpreted as formed by the dilute current derived from lateral  
616 evolution of the flow after impaction on the ground. Along the gullies departing from the volcano,  
617 the thick massive deposit, aided by the steeper slope, moved as a highly concentrated debris flow,  
618 independent from the faster, more dilute, overcurrent. The contemporaneous occurrence of a massive  
619 underflow together with a dilute overcurrent has already been reported at other localities (Druitt et  
620 al., 2002; Gernon et al., 2013; Breard and Lube, 2017).

621 The massive facies terminates at the base of the cone where, as a response of the drastic decrease of  
622 the slope angle, it “freezes” abruptly in a 2 m thick layer (Fig.14 d). This suggests that the flow’s  
623 internal yield strength did not allow down-slope flowage until a minimum shear stress was overcome,  
624 similarly to a Bingham-plastic fluid (Furbish, 1987). Assuming that flow density was not much  
625 different than deposit density  $\rho_{dep}$ , i.e. 1400 kg/m<sup>3</sup> (as it was calculated by considering a known  
626 volume of deposits and weighting it), the yield strength  $\tau_0$ , e.g. the minimum shear stress for the  
627 down-slope movement of the debris flow, was evaluated as 700 Pa by means of the  
628 relation  $\tau_0 = \rho_{dep} g \sin \alpha d$ , with  $\alpha = 1.5^\circ$  as slope angle at deposit freezing and  $d = 2\text{m}$  as deposit  
629 thickness.

630 The stratified layer, when cropping out on the gentle slope of the volcano flank, shows a thickness  
 631 of 0.5 m; whereas, when it occurs along the steeper slopes of the gullies departing from the crater  
 632 rim, is split into two parts by the intercalated massive layer of the underflow, resulting in a 0.15 m  
 633 thick basal part (Fig. 14e) and a 0.35 m thick top part (Fig. 14f). It exhibits small dunes of lapilli and  
 634 ash of 1 m wavelength and 0.1 m in height (Fig. 14e). The intercalated massive layer is 1.6 m thick.  
 635 Such interpenetrating stratigraphy between the massive and the tractional layer, which can be  
 636 followed along the gullies, can be explained by the different speed of the two flows. The dilute current  
 637 overtook the slower basal flow, initiating sedimentation of the tractional layer. After the massive flow  
 638 passed over, the aggradation of the tractional layer continued as long as the current was fed from the  
 639 source. Lots of examples of such a kind of stratigraphy are reported for turbidites and debris flows  
 640 (Talling et al., 2004). Basing on the assumption that the underflow and the overcurrent were formed  
 641 almost contemporaneously by the eruption column impact at the crater rim, the speed of the underflow  
 642 can be estimated by dividing the distance from the crater rim by the total time the underflow took to  
 643 reach the tractive layer at the “split” location. The total time was obtained by summing up the time  
 644 that the overcurrent took to reach the “split” location, and the time that the overcurrent took to form,  
 645 by aggradation, the stratified layer found under the massive one, which is 0.15 m thick. The former  
 646 time, 52 s, was calculated by the distance, 1.6 km, divided by the height integrated average speed of  
 647 the overcurrent, ca. 30 ms<sup>-1</sup>. The average speed of the overcurrent was calculated by means of the  
 648 turbulent boundary layer shear flow approximation  $\overline{u(H)} = u_* \left( \frac{1}{k} \ln \frac{H}{k_s} + 6 \right)$ , where  $k_s$  is ground  
 649 roughness measured as 0.5 m in the field (as resulted from averaging the height of roughness elements  
 650 of the substrate) and  $\overline{u(H)}$  is the height integrated average speed.  $u_*$  and  $H$  were obtained by data of  
 651 Mele et al. (2011), see Tab. 2. The time of aggradation of the stratified layer found under the massive  
 652 layer, 375 s, resulted from  $d/S_r/\rho_{dep}$ , where  $S_r$  is the sedimentation rate of the dilute overcurrent at  
 653 the “split” location. Step by step calculation of  $S_r$  of the overcurrent is found in the supporting file G.  
 654 The total time that the underflow took to reach the split location was 428 s, corresponding in a height-

655 averaged velocity of the massive undercurrent of 3.7 m/s, which is much slower than the that of the  
 656 overcurrent (about 30 m/s) as it is expected for a highly concentrated debris flow moving downslope.  
 657 By inverting the equation of the height-averaged velocity of a Bingham-plastic fluid:

$$658 \quad \bar{u}(y) = d \left( \frac{\rho g \sin \alpha d}{3\mu} - \frac{\tau_0}{2\mu} \right) \quad (19)$$

659 and using in Equation (19) the yield strength  $\tau_0$  previously obtained, the thickness  $d$  of the massive  
 660 layer and a slope angle at the “split” location of 6.5°, a viscosity  $\mu$  of 200 Pas is obtained, which  
 661 completes the rheological characterization of the massive underflow. Such a rheology is compatible  
 662 with that of sedimentary debris flows, to which massive pyroclastic flows have been sometimes  
 663 compared in the literature (Fink et al., 1981; Major and Pierson, 1992; Major and Iverson, 1999;  
 664 Palladino and Valentine, 1995; Capra et al., 2018). While such flows maintain mobility only inside  
 665 channels and stop at the gully apron, where the slope abruptly decreases, still they are highly  
 666 hazardous in the proximal zones at the foot of the volcano, because of the high mass load, thickness  
 667 and temperature, that render them highly destructive.

668  $Q_b$  and  $S_{rw}$  were calculated for the stratified layer of Mercato cropping below the massive one (fig.  
 669 14e). Equations (5) and (10) were used for the calculation of  $Q_b$ , with a resulting value 0.011 m<sup>2</sup>s<sup>-1</sup>.  
 670 Equation (14) was used for the calculation of  $S_r$ . As input, data from Mele et al. (2011) were used,  
 671 see Tab. 2. By means of Equations (7) and (12)  $S_r$  was transformed into  $S_{rw}$  with a value of 0.0004  
 672 m<sup>2</sup>s<sup>-1</sup>. With these values of  $Q_b$  and  $S_{rw}$  the layer with small-dune bedforms of Mercato plots, in the  
 673 diagram of Fig. 10, in the field of stratified deposits, with a lower value of  $S_{rw}/Q_b$  when compared to  
 674 experimental deposits. The lapilli and ash small-dune structures of Mercato stratified deposit, with a  
 675 wavelength about 1 m, and a thickness of about 0.1m (Fig. 14e) further confirm, by the position on  
 676 the diagram, the inverse relationship between  $S_{rw}/Q_b$  and bedform scale.

677

## 678 **6. Summary and conclusion**

679 The present paper aims at defining the threshold ratio of the sedimentation and bedload transportation  
680 rates that marks the limit between the formation of massive vs stratified deposits by suspension  
681 currents. By the result of experiments, the threshold ratio was set at  $S_{rw}/Q_b = 100$ . The same ratio  
682 separates structureless fine-ash beds from layers with shear structures. By the experiments it was  
683 demonstrated that, in the range of tested fluid dynamics, fine ash never shows shear structures. It is  
684 in fact difficult to move by traction very fine material that exposes a limited surface to the basal shear  
685 stress of a current.

686 The paucity, in the literature, of data of the kind of those proposed here depends on the relatively few  
687 experiments on deposition from suspension currents. Such experiments need a scale and equipment  
688 quite different from those of the more frequently designed “flume” type, which generally do not  
689 include the contemporaneous measurement of suspension sedimentation and bedload transportation.  
690 For the same reason, the experiments carried out in this paper, due to the time length of design,  
691 construction, execution, analysis and post processing, cannot be repeated multiple times, in the life  
692 span of a research-project, as to cover the complete range of eruption conditions leading to the  
693 formation of pyroclastic deposits. A comparison with the natural cases shows that the main physical  
694 parameters of experiments are in the same order of magnitude of real pyroclastic density currents and  
695 that the facies model of the experimental deposits of Fig. 8 is similar to what has been reported for  
696 some ignimbrites, including the deposits of the Agnano Monte-Spina eruption of Campi Flegrei  
697 studied in this paper. The facies model could be of help in the interpretation of the lateral facies  
698 architecture of pyroclastic density currents deposits.

699 The  $S_{rw}/Q_b$  ratio has been calculated also for natural deposits of Campi Flegrei and Vesuvius, which  
700 were plotted in the diagram of Fig. 10. The diagram was able to correctly discriminate between  
701 massive and stratified deposits, legitimating the scaling between experimental and natural deposits.  
702 The deposits with shear structures of the Mercato eruption at Vesuvius and of the Sbirri and Verdolino  
703 localities at Campi Flegrei plot, on Fig. 10, in the field of stratified deposits with lower  $S_{rw}/Q_b$  values  
704 with respect to the experimental deposits, and allow to extend the range of structures covered in the

705 diagram. Starting from the smaller ripple-like structures and, by gradually increasing bedform  
706 wavelength and height, going up to deposits with small dunes, the larger are the structures, the lower  
707 is the  $S_{rw}/Q_b$  value, with a regular decrease at increasing bedform scale. This may suggest that it  
708 happens in natural deposits, as it happens in the experimental ones, that the scale of bedforms is  
709 inversely related to  $S_{rw}/Q_b$ .

710 The application of the classificatory diagram to natural deposits demonstrated also the potentiality in  
711 the reconstruction of physical characteristics of the eruptions, which have major implications on  
712 hazard assessment. Application to the proximal facies of B2 deposits of Campi Flegrei represents a  
713 new way of reconstructing *MER* feeding the pyroclastic density currents of past eruptions.

714 Application to the distal fine-ash deposits of Campi Flegrei allowed the determination, by means of  
715 the threshold limit between massive vs tractional sedimentation, of the duration of passage of the  
716 unbreathable current in distal localities, which could last for almost an hour in the case of the Agnano  
717 Monte-Spina eruption.

718 Application to the deposits of the Mercato eruption of Vesuvius, has helped demonstrating that the  
719 eruptive column can lead to the formation of both a slow, concentrated, underflow and a dilute  
720 overcurrent. It has been demonstrated that the underflow behaves as a debris flow, controlled by a  
721 Bingham rheology with a yield strength of 700 Pa and viscosity of 200 Pas, and moves at a few meters  
722 per second, independent from the dilute, faster, overcurrent.

723 It is reasonable to hypothesize that the numerical limits of the discriminatory diagram of Fig. 10  
724 would likely be more precisely defined by additional applications to a larger number of pyroclastic  
725 deposits, with the aim to also further extend the range up to larger-scale dunes of more intense  
726 pyroclastic density currents. Further applications of the diagram could also prove whether it can be  
727 of some use also in the study of other density currents as, for example, turbiditic currents.

728

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734

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1009

#### 1010 **Figure captions**

1011

1012 Fig. 1. Sketch design of the experimental apparatus with description of the main parts. Modified after  
1013 Dellino et al., 2019. A: Packs of interconnected compressed-gas bottles. B: hub connecting the hoses  
1014 to the base plate of the conduit. C: the conduit.

1015

1016 Fig. 2. Grain-size distribution of the material used for experiments. The size is expressed in  $\phi$  units  
1017 ( $\phi = -\log_2 D$ , where  $D$  is size in mm). The relative fractions of components are shown. The xx symbol  
1018 means crystals. A: multicomponent grain-size distribution of the coarse material coming from the  
1019 Veusvius composition. B: grain-size distribution of the fine glassy material coming from the Campi  
1020 Flegrei composition. Modified after Dellino et al., 2019.

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1022

1023 Fig. 3. On the left, a photo showing the passage of an experimental current, with values of average  
1024 velocity  $u_{avg}$ , shear velocity,  $u_*$ , thickness and dynamics pressure in the inset. On the right, the  
1025 sensor network recording pressure.

1026 Fig. 4. General view of deposit dispersal area with traces of the sampling areas.

1027 Fig. 5. Display mount showing phases of the experiment and massive deposits. A: formation of the  
1028 dense gas-particle fountain at the conduit exit. B: Collapse of the fountain and impact on the ground.  
1029 C: The massive deposit formed by an experiment fed with the coarse grain size at the centre of the  
1030 impact. The lateral undulations (marked by arrows) formed by the incipient shear at the lateral zone  
1031 of the impact are visible. D: Cuttings showing the internal sections of the massive deposit. The field  
1032 of view is 0.5 m across. E. A massive deposit formed at the centre of the impact of an experiment fed  
1033 with the fine material. F: The structureless deposit formed by an experiment fed with the fine grain  
1034 size.

1035

1036 Fig. 6. Display mount showing the fully turbulent density current and deposits structures. A: the  
1037 density current fed from the fountain collapse. B: widely spaced coarse ash cm thick ripple-like  
1038 structures. The field of view is 1m across. C: close up view of the passage of a density current showing  
1039 the shearing effect on the particles at the base. D: close up view of the successive frame of figure 6c  
1040 with continuing shear of particles at the base of the current. E: Narrowly spaced coarse ash cm thick  
1041 ripple-like structures. F: widely spaced mm thick medium ash ripple-like structures.

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1043 Fig. 7. Far reaches of the current. A: end of the fountain feeding a current, resulting in washload  
1044 transportation. B: the submillimetric veil of ash representing the washload.

1045

1046 Fig. 8. Sketch of fountain/current evolution along runout with the lateral facies evolution of deposits.

1047

1048 Fig. 9. Diagram showing the functional relationship of the ratio between the concentration in the  
 1049 fountain and that in the suspension current ( $C1/C2$ ) against the ratio of normalized bedload  
 1050 transportation and total fallout rate ( $Q_b^*/S_{fr}$ ). The equation of the best-fit law interpolating data, and  
 1051 the correlation coefficient, are inset.

1052

1053 Fig. 10.  $S_{rw} - Q_b$  discrimination diagram between massive and stratified deposits. The separation line  
 1054 between massive and stratified deposits is set at the threshold of  $S_{rw}=100Q_b$ . Both the experimental  
 1055 and natural deposits are plot. The relationship between structures spacing and  $S_{fr}/Q_b$  is shown. The  
 1056 dashed line ( $S_{fr}=4Q_b$ ) represents the limit between narrowly spaced and widely spaced ripple like  
 1057 structures of the experimental deposits. FAMass = Fine ash massive deposit; LAMass = Lapilli and  
 1058 ash massive deposit; CARip.narr = Coarse Ash narrowly spaced ripple-like structures; CARip.wide  
 1059 = Coarse ash widely spaced ripple like structures; MARip.wide = Medium Ash widely spaced ripple-  
 1060 like structures; LASmDune = Lapilli and Ash small dune; FASrless = Fine Ash structureless beds.

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1062 Fig. 11. Digital elevation model, rendered as shaded relief, with illumination from the North of  
 1063 Vesuvius volcano and Campi Flegrei caldera in southern Italy.

1064

1065 Fig.12. Display mount showing the Agnano Monte-Spina B2 deposits used as a model application.  
 1066 A: Digital Elevation Model map showing the Campi Flegrei caldera and the location of deposits. B:  
 1067 A massive deposit with alignments of coarse clasts, representing the impact of a gas-particle  
 1068 fountain, cropping out at Vallone del Corvo. C: Massive (bottom) and stratified (top) deposits,  
 1069 cropping at Sbirri, formed by the lateral zone of the fountain impact. The small dunes circled in the  
 1070 image have a 2 m spacing and 0.15 m height. D: On top of the hammer B2 deposits at Verdolino  
 1071 showing widely-spaced ripple-like structures (circled). Structure spacing is about 0.5 m and height  
 1072 about 0.05 m. E: The coin points to the fine-ash layer at Camaldoli.

1073

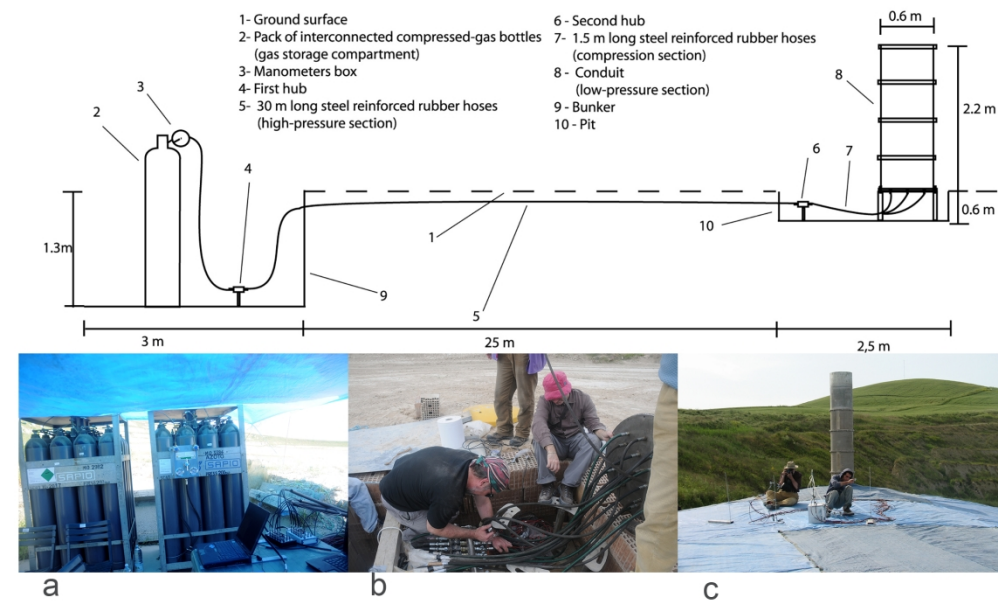
1074 Fig. 13. Diagram showing the relationship between Mass Eruption Rate (MER) and Desimetric  
1075 Froude number ( $Fr_0'$ ) for the B2 phase of the Agnano Monte-Spina eruption, as obtained by model  
1076 application. The line intersecting the curve represents the value of MER calculated by the deposit  
1077 volume data of Dell'Erba (2004).

1078

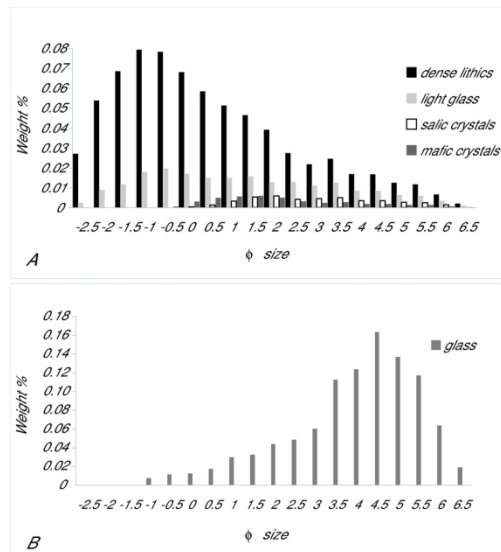
1079 Fig. 14. Display mount showing the Mercato eruption deposits used as a case study for model  
1080 application. The hammer used as a scale for field photographs, is 30 cm long. A: Digital Elevation  
1081 Model map of Vesuvius. The location of deposits are marked, together with the zone of impact of the  
1082 eruptive fountain. B: the massive layer formed at the impact of the eruptive fountain. C: the stratified  
1083 layer formed at the bedload of the suspension current. D: the massive deposit emplaced by freezing  
1084 at the break-in slope. E: the stratified layer at the base of the massive layer at the split location. Small  
1085 dune structures are circled. Structure spacing is avbout 1 m, height about 0.1 m. F: the stratified layer  
1086 on top of the massive layer at the split location

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Sketch design of the experimental apparatus with description of the main parts. Modified after Dellino et al., 2019. A: Packs of interconnected compressed-gas bottles. B: hub connecting the hoses to the base plate of the conduit. C: the conduit.

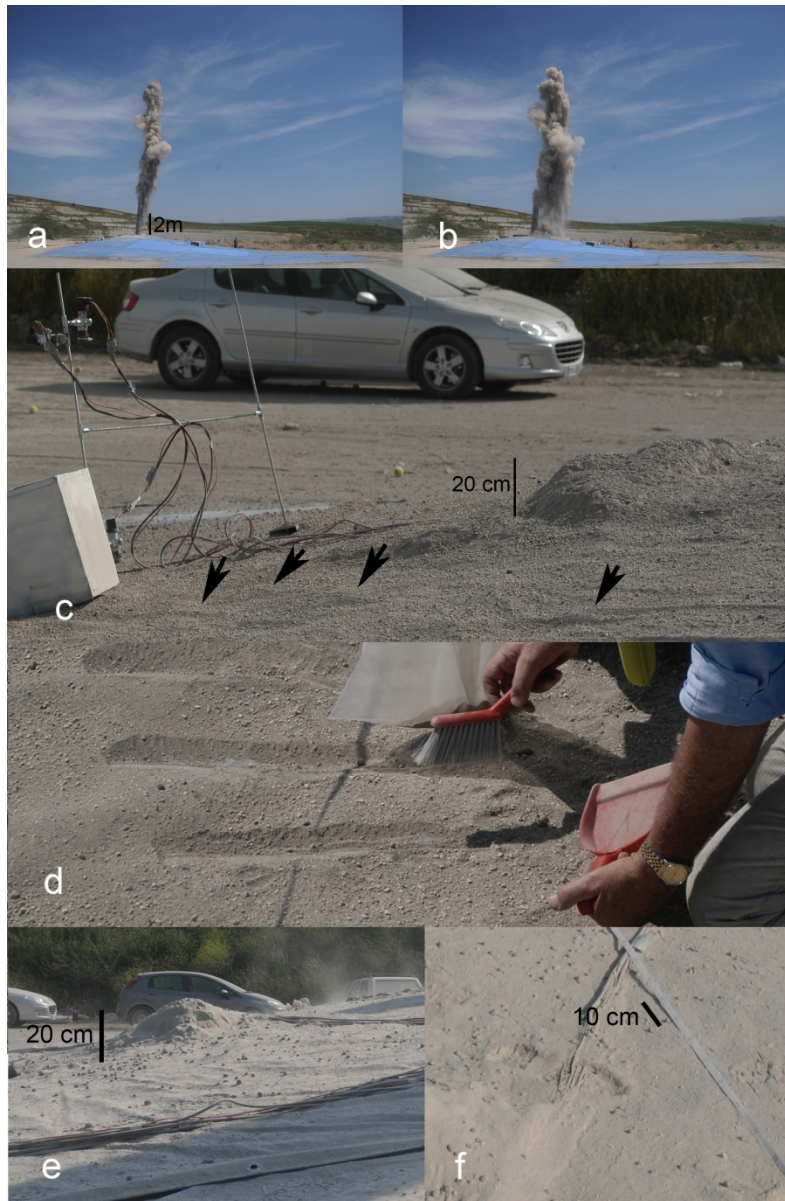


Grain-size distribution of the material used for experiments. The size is expressed in  $\Phi$  units ( $\Phi = -\log_{2}D$ , where  $D$  is size in mm). The relative fractions of components are shown. The xx symbol means crystals. A: multicomponent grain-size distribution of the coarse material coming from the Veusvius composition. B: grain-size distribution of the fine glassy material coming from the Campi Flegrei composition. Modified after Dellino et al., 2019.

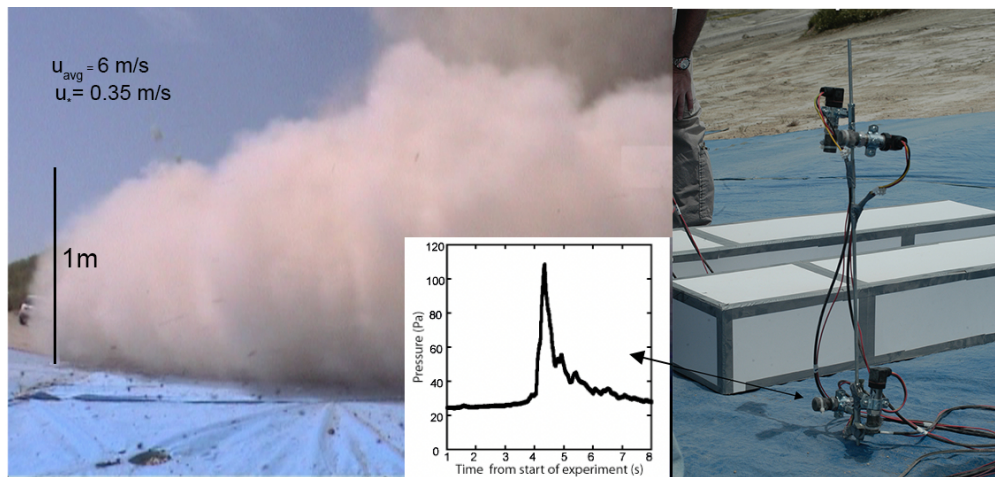
185x102mm (300 x 300 DPI)



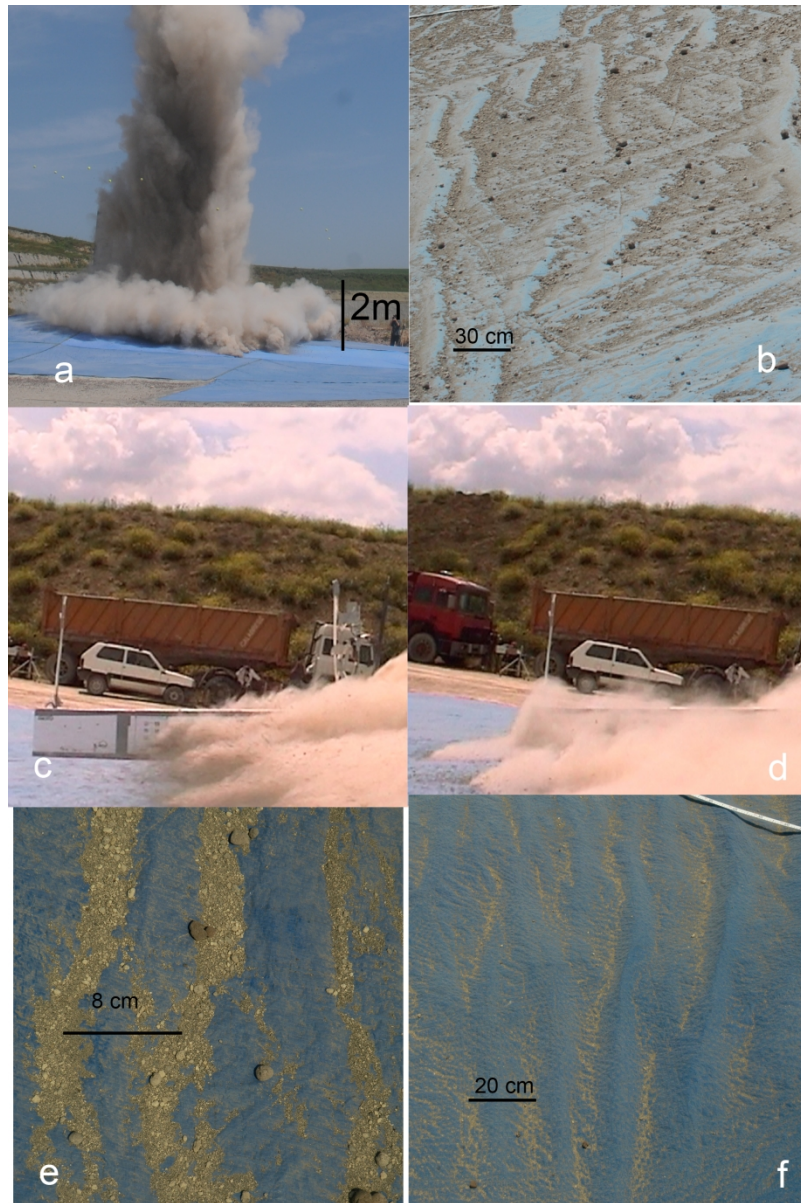
General view of deposit dispersal area with traces of the sampling areas.



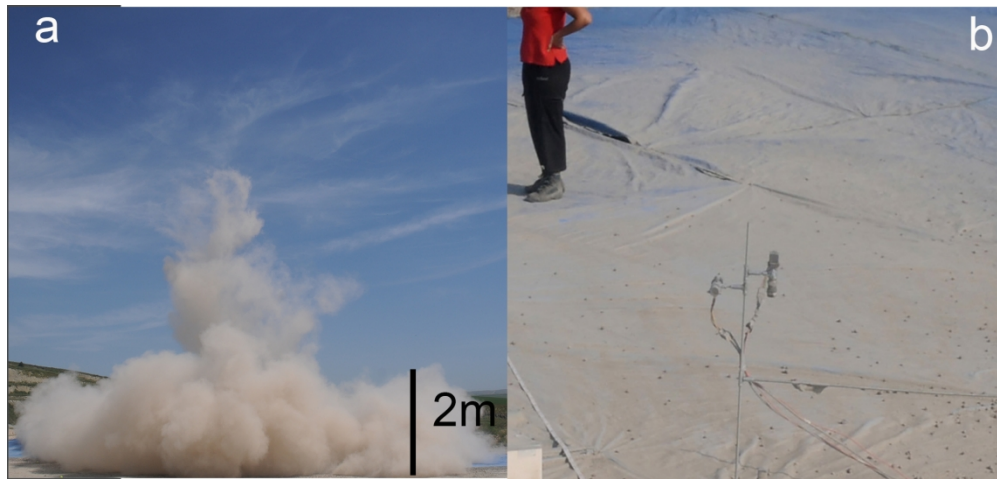
Display mount showing phases of the experiment and massive deposits. A: formation of the dense gas-particle fountain at the conduit exit. B: Collapse of the fountain and impact on the ground. C: The massive deposit formed by an experiment fed with the coarse grain size at the centre of the impact. The lateral undulations (marked by arrows) formed by the incipient shear at the lateral zone of the impact are visible. D: Cuttings showing the internal sections of the massive deposit. The field of view is 0.5 m across. E: A massive deposit formed at the centre of the impact of an experiment fed with the fine material. F: The structureless deposit formed by an experiment fed with the fine grain size.



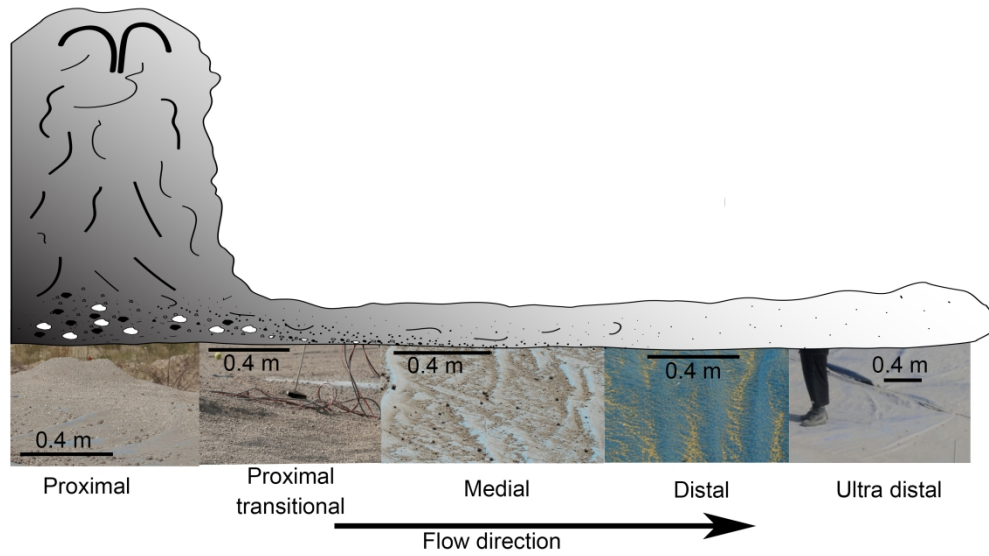
On the left, a photo showing the passage of an experimental current, with values of average velocity  $u_{avg}$ , shear velocity,  $u_*$ , thickness and dynamics pressure in the inset. On the right, the sensor network recording pressure.



Display mount showing the fully turbulent density current and deposits structures. A: the density current fed from the fountain collapse. B: widely spaced coarse ash cm thick ripple-like structures. The field of view is 1m across. C: close up view of the passage of a density current showing the shearing effect on the particles at the base. D: close up view of the successive frame of figure 6c with continuing shear of particles at the base of the current. E: Narrowly spaced coarse ash cm thick ripple-like structures. F: widely spaced mm thick medium ash ripple-like structures.



Far reaches of the current. A: end of the column feeding a current, resulting in washload transportation. B: the submillimetric veil of ash representing the washload.



Sketch of column/current evolution along runout with the lateral facies evolution of deposits.

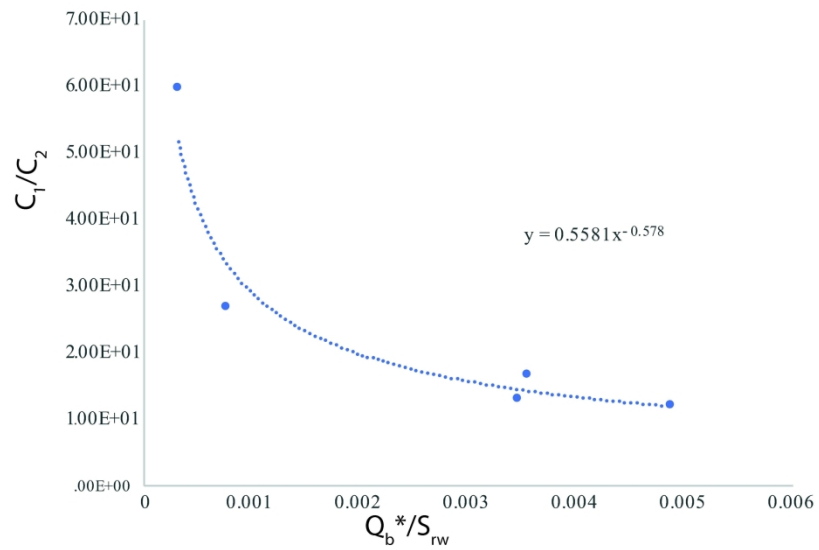
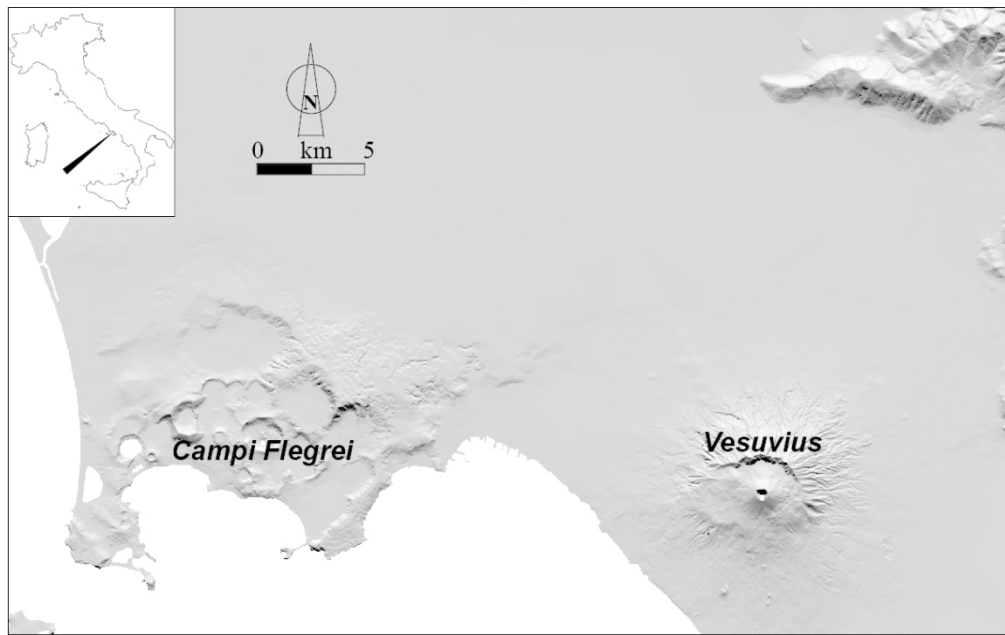


Diagram showing the functional relationship of the ratio between the concentration in the column and that in the suspension current ( $C_1/C_2$ ) against the ratio of normalized bedload transportation and total fallout rate ( $Q_b^*/S_{fr}$ ). The equation of the best-fit law interpolating data, and the correlation coefficient, are inset.

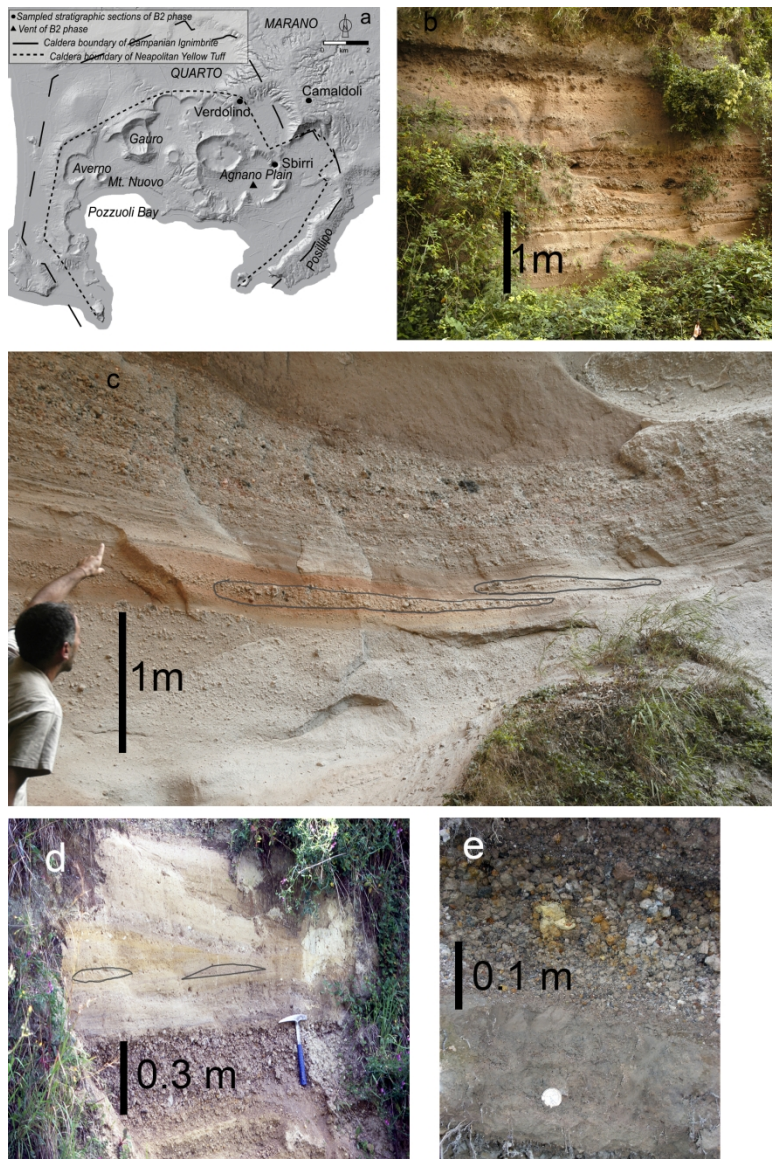
219x139mm (300 x 300 DPI)





Digital elevation model, rendered as shaded relief, with illumination from the North of Vesuvius volcano and Campi Flegrei caldera in southern Italy.

102x64mm (300 x 300 DPI)



Display mount showing the Agnano Monte-Spina B2 deposits used as an application. A: Digital Elevation Model map showing the Campi Flegrei caldera and the location of deposits. B: A massive deposit with alignments of coarse clasts, representing the impact of a gas-particle column, cropping out at Vallone del Corvo. C: Massive (bottom) and stratified (top) deposits, cropping at Sbirri, formed by the lateral zone of the fountain impact. The small dunes circled in the image have a 2 m spacing and 0.15 m height. D: On top of the hammer B2 deposits at Verdolino showing widely-spaced ripple-like structures (circled). Structure spacing is about 0.5 m and height about 0.05 m. E: The coin points to the fine-ash layer at Camaldoli.

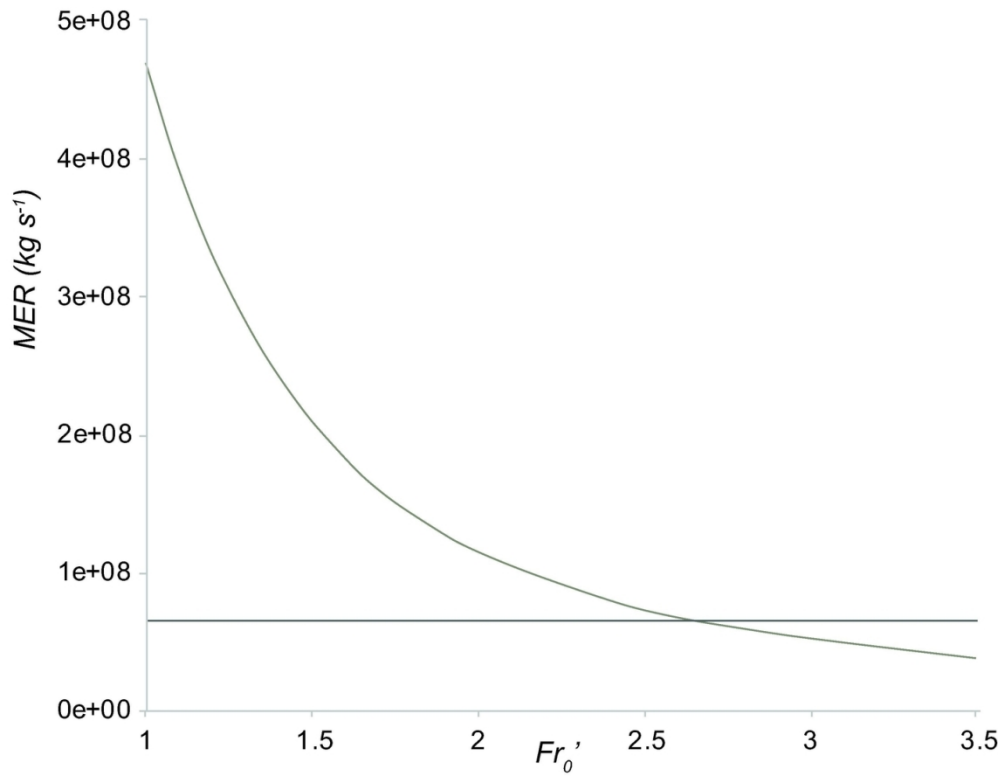
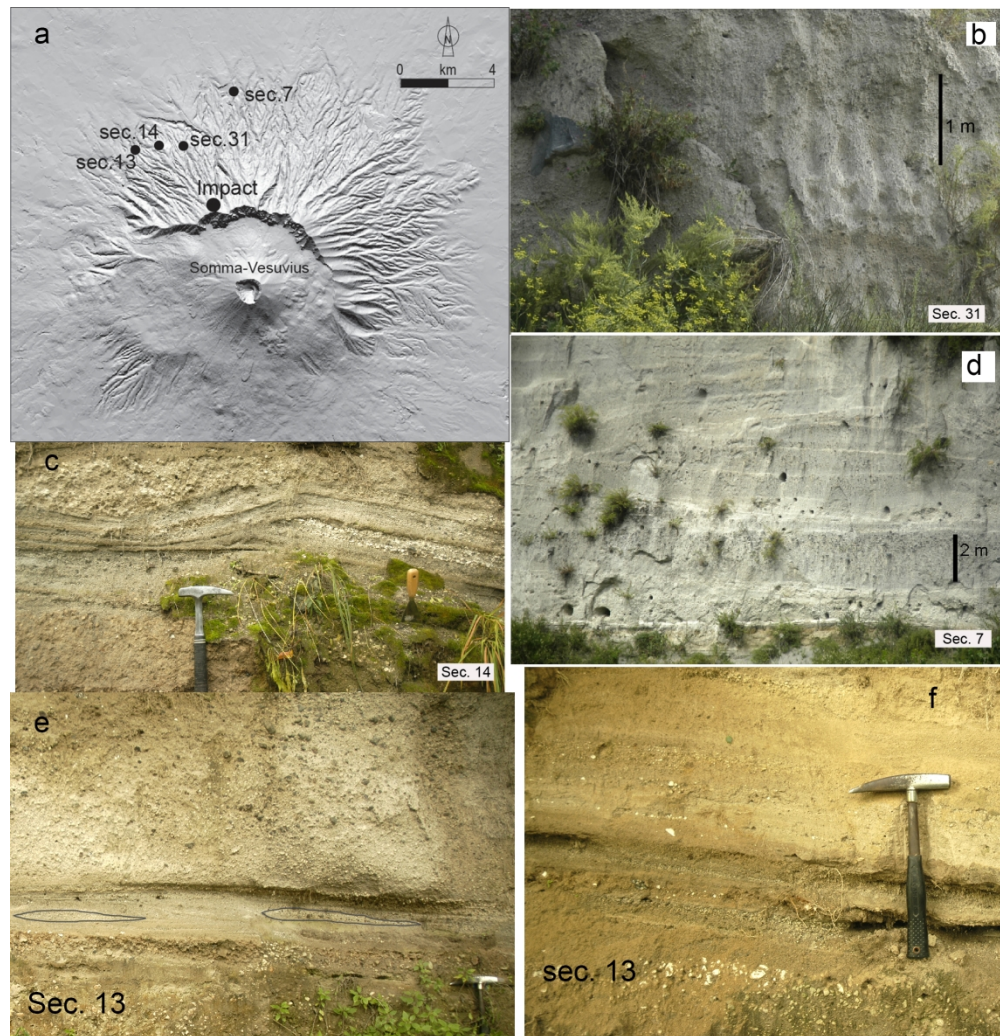


Diagram showing the relationship between Mass Eruption Rate (MER) and Desimetric Froude number ( $Fr_0'$ ) for the B2 phase of the Agnano Monte-Spina eruption, as obtained by model application. The line intersecting the curve represents the value of MER calculated by the deposit volume data of Dell'Erba (2004).

120x94mm (300 x 300 DPI)



Display mount showing the Mercato eruption deposits used as a case study. The hammer used as a scale for field photographs, is 30 cm long. A: Digital Elevation Model map of Vesuvius. The location of deposits are marked, together with the zone of impact of the eruptive fountain. B: the massive layer formed at the impact of the eruptive fountain. C: the stratified layer formed at the bedload of the suspension current. D: the massive deposit emplaced by freezing at the break-in slope. E: the stratified layer at the base of the massive layer at the split location. Small dune structures are circled. Structure spacing is about 1 m, height about 0.1 m. F: the stratified layer on top of the massive layer at the split location

<i>Nr</i>	<i>S<sub>r</sub></i> ( <i>kgm<sup>-2</sup>s<sup>-1</sup></i> )	<i>Q<sub>b</sub></i> ( <i>m<sup>2</sup>s<sup>-1</sup></i> )	<i>S<sub>rw</sub></i> ( <i>m<sup>2</sup>s<sup>-1</sup></i> )	<i>S<sub>rw</sub>/Q<sub>b</sub></i>	<i>u<sub>avg</sub></i> ( <i>ms<sup>-1</sup></i> )	<i>u*</i> ( <i>ms<sup>-1</sup></i> )	<i>D<sub>sm</sub></i> ( <i>mm</i> )	<i>C<sub>1</sub></i> (-)	<i>C<sub>2</sub></i> (-)	<i>Qb*</i> ( <i>m<sup>2</sup>s<sup>-1</sup></i> )	<i>structure</i>	<i>Origin</i>
1	1.10E+02	4.70E-04	1.00E-01	2.13E+02	-	0.26	0.75	5.95E-02	4.54E-03	3.55E-03	LAMass	Exp
2	2.20E+02	1.52E-03	2.00E-01	1.32E+02	-	0.44	0.65	3.19E-02	2.60E-03	4.92E-03	LAMass	Exp
3	1.38E+02	9.27E-05	1.25E-01	1.35E+03	-	0.14	0.53	6.28E-02	8.83E-04	4.67E-04	LAMass	Exp
4	3.90E+01	1.90E-04	3.00E-02	1.58E+02	-	0.26	0.57	3.35E-02	2.00E-03	3.61E-03	LAMass	Exp
5	1.60E+02	2.00E-03	2.00E-01	1.00E+02	-	0.4	0.08	1.89E-02	7.00E-04	6.96E-04	FAMass	Exp
6	7.80E-01	6.07E-05	3.50E-04	5.77E+00	5.52	0.36	0.64	-	3.54E-03	-	CARip.narr	Exp
7	1.17E+00	1.07E-03	9.75E-03	9.11E+00	9.74	0.66	0.63	-	4.60E-03	-	CARip.wide	Exp
8	6.80E-01	1.02E-04	2.60E-04	2.55E+00	8.83	0.55	0.64	-	1.30E-03	-	CARip.wide	Exp
9	1.90E-01	2.36E-06	8.20E-05	3.47E+01	4.77	0.30	0.64	-	7.83E-04	-	CARip.narr	Exp
10	6.10E-01	7.40E-06	2.80E-04	3.78E+01	3.05	0.27	0.77	-	2.41E-03	-	CARip.narr	Exp
11	1.90E-02	1.54E-05	1.50E-05	9.74E-01	4.8	0.27	0.073	-	5.950E-04	-	MARip.wide	Exp
12	4.20E-01	1.15E-03	2.80E-04	2.43E-01	13	0.73	0.79	-	3.50E-03	-	CARip.wide	Exp
13	2.20E-02	1.77E-07	1.78E-05	1.01E+02	3.56	0.17	0.044	-	5.950E-04	-	FAStrless	Exp
14	1.91E+03	1.53E-02	1.53E-00	1.00E+02	-	1.18	0.9	-	9.00E-03	-	LAMass	Nat (Sbirri)
15	5.09E-01	2.20E-02	4.70E-04	2.14E-02	63,29	2.72	0.8	-	1.00 E-03	-	LASmDune	Nat (Sbirri)
16	1.90E-01	2.90 E-03	1.53E-04	5.28E-02	31.3	1.43	0.24	-	1.00 E-03	-	CARip.wide	Nat (Verdolino)
17	8.20E-02	6.66E-07	6.66E-05	1.00E+02	9.50	0.43	0.073	-	6.00E-04	-	FAStrless	Nat (Camaldoli)
18	5.89E-01	1.10E-02	4.00E-04	3.64E-02	30	2.00	0.5	-	1.75 E-03	-	LASmDune	Nat (Mercato)

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Table 2. Parameters used for the calculation of sedimentation and bedload transportation rates. For symbols see symbol list.

<i>Eruption</i>	<i>Location</i>	$H_{tot}$ (m)	$u^*$ (m/s)	$\tau$ (Pa)	$\rho_f$ (kgm <sup>-3</sup> )	$C_2$ (-)
Agnano-Monte Spina	Sbirri	1685	2.72	18.63	2.52	0.0010
	Verdolino	290	1.43	6.02	2.94	0.0010
Mercato	Sandwich location	380	1.40	14.92	3.54	0.0013

Parameters of the dilute currents used as input for the calculation of sedimentation and transportation rates for the deposits of Campi Flegrei and Vesuvius eruptions used as case studies.

Table 2. Results of model calculations: flow field variables of PDCs. For symbol description see Table A.1.

Symbol	description	dimension
$A_r$	Rate of deposit aggradation	$\text{ms}^{-1}$
$C_0$	Reference known concentration (0.75)	-
$C_d$	Particle drag coefficient	-
$C_i$	Particle volumetric concentration of the $i_{\text{th}}$ size	-
$C_{\text{tot}}$	Total particle volumetric concentration	-
$C_1$	Particle volumetric concentration in the fountain	-
$C_2$	Particle volumetric concentration in the dilute current	-
$D$	particle size	mm
$D_i$	diameter of the $i_{\text{th}}$ size fraction	mm
$D_{sm}$	diameter of the median size	mm
$d$	Layer thickness	mm
$Fr'_0$	Densimetric Froude number	-
$g$	Gravity acceleration (9.81)	$\text{ms}^{-2}$
$g'$	Reduced gravity	$\text{ms}^{-2}$
$H_{dep}$	Deposit thickness	m
$H_{\text{tot}}$	Total flow thickness	m
$k$	Von Karman constant (0.4)	-
$k_s$	Ground roughness	m
$MER$	Mass Eruption rate	$\text{kgs}^{-1}$
$P_n$	Particle Rouse number	-
$P_n^*$	Normalized Rouse number	-
$P_{\text{avg}}$	Average Rouse number of the particle load	-
$P_{ni}$	Rouse number of the $i_{\text{th}}$ particle-size class	-
$P_{\text{nsusp}}$	Rouse number at maximum suspension capacity	-
$Q_b$	Bedload transportation rate per unit width	$\text{m}^2\text{s}^{-1}$
$Q_b^*$	Bedload transportation rate per unit width normalized to one mm	$\text{m}^2\text{s}^{-1}$
$Q_s$	Area specific mass flow rate	$\text{Kgm}^{-2}\text{s}^{-1}$
$r$	Vent diameter	m
$Re^*$	Reynolds' number	-
$S_r$	Sedimentation rate	$\text{kgm}^{-2}\text{s}^{-1}$
$S_{ri}$	Sedimentation rate of the $i_{\text{th}}$ particle-size fraction	$\text{kgm}^{-2}\text{s}^{-1}$
$S_{rw}$	Sedimentation rate per unit width	$\text{m}^2\text{s}^{-1}$
$t$	Aggradation time	s
$u^*$	Flow shear velocity	$\text{ms}^{-1}$
$u_{\text{avg}}$	Height averaged flow velocity	$\text{ms}^{-1}$
$U_f$	Impact speed of the fountain	$\text{ms}^{-1}$
$V_0$	Exit velocity at the vent	$\text{ms}^{-1}$
$w_t$	Particle terminal velocity	$\text{ms}^{-1}$
$w_{ti}$	Terminal velocity of the $i_{\text{th}}$ particle-size class	$\text{ms}^{-1}$
$y$	Flow vertical coordinate	m
$y_0$	Specific height of $C_0$	-
$\phi$	Unit of grain-size distribution ( $\phi = -\log_2 d$ )	-
$\phi_i$	Weight fraction of the $i_{\text{th}}$ size class	Weight%
$\gamma$	Proportion of particles settling from suspension	-
$\gamma_i$	Proportion of particles settling from suspension of the $i_{\text{th}}$ size class	-

$\theta$	Shields' number	-
$\theta_c$	Critical Shield's number	-
$\theta_{median}$	Shield's number of the median size	-
$\mu$	viscosity	Pas
$\rho_{atm}$	Atmosphere density	kgm <sup>-3</sup>
$\rho_{dep}$	Deposit density ( $0.6\rho_s$ )	kgm <sup>-3</sup>
$\rho_{mix}$	Density of the gas-particle mixture	kgm <sup>-3</sup>
$\rho_g$	Gas density	kgm <sup>-3</sup>
$\rho_s$	Particle density	kgm <sup>-3</sup>
$\rho_{si}$	Density of the $i$ th particle-size class	kgm <sup>-3</sup>
$\tau$	Shear stress at the base of the current	Pa
$\tau_{ri}$	minimum shear of the $i_{th}$ size fraction	Pa
$\tau_{rsm}$	minimum shear stress of the particle mixture	Pa
$\tau_{rsm}^*$	Shield number normalized to the mixture	-

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List of symbols, with description and physical dimension

$u_*$ (ms <sup>-1</sup> )	0.137									
$\rho_s$ average (kgm <sup>-3</sup> )	1815	Component	$D_i$ (m)	$D_i/D_{sm}$	$\tau_{ri}$ (Pa)	$v/\tau_{ri}$	$W_i^*$	$Q_{bi}$ (m <sup>2</sup> s <sup>-1</sup> )	$\phi_i$	$\rho_{si}$ (kgm <sup>-3</sup> )
$C_{tot}$ (-)	0.0628	dense lithic	5.66E-03	9.725	0.603	4.188	1.057	1.814E-07	0.0083	2588
$\rho_{mix}$ (kgm <sup>-3</sup> )	125.6	dense lithic	4.00E-03	6.873	0.475	5.314	1.539	3.709E-07	0.0116	2588
$\tau$ (Pa)	2.357	dense lithic	2.83E-03	4.863	0.366	6.902	2.154	1.104E-06	0.0247	2588
$D_{sm}$ (mm)	0.53	dense lithic	2.00E-03	3.436	0.271	9.329	2.947	1.742E-06	0.0285	2588
$Re_*$ (-)	459	dense lithic	1.41E-03	2.423	0.201	12.574	3.788	2.265E-06	0.0288	2588
$\theta_i$ (-)	0.015	dense lithic	1.00E-03	1.718	0.161	15.719	4.434	5.425E-06	0.0590	2588
$\tau_{rsm}^*$ (-)	0.0112	dense lithic	7.07E-04	1.215	0.139	18.173	4.855	5.313E-06	0.0527	2588
$\tau_{rsm}$ (Pa)	0.1313	dense lithic	5.00E-04	0.859	0.127	19.906	5.118	4.713E-06	0.0444	2588
$S_r$ (kgm <sup>-2</sup> s <sup>-1</sup> )	138	dense lithic	3.54E-04	0.608	0.119	21.174	5.296	5.115E-06	0.0465	2588
$\rho_{dep}$ (kgm <sup>-3</sup> )	1100	dense lithic	2.50E-04	0.430	0.114	22.208	5.433	4.102E-06	0.0364	2588
$A_r$ (ms <sup>-1</sup> )	0.125	dense lithic	1.77E-04	0.304	0.109	23.129	5.549	3.397E-06	0.0295	2588
		dense lithic	1.25E-04	0.215	0.105	24.027	5.657	3.189E-06	0.0272	2588
		dense lithic	8.84E-05	0.152	0.101	24.934	5.762	3.549E-06	0.0297	2588
		dense lithic	6.25E-05	0.107	0.098	25.878	5.867	2.845E-06	0.0234	2588
		dense lithic	4.42E-05	0.076	0.094	26.870	5.973	2.946E-06	0.0238	2588
		dense lithic	3.13E-05	0.054	0.090	27.914	6.079	2.304E-06	0.0183	2588
		dense lithic	2.21E-05	0.038	0.087	29.028	6.188	2.535E-06	0.0197	2588
		dense lithic	1.56E-05	0.027	0.084	30.206	6.298	1.843E-06	0.0141	2588
		dense lithic	1.10E-05	0.019	0.080	31.453	6.409	7.305E-07	0.0055	2588
		vesicular glass	5.66E-03	9.725	0.603	4.188	1.057	3.242E-07	0.0148	686
		vesicular glass	4.00E-03	6.873	0.475	5.314	1.539	1.180E-06	0.0369	733
		vesicular glass	2.83E-03	4.863	0.366	6.902	2.154	1.477E-06	0.0330	783
		vesicular glass	2.00E-03	3.436	0.271	9.329	2.947	3.134E-06	0.0512	837
		vesicular glass	1.41E-03	2.423	0.201	12.574	3.788	4.127E-06	0.0525	894
		vesicular glass	1.00E-03	1.718	0.161	15.719	4.434	2.105E-06	0.0229	956
		vesicular glass	7.07E-04	1.215	0.139	18.173	4.855	1.812E-06	0.0180	1021
		vesicular glass	5.00E-04	0.859	0.127	19.906	5.118	2.597E-06	0.0244	1091

Sedimentology

vesicular glass	3.54E-04	0.608	0.119	21.174	5.296	1.580E-06	0.0144	1166
vesicular glass	2.50E-04	0.430	0.114	22.208	5.433	1.979E-06	0.0175	1245
vesicular glass	1.77E-04	0.304	0.109	23.129	5.549	1.273E-06	0.0111	1331
vesicular glass	1.25E-04	0.215	0.105	24.027	5.657	1.194E-06	0.0102	1422
vesicular glass	8.84E-05	0.152	0.101	24.934	5.762	1.329E-06	0.0111	1519
vesicular glass	6.25E-05	0.107	0.098	25.878	5.867	1.066E-06	0.0087	1623
vesicular glass	4.42E-05	0.076	0.094	26.870	5.973	1.103E-06	0.0089	1734
vesicular glass	3.13E-05	0.054	0.090	27.914	6.079	8.629E-07	0.0068	1853
vesicular glass	2.21E-05	0.038	0.087	29.028	6.188	9.494E-07	0.0074	1979
vesicular glass	1.56E-05	0.027	0.084	30.206	6.298	6.902E-07	0.0053	2115
vesicular glass	1.10E-05	0.019	0.080	31.453	6.409	2.736E-07	0.0021	2259
salic xx	1.00E-03	1.718	0.161	15.719	4.434	8.907E-08	0.0010	2400
salic xx	7.07E-04	1.215	0.139	18.173	4.855	2.483E-07	0.0025	2400
salic xx	5.00E-04	0.859	0.127	19.906	5.118	4.730E-07	0.0045	2400
salic xx	3.54E-04	0.608	0.119	21.174	5.296	6.500E-07	0.0059	2400
salic xx	2.50E-04	0.430	0.114	22.208	5.433	5.596E-07	0.0050	2400
salic xx	1.77E-04	0.304	0.109	23.129	5.549	4.702E-07	0.0041	2400
salic xx	1.25E-04	0.215	0.105	24.027	5.657	3.286E-07	0.0028	2400
salic xx	8.84E-05	0.152	0.101	24.934	5.762	3.658E-07	0.0031	2400
salic xx	6.25E-05	0.107	0.098	25.878	5.867	2.932E-07	0.0024	2400
salic xx	4.42E-05	0.076	0.094	26.870	5.973	3.036E-07	0.0024	2400
salic xx	3.13E-05	0.054	0.090	27.914	6.079	2.375E-07	0.0019	2400
salic xx	2.21E-05	0.038	0.087	29.028	6.188	2.613E-07	0.0020	2400
salic xx	1.56E-05	0.027	0.084	30.206	6.298	1.899E-07	0.0015	2400
salic xx	1.10E-05	0.019	0.080	31.453	6.409	7.528E-08	0.0006	2400
femic xx	1.00E-03	1.718	0.161	15.719	4.434	1.700E-07	0.0018	3280
femic xx	7.07E-04	1.215	0.139	18.173	4.855	6.615E-07	0.0066	3280
femic xx	5.00E-04	0.859	0.127	19.906	5.118	8.095E-07	0.0076	3280
femic xx	3.54E-04	0.608	0.119	21.174	5.296	1.228E-06	0.0112	3280
femic xx	2.50E-04	0.430	0.114	22.208	5.433	1.139E-06	0.0101	3280

femic xx	1.77E-04	0.304	0.109	23.129	5.549	7.783E-07	0.0068	3280
femic xx	1.25E-04	0.215	0.105	24.027	5.657	5.393E-07	0.0046	3280
femic xx	8.84E-05	0.152	0.101	24.934	5.762	6.005E-07	0.0050	3280
femic xx	6.25E-05	0.107	0.098	25.878	5.867	4.814E-07	0.0040	3280
femic xx	4.42E-05	0.076	0.094	26.870	5.973	4.985E-07	0.0040	3280
femic xx	3.13E-05	0.054	0.090	27.914	6.079	3.898E-07	0.0031	3280
femic xx	2.21E-05	0.038	0.087	29.028	6.188	4.289E-07	0.0033	3280
femic xx	1.56E-05	0.027	0.084	30.206	6.298	3.118E-07	0.0024	3280
femic xx	1.10E-05	0.019	0.080	31.453	6.409	1.236E-07	0.0009	3280

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$Q_b$  (m<sup>2</sup>s<sup>-1</sup>) 9.543E-05

Supporting file A. Step-by-step calculation of the bedload transportation rate  $Q_b$  of an experimental massive deposit formed at the column impact. In yellow the main parametres. For symbols see the symbol list.

$u^*$ (ms <sup>-1</sup> )	0.66									
$\rho_s$ average (kgm <sup>-3</sup> )	2367	Component	$D_i$ (m)	$D_i/D_{sm}$	$\tau_{ri}$ (Pa)	$\tau/\tau_{ri}$	$W_i^*$	$Q_{bi}$ (m <sup>2</sup> s <sup>-1</sup> )	$\phi_i$	$\rho_{si}$ (kgm <sup>-3</sup> )
$C_{tot}$ (-)	0.0046	dense lithic	5.66E-03	8.998	0.475	11.114	3.436	8.332E-06	0.0161	2588
$\rho_{mix}$ (kgm <sup>-3</sup> )	12.11	dense lithic	4.00E-03	6.359	0.373	14.147	4.128	2.082E-05	0.0335	2588
$\tau$ (Pa)	5.274	dense lithic	2.83E-03	4.499	0.285	18.535	4.912	3.706E-05	0.0501	2588
$D_{sm}$ (mm)	0.629	dense lithic	2.00E-03	3.180	0.209	25.183	5.790	5.872E-05	0.0673	2588
$Re^*$ (-)	256	dense lithic	1.41E-03	2.242	0.157	33.552	6.584	7.089E-05	0.0715	2588
$\theta_i$ (-)	0.01	dense lithic	1.00E-03	1.590	0.128	41.147	7.125	8.101E-05	0.0754	2588
$\tau_{rsm}^*$ (-)	0.0075	dense lithic	7.07E-04	1.124	0.113	46.878	7.459	7.881E-05	0.0701	2588
$\tau_{rsm}$ (Pa)	0.1090	dense lithic	5.00E-04	0.795	0.104	50.921	7.665	7.359E-05	0.0637	2588
$S_r$ (kgm <sup>-2</sup> s <sup>-1</sup> )	1.17	dense lithic	3.54E-04	0.563	0.098	53.938	7.807	6.150E-05	0.0523	2588
$\rho_{dep}$ (kgm <sup>-3</sup> )	1200	dense lithic	2.50E-04	0.397	0.093	56.459	7.917	4.996E-05	0.0419	2588
$A_r$ (ms <sup>-1</sup> )	0.000975	dense lithic	1.77E-04	0.281	0.090	58.750	8.012	3.261E-05	0.0270	2588
		dense lithic	1.25E-04	0.199	0.086	61.011	8.102	2.791E-05	0.0229	2588
		dense lithic	8.84E-05	0.141	0.083	63.311	8.188	3.609E-05	0.0293	2588
		dense lithic	6.25E-05	0.099	0.080	65.713	8.275	2.300E-05	0.0184	2588
		dense lithic	4.42E-05	0.070	0.077	68.241	8.362	2.124E-05	0.0169	2588
		dense lithic	3.13E-05	0.050	0.074	70.904	8.448	1.628E-05	0.0128	2588
		dense lithic	2.21E-05	0.035	0.072	73.744	8.537	1.531E-05	0.0119	2588
		dense lithic	1.56E-05	0.025	0.069	76.747	8.625	9.563E-06	0.0074	2588
		dense lithic	1.10E-05	0.017	0.066	79.925	8.714	3.203E-06	0.0024	2588
		vesicular glass	5.66E-03	8.998	0.475	11.114	3.436	1.088E-06	0.0021	686
		vesicular glass	4.00E-03	6.359	0.373	14.147	4.128	4.336E-06	0.0070	733
		vesicular glass	2.83E-03	4.499	0.285	18.535	4.912	8.042E-06	0.0109	783
		vesicular glass	2.00E-03	3.180	0.209	25.183	5.790	1.395E-05	0.0160	837
		vesicular glass	1.41E-03	2.242	0.157	33.552	6.584	1.607E-05	0.0162	894
		vesicular glass	1.00E-03	1.590	0.128	41.147	7.125	1.200E-05	0.0112	956
		vesicular glass	7.07E-04	1.124	0.113	46.878	7.459	1.048E-05	0.0093	1021
		vesicular glass	5.00E-04	0.795	0.104	50.921	7.665	1.575E-05	0.0136	1091

vesicular glass	3.54E-04	0.563	0.098	53.938	7.807	1.230E-05	0.0105	1166
vesicular glass	2.50E-04	0.397	0.093	56.459	7.917	1.339E-05	0.0112	1245
vesicular glass	1.77E-04	0.281	0.090	58.750	8.012	1.801E-05	0.0149	1331
vesicular glass	1.25E-04	0.199	0.086	61.011	8.102	1.708E-05	0.0140	1422
vesicular glass	8.84E-05	0.141	0.083	63.311	8.188	2.209E-05	0.0179	1519
vesicular glass	6.25E-05	0.099	0.080	65.713	8.275	1.408E-05	0.0113	1623
vesicular glass	4.42E-05	0.070	0.077	68.241	8.362	1.300E-05	0.0103	1734
vesicular glass	3.13E-05	0.050	0.074	70.904	8.448	9.968E-06	0.0078	1853
vesicular glass	2.21E-05	0.035	0.072	73.744	8.537	9.369E-06	0.0073	1979
vesicular glass	1.56E-05	0.025	0.069	76.747	8.625	5.853E-06	0.0045	2115
vesicular glass	1.10E-05	0.017	0.066	79.925	8.714	1.961E-06	0.0015	2259
salic xx	1.41E-03	2.242	0.157	33.552	6.584	6.102E-07	0.0006	2400
salic xx	7.07E-04	1.124	0.113	46.878	7.459	3.514E-06	0.0031	2400
salic xx	5.00E-04	0.795	0.104	50.921	7.665	5.269E-06	0.0046	2400
salic xx	3.54E-04	0.563	0.098	53.938	7.807	6.150E-06	0.0052	2400
salic xx	2.50E-04	0.397	0.093	56.459	7.917	6.604E-06	0.0055	2400
salic xx	1.77E-04	0.281	0.090	58.750	8.012	4.208E-06	0.0035	2400
salic xx	1.25E-04	0.199	0.086	61.011	8.102	4.130E-06	0.0034	2400
salic xx	8.84E-05	0.141	0.083	63.311	8.188	5.341E-06	0.0043	2400
salic xx	6.25E-05	0.099	0.080	65.713	8.275	3.403E-06	0.0027	2400
salic xx	4.42E-05	0.070	0.077	68.241	8.362	3.143E-06	0.0025	2400
salic xx	3.13E-05	0.050	0.074	70.904	8.448	2.410E-06	0.0019	2400
salic xx	2.21E-05	0.035	0.072	73.744	8.537	2.265E-06	0.0018	2400
salic xx	1.56E-05	0.025	0.069	76.747	8.625	1.415E-06	0.0011	2400
salic xx	1.10E-05	0.017	0.066	79.925	8.714	4.740E-07	0.0004	2400
femic xx	1.00E-03	1.590	0.128	41.147	7.125	3.027E-06	0.0028	3280
femic xx	7.07E-04	1.124	0.113	46.878	7.459	8.757E-06	0.0078	3280
femic xx	5.00E-04	0.795	0.104	50.921	7.665	5.269E-06	0.0046	3280
femic xx	3.54E-04	0.563	0.098	53.938	7.807	1.351E-05	0.0115	3280
femic xx	2.50E-04	0.397	0.093	56.459	7.917	1.082E-05	0.0091	3280

femic xx	1.77E-04	0.281	0.090	58.750	8.012	7.178E-06	0.0059	3280
femic xx	1.25E-04	0.199	0.086	61.011	8.102	6.758E-06	0.0055	3280
femic xx	8.84E-05	0.141	0.083	63.311	8.188	8.740E-06	0.0071	3280
femic xx	6.25E-05	0.099	0.080	65.713	8.275	5.569E-06	0.0045	3280
femic xx	4.42E-05	0.070	0.077	68.241	8.362	5.143E-06	0.0041	3280
femic xx	3.13E-05	0.050	0.074	70.904	8.448	3.943E-06	0.0031	3280
femic xx	2.21E-05	0.035	0.072	73.744	8.537	3.707E-06	0.0029	3280
femic xx	1.56E-05	0.025	0.069	76.747	8.625	2.316E-06	0.0018	3280
femic xx	1.10E-05	0.017	0.066	79.925	8.714	7.756E-07	0.0006	3280

$Q_b$  (m<sup>2</sup>s<sup>-1</sup>) 1.079E-03

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Supporting file B. Step-by-step calculation of the bedload transportation rate  $Q_b$  of an experimental layer with ripple-like structures sedimented by the dilute turbulent suspension current formed as a lateral evolution of the column impact. In yellow the main parametres. For symbols see the symbol list.

$u_*$ (ms <sup>-1</sup> )	0.17								
$\rho_s$ average (kgm <sup>-3</sup> )	2241	Component	$D_i$ (m)	$D_i/D_{sm}$	$\tau_{ri}$ (Pa)	$\tau/\tau_{ri}$	$W_i^*$	$Q_{bi}$ (m <sup>2</sup> s <sup>-1</sup> )	$\phi_i$
$C_{tot}$ (-)	0.00060	vesicular glass	2.00E-03	45.455	0.323	0.229	2.607E-06	4.307E-14	0.0289
$\rho_{mix}$ (kgm <sup>-3</sup> )	2.55	vesicular glass	1.41E-03	32.045	0.255	0.289	7.480E-06	6.504E-14	0.0152
$\tau$ (Pa)	0.074	vesicular glass	1.00E-03	22.727	0.203	0.364	2.108E-05	2.016E-13	0.0167
$D_{sm}$ (mm)	0.044	vesicular glass	7.07E-04	16.068	0.161	0.459	5.995E-05	6.255E-13	0.0183
$Re_*$ (-)	0.90	vesicular glass	5.00E-04	11.364	0.128	0.579	1.704E-04	2.371E-12	0.0244
$\theta_i$ (-)	0.019	vesicular glass	3.54E-04	8.045	0.101	0.731	4.869E-04	8.044E-12	0.0289
$\tau_{rsm}^*$ (-)	0.0143	vesicular glass	2.50E-04	5.682	0.079	0.937	1.490E-03	4.016E-11	0.0472
$\tau_{rsm}$ (Pa)	0.0250	vesicular glass	1.77E-04	4.023	0.059	1.243	5.327E-03	1.714E-10	0.0563
$S_r$ (kgm <sup>-2</sup> s <sup>-1</sup> )	0.024	vesicular glass	1.25E-04	2.841	0.044	1.693	7.510E-02	3.527E-09	0.0822
$\rho_{dep}$ (kgm <sup>-3</sup> )	1344.6	vesicular glass	8.84E-05	2.009	0.034	2.201	2.203E-01	1.321E-08	0.1050
$A_r$ (ms <sup>-1</sup> )	1.78E-05	vesicular glass	6.25E-05	1.420	0.028	2.633	3.815E-01	3.062E-08	0.1405
		vesicular glass	4.42E-05	1.005	0.025	2.944	5.101E-01	5.087E-08	0.1746
		vesicular glass	3.13E-05	0.711	0.023	3.166	6.051E-01	4.503E-08	0.1303
		vesicular glass	2.21E-05	0.502	0.022	3.338	6.804E-01	3.340E-08	0.0859
		vesicular glass	1.56E-05	0.355	0.021	3.486	7.456E-01	1.540E-08	0.0361
		vesicular glass	1.10E-05	0.250	0.020	3.625	8.073E-01	4.372E-09	0.0095
							$Q_b$ (m <sup>2</sup> s <sup>-1</sup> )	1.77E-07	1

Supporting file C. Step-by-step calculation of the bedload transportation rate  $Q_b$  of an experimental fine-ash structureless thin layer formed by a dilute suspension current fed with the fine material. In yellow the main parameters. For symbols see the symbol list.

$u_*$ (ms <sup>-1</sup> )	1.176									
$\rho_s$ average (kgm <sup>-3</sup> )	1757	Component	$D_i$ (m)	$D_i/D_{sm}$	$\tau_{ri}$ (Pa)	$\tau/\tau_{ri}$	$W_i^*$	$Q_{bi}$ (m <sup>2</sup> s <sup>-1</sup> )	$\phi_i$	$\rho_{si}$ (kgm <sup>-3</sup> )
$C_{tot}$ (-)	0.009	vesicular glass	3.20E-02	35.556	0.882	25.422	5.817	1.789E-04	0.0199	474
$\rho_{mix}$ (kgm <sup>-3</sup> )	16.19	vesicular glass	2.24E-02	24.889	0.694	32.284	6.480	3.167E-04	0.0317	474
$\tau$ (Pa)	22.41	vesicular glass	1.60E-02	17.778	0.554	40.448	7.081	2.704E-04	0.0247	474
$D_{sm}$ (mm)	0.9	vesicular glass	1.13E-02	12.556	0.439	51.063	7.672	3.420E-04	0.0289	474
$Re_*$ (-)	856	vesicular glass	8.00E-03	8.889	0.348	64.414	8.229	4.169E-04	0.0328	531
$\theta_i$ (-)	0.01	vesicular glass	5.66E-03	6.289	0.273	81.975	8.769	5.732E-04	0.0423	596
$\tau_{rsm}^*$ (-)	0.0075	vesicular glass	4.00E-03	4.444	0.208	107.641	9.333	5.048E-04	0.0350	669
$\tau_{rsm}$ (Pa)	0.081	vesicular glass	2.83E-03	3.144	0.153	146.189	9.908	4.438E-04	0.0290	751
		vesicular glass	2.00E-03	2.222	0.115	194.053	10.387	6.202E-04	0.0387	842
		vesicular glass	1.41E-03	1.567	0.094	238.133	10.703	6.324E-04	0.0383	945
		vesicular glass	1.00E-03	1.111	0.083	270.351	10.886	4.042E-04	0.0240	1060
		vesicular glass	7.07E-04	0.786	0.076	293.354	10.999	4.126E-04	0.0243	1189
		vesicular glass	5.00E-04	0.556	0.072	310.601	11.076	4.050E-04	0.0237	1334
		vesicular glass	3.54E-04	0.393	0.069	324.923	11.136	4.209E-04	0.0245	1497
		vesicular glass	2.50E-04	0.278	0.066	338.161	11.187	4.117E-04	0.0238	1679
		vesicular glass	1.77E-04	0.197	0.064	351.067	11.235	3.923E-04	0.0226	1884
		vesicular glass	1.25E-04	0.139	0.062	364.357	11.282	3.651E-04	0.0210	2114
		vesicular glass	8.84E-05	0.098	0.059	378.174	11.328	2.129E-04	0.0122	2372
		vesicular glass	6.25E-05	0.069	0.057	392.744	11.374	2.388E-04	0.0136	2560
		vesicular glass	4.42E-05	0.049	0.055	408.138	11.420	2.888E-04	0.0164	2560
		vesicular glass	3.13E-05	0.035	0.053	424.353	11.466	2.781E-04	0.0157	2560
		vesicular glass	2.21E-05	0.025	0.051	441.629	11.512	3.698E-04	0.0208	2560
		vesicular glass	1.56E-05	0.017	0.049	459.868	11.559	3.405E-04	0.0191	2560
		vesicular glass	1.10E-05	0.012	0.047	479.133	11.605	3.025E-04	0.0169	2560
		vesicular glass	8.00E-06	0.009	0.045	497.556	11.646	2.061E-04	0.0115	2560
		vesicular glass	5.60E-06	0.006	0.043	519.185	11.692	1.345E-04	0.0074	2560
		vesicular glass	4.00E-06	0.004	0.041	540.573	11.735	7.021E-05	0.0039	2560
		vesicular glass	2.80E-06	0.003	0.040	564.315	11.780	3.957E-05	0.0022	2560

vesicular glass	2.00E-06	0.002	0.038	587.748	11.822	2.084E-05	0.0011	2560
salic xx	2.00E-03	2.222	0.115	194.053	10.387	1.227E-05	0.0008	2400
salic xx	1.41E-03	1.567	0.094	238.133	10.703	3.314E-05	0.0020	2400
salic xx	1.00E-03	1.111	0.083	270.351	10.886	1.110E-04	0.0066	2400
salic xx	7.07E-04	0.786	0.076	293.354	10.999	1.923E-04	0.0113	2400
salic xx	5.00E-04	0.556	0.072	310.601	11.076	3.004E-04	0.0176	2400
salic xx	3.54E-04	0.393	0.069	324.923	11.136	2.343E-04	0.0136	2400
salic xx	2.50E-04	0.278	0.066	338.161	11.187	2.488E-04	0.0144	2400
salic xx	1.77E-04	0.197	0.064	351.067	11.235	1.541E-04	0.0089	2400
salic xx	1.25E-04	0.139	0.062	364.357	11.282	1.041E-04	0.0060	2400
salic xx	8.84E-05	0.098	0.059	378.174	11.328	5.787E-05	0.0033	2400
salic xx	6.25E-05	0.069	0.057	392.744	11.374	5.709E-05	0.0033	2400
salic xx	4.42E-05	0.049	0.055	408.138	11.420	6.158E-05	0.0035	2400
salic xx	3.13E-05	0.035	0.053	424.353	11.466	5.352E-05	0.0030	2400
salic xx	2.21E-05	0.025	0.051	441.629	11.512	6.484E-05	0.0036	2400
salic xx	1.56E-05	0.017	0.049	459.868	11.559	5.482E-05	0.0031	2400
salic xx	1.10E-05	0.012	0.047	479.133	11.605	4.501E-05	0.0025	2400
salic xx	8.00E-06	0.009	0.045	497.556	11.646	2.851E-05	0.0016	2400
salic xx	5.60E-06	0.006	0.043	519.185	11.692	1.737E-05	0.0010	2400
salic xx	4.00E-06	0.004	0.041	540.573	11.735	8.511E-06	0.0005	2400
salic xx	2.80E-06	0.003	0.040	564.315	11.780	4.518E-06	0.0002	2400
salic xx	2.00E-06	0.002	0.038	587.748	11.822	2.247E-06	0.0001	2400
dense lithic	3.20E-02	35.556	0.882	25.422	5.817	1.106E-06	0.0001	2570
dense lithic	2.24E-02	24.889	0.694	32.284	6.480	1.232E-06	0.0001	2570
dense lithic	1.60E-02	17.778	0.554	40.448	7.081	1.346E-06	0.0001	2570
dense lithic	1.13E-02	12.556	0.439	51.063	7.672	2.166E-05	0.0018	2570
dense lithic	8.00E-03	8.889	0.348	64.414	8.229	6.296E-05	0.0050	2570
dense lithic	5.66E-03	6.289	0.273	81.975	8.769	2.280E-04	0.0168	2570
dense lithic	4.00E-03	4.444	0.208	107.641	9.333	2.128E-04	0.0148	2570
dense lithic	2.83E-03	3.144	0.153	146.189	9.908	3.335E-04	0.0218	2570

dense lithic	2.00E-03	2.222	0.115	194.053	10.387	3.428E-04	0.0214	2570
dense lithic	1.41E-03	1.567	0.094	238.133	10.703	2.996E-04	0.0181	2570
dense lithic	1.00E-03	1.111	0.083	270.351	10.886	4.593E-04	0.0273	2570
dense lithic	7.07E-04	0.786	0.076	293.354	10.999	4.294E-04	0.0253	2570
dense lithic	5.00E-04	0.556	0.072	310.601	11.076	3.600E-04	0.0211	2570
dense lithic	3.54E-04	0.393	0.069	324.923	11.136	3.813E-04	0.0222	2570
dense lithic	2.50E-04	0.278	0.066	338.161	11.187	2.835E-04	0.0164	2570
dense lithic	1.77E-04	0.197	0.064	351.067	11.235	1.897E-04	0.0109	2570
dense lithic	1.25E-04	0.139	0.062	364.357	11.282	1.493E-04	0.0086	2570
dense lithic	8.84E-05	0.098	0.059	378.174	11.328	7.863E-05	0.0045	2570
dense lithic	6.25E-05	0.069	0.057	392.744	11.374	7.812E-05	0.0044	2570
dense lithic	4.42E-05	0.049	0.055	408.138	11.420	8.479E-05	0.0048	2570
dense lithic	3.13E-05	0.035	0.053	424.353	11.466	7.408E-05	0.0042	2570
dense lithic	2.21E-05	0.025	0.051	441.629	11.512	9.017E-05	0.0051	2570
dense lithic	1.56E-05	0.017	0.049	459.868	11.559	7.655E-05	0.0043	2570
dense lithic	1.10E-05	0.012	0.047	479.133	11.605	6.309E-05	0.0035	2570
dense lithic	8.00E-06	0.009	0.045	497.556	11.646	4.009E-05	0.0022	2570
dense lithic	5.60E-06	0.006	0.043	519.185	11.692	2.451E-05	0.0014	2570
dense lithic	4.00E-06	0.004	0.041	540.573	11.735	1.204E-05	0.0007	2570
dense lithic	2.80E-06	0.003	0.040	564.315	11.780	6.405E-06	0.0004	2570
dense lithic	2.00E-06	0.002	0.038	587.748	11.822	3.194E-06	0.0002	2570

$Q_b$ (m <sup>2</sup> s <sup>-1</sup> )	0.0153
$S_{rw}$ (m <sup>2</sup> s <sup>-1</sup> )	1.53
$A_r$ (ms <sup>-1</sup> )	1.53
$U_f$ (ms <sup>-1</sup> )	170.60

Supporting file D. Step-by-step calculation of the bedload transportation rate  $Q_b$ , sedimentation rate per unit width  $S_{rw}$ , aggradation rate  $A_r$  and collapse falling velocity  $U_f$  for the massive layer at Sbirri. In yellow the main parametres. For symbols see the symbol list.

Input data	
$u^*$ ( $\text{ms}^{-1}$ )	2.72
$C_{\text{tot}}$	0.00103
$P_{\text{nsusp}}$	0.6

Component	$D_i$ (m)	$P_{ni}$	$P_{ni} * C_{\text{transi}}$	$P_{ni}/P_n^*$	$\rho_{si}$ ( $\text{kgm}^{-3}$ )	$w_i$ ( $\text{ms}^{-1}$ )	$\phi_i$	Vol. fraction ( $\text{m}^3\text{kg}^{-1}$ )	$C_{\text{transi}}$	$\gamma_i$	$\gamma_{\text{truei}}$	$S_{ri}$ ( $\text{kgm}^{-2}\text{s}^{-1}$ )	$\rho_{si} * \phi_i$ ( $\text{kgm}^{-3}$ )
vesicular glass	1.60E-02	4.266	2.52E-04	1.31E+00	422	4.76	1.43E-02	3.39E-05	5.91E-05	4.42E-06	6.03E-06	1.21E-02	6.037
vesicular glass	1.13E-02	3.663	2.26E-04	1.13E+00	474	4.08	1.68E-02	3.55E-05	6.18E-05	5.38E-06	5.12E-06	9.91E-03	7.958
vesicular glass	8.00E-03	3.137	1.40E-04	9.65E-01	531	3.50	1.36E-02	2.56E-05	4.46E-05	4.52E-06	4.63E-06	8.60E-03	7.227
vesicular glass	5.66E-03	2.681	1.11E-04	8.25E-01	596	2.99	1.41E-02	2.36E-05	4.12E-05	4.87E-06	4.61E-06	8.22E-03	8.407
vesicular glass	4.00E-03	2.286	6.61E-05	7.03E-01	669	2.55	1.11E-02	1.66E-05	2.89E-05	4.00E-06	4.93E-06	8.41E-03	7.425
vesicular glass	2.83E-03	1.944	8.54E-05	5.98E-01	751	2.17	1.89E-02	2.52E-05	4.39E-05	7.11E-06	8.70E-06	1.42E-02	14.184
vesicular glass	2.00E-03	1.649	1.08E-04	5.07E-01	842	1.84	3.16E-02	3.75E-05	6.54E-05	1.24E-05	1.32E-05	2.04E-02	26.608
vesicular glass	1.41E-03	1.395	9.32E-05	4.29E-01	945	1.56	3.62E-02	3.83E-05	6.68E-05	1.49E-05	1.44E-05	2.11E-02	34.198
vesicular glass	1.00E-03	1.175	5.86E-05	3.62E-01	1060	1.31	3.03E-02	2.86E-05	4.98E-05	1.31E-05	1.38E-05	1.91E-02	32.115
vesicular glass	7.07E-04	0.984	4.83E-05	3.03E-01	1189	1.10	3.35E-02	2.82E-05	4.91E-05	1.53E-05	1.48E-05	1.93E-02	39.836
vesicular glass	5.00E-04	0.817	2.96E-05	2.51E-01	1334	0.91	2.77E-02	2.08E-05	3.62E-05	1.35E-05	1.39E-05	1.69E-02	36.956
vesicular glass	3.54E-04	0.670	2.22E-05	2.06E-01	1497	0.75	2.85E-02	1.90E-05	3.32E-05	1.49E-05	1.55E-05	1.73E-02	42.660
vesicular glass	2.50E-04	0.540	1.67E-05	1.66E-01	1679	0.60	2.97E-02	1.77E-05	3.08E-05	1.68E-05	1.57E-05	1.59E-02	49.878
vesicular glass	1.77E-04	0.425	8.23E-06	1.31E-01	1884	0.47	2.09E-02	1.11E-05	1.93E-05	1.31E-05	1.30E-05	1.16E-02	39.379
vesicular glass	1.25E-04	0.325	4.82E-06	1.00E-01	2114	0.36	1.80E-02	8.51E-06	1.48E-05	1.28E-05	1.20E-05	9.18E-03	38.051
vesicular glass	8.84E-05	0.238	2.17E-06	7.32E-02	2372	0.27	1.24E-02	5.23E-06	9.11E-06	1.02E-05	1.08E-05	6.78E-03	29.409
vesicular glass	6.25E-05	0.159	1.26E-06	4.89E-02	2560	0.18	1.16E-02	4.53E-06	7.90E-06	1.21E-05	1.40E-05	6.36E-03	29.696
vesicular glass	4.42E-05	0.090	7.24E-07	2.77E-02	2560	0.10	1.18E-02	4.61E-06	8.04E-06	1.84E-05	1.85E-05	4.75E-03	30.208
vesicular glass	3.13E-05	0.044	2.40E-07	1.35E-02	2560	0.05	8.04E-03	3.14E-06	5.48E-06	1.86E-05	1.89E-05	2.37E-03	20.582
vesicular glass	2.21E-05	0.018	7.61E-08	5.60E-03	2560	0.02	6.14E-03	2.40E-06	4.18E-06	1.95E-05	1.74E-05	9.05E-04	15.718

## Sedimentology

vesicular glass	1.56E-05	0.007	1.52E-08	2.09E-03	2560	0.01	3.29E-03	1.29E-06	2.24E-06	1.25E-05	1.00E-05	1.95E-04	8.422
vesicular glass	1.10E-05	0.002	1.63E-09	7.38E-04	2560	0.00	9.98E-04	3.90E-07	6.80E-07	4.11E-06	2.88E-06	1.99E-05	2.555
salic xx	2.83E-03	5.161	1.15E-06	1.59E+00	2400	5.76	3.08E-04	1.28E-07	2.24E-07	1.39E-08	1.91E-08	2.64E-04	0.739
salic xx	2.00E-03	4.176	1.72E-06	1.28E+00	2400	4.66	5.67E-04	2.36E-07	4.12E-07	3.15E-08	9.33E-08	1.04E-03	1.361
salic xx	1.41E-03	3.363	8.45E-06	1.03E+00	2400	3.75	3.46E-03	1.44E-06	2.51E-06	2.38E-07	6.22E-07	5.60E-03	8.304
salic xx	1.00E-03	2.683	3.45E-05	8.25E-01	2400	2.99	1.77E-02	7.38E-06	1.29E-05	1.52E-06	1.65E-06	1.18E-02	42.480
salic xx	7.07E-04	2.108	2.74E-05	6.48E-01	2400	2.35	1.79E-02	7.46E-06	1.30E-05	1.95E-06	2.60E-06	1.47E-02	42.960
salic xx	5.00E-04	1.623	3.48E-05	4.99E-01	2400	1.81	2.95E-02	1.23E-05	2.14E-05	4.13E-06	4.32E-06	1.88E-02	70.800
salic xx	3.54E-04	1.222	2.29E-05	3.76E-01	2400	1.36	2.58E-02	1.08E-05	1.87E-05	4.75E-06	4.86E-06	1.59E-02	61.920
salic xx	2.50E-04	0.898	1.34E-05	2.76E-01	2400	1.00	2.06E-02	8.58E-06	1.50E-05	5.09E-06	5.09E-06	1.22E-02	49.440
salic xx	1.77E-04	0.640	6.97E-06	1.97E-01	2400	0.71	1.50E-02	6.25E-06	1.09E-05	5.09E-06	5.65E-06	9.68E-03	36.000
salic xx	1.25E-04	0.438	4.61E-06	1.35E-01	2400	0.49	1.45E-02	6.04E-06	1.05E-05	6.95E-06	6.97E-06	8.17E-03	34.800
salic xx	8.84E-05	0.284	2.07E-06	8.75E-02	2400	0.32	1.00E-02	4.17E-06	7.26E-06	6.99E-06	7.85E-06	5.97E-03	24.000
salic xx	6.25E-05	0.172	1.17E-06	5.29E-02	2400	0.19	9.35E-03	3.90E-06	6.79E-06	9.84E-06	1.15E-05	5.28E-03	22.440
salic xx	4.42E-05	0.094	6.50E-07	2.90E-02	2400	0.11	9.48E-03	3.95E-06	6.89E-06	1.53E-05	1.54E-05	3.89E-03	22.752
salic xx	3.13E-05	0.046	2.14E-07	1.40E-02	2400	0.05	6.46E-03	2.69E-06	4.69E-06	1.57E-05	1.60E-05	1.94E-03	15.504
salic xx	2.21E-05	0.019	6.80E-08	5.85E-03	2400	0.02	4.93E-03	2.05E-06	3.58E-06	1.65E-05	1.48E-05	7.51E-04	11.832
salic xx	1.56E-05	0.007	1.38E-08	2.22E-03	2400	0.01	2.64E-03	1.10E-06	1.92E-06	1.06E-05	8.51E-06	1.63E-04	6.336
salic xx	1.10E-05	0.003	1.52E-09	8.00E-04	2400	0.00	8.06E-04	3.36E-07	5.85E-07	3.53E-06	2.49E-06	1.73E-05	1.934
dense lithic	1.60E-02	13.785	3.13E-05	4.24E+00	2570	15.37	3.35E-03	1.30E-06	2.27E-06	5.30E-08	1.15E-07	4.54E-03	8.610
dense lithic	1.13E-02	11.191	1.01E-04	3.44E+00	2570	12.48	1.33E-02	5.18E-06	9.02E-06	2.59E-07	2.23E-07	7.16E-03	34.181
dense lithic	8.00E-03	9.061	3.57E-05	2.79E+00	2570	10.11	5.81E-03	2.26E-06	3.94E-06	1.40E-07	2.62E-07	6.80E-03	14.932
dense lithic	5.66E-03	7.318	9.13E-05	2.25E+00	2570	8.16	1.84E-02	7.16E-06	1.25E-05	5.47E-07	6.96E-07	1.46E-02	47.288
dense lithic	4.00E-03	5.897	1.13E-04	1.81E+00	2570	6.58	2.83E-02	1.10E-05	1.92E-05	1.04E-06	1.08E-06	1.82E-02	72.731
dense lithic	2.83E-03	4.741	8.20E-05	1.46E+00	2570	5.29	2.55E-02	9.92E-06	1.73E-05	1.17E-06	1.38E-06	1.87E-02	65.535
dense lithic	2.00E-03	3.805	8.52E-05	1.17E+00	2570	4.24	3.30E-02	1.28E-05	2.24E-05	1.87E-06	2.03E-06	2.22E-02	84.810
dense lithic	1.41E-03	3.045	7.00E-05	9.37E-01	2570	3.40	3.39E-02	1.32E-05	2.30E-05	2.40E-06	2.48E-06	2.17E-02	87.123
dense lithic	1.00E-03	2.424	4.98E-05	7.46E-01	2570	2.70	3.03E-02	1.18E-05	2.06E-05	2.68E-06	3.08E-06	2.14E-02	77.871
dense lithic	7.07E-04	1.912	4.66E-05	5.88E-01	2570	2.13	3.59E-02	1.40E-05	2.44E-05	4.01E-06	3.88E-06	2.13E-02	92.263

dense lithic	5.00E-04	1.488	2.54E-05	4.58E-01	2570	1.66	2.52E-02	9.81E-06	1.71E-05	3.59E-06	3.75E-06	1.60E-02	64.764
dense lithic	3.54E-04	1.135	1.72E-05	3.49E-01	2570	1.27	2.23E-02	8.68E-06	1.51E-05	4.12E-06	4.11E-06	1.34E-02	57.311
dense lithic	2.50E-04	0.846	9.59E-06	2.60E-01	2570	0.94	1.67E-02	6.50E-06	1.13E-05	4.08E-06	4.13E-06	1.00E-02	42.919
dense lithic	1.77E-04	0.613	5.37E-06	1.89E-01	2570	0.68	1.29E-02	5.02E-06	8.75E-06	4.25E-06	4.27E-06	7.50E-03	33.153
dense lithic	1.25E-04	0.428	2.72E-06	1.32E-01	2570	0.48	9.38E-03	3.65E-06	6.36E-06	4.29E-06	4.28E-06	5.24E-03	24.107
dense lithic	8.84E-05	0.283	1.25E-06	8.72E-02	2570	0.32	6.48E-03	2.52E-06	4.40E-06	4.25E-06	4.73E-06	3.84E-03	16.654
dense lithic	6.25E-05	0.175	7.17E-07	5.38E-02	2570	0.19	6.05E-03	2.35E-06	4.10E-06	5.87E-06	6.83E-06	3.42E-03	15.549
dense lithic	4.42E-05	0.097	4.06E-07	3.00E-02	2570	0.11	6.14E-03	2.39E-06	4.17E-06	9.06E-06	9.14E-06	2.55E-03	15.780
dense lithic	3.13E-05	0.047	1.34E-07	1.46E-02	2570	0.05	4.18E-03	1.63E-06	2.84E-06	9.32E-06	9.50E-06	1.29E-03	10.743
dense lithic	2.21E-05	0.020	4.28E-08	6.06E-03	2570	0.02	3.20E-03	1.25E-06	2.17E-06	9.92E-06	8.87E-06	5.01E-04	8.224
dense lithic	1.56E-05	0.007	8.58E-09	2.28E-03	2570	0.01	1.71E-03	6.65E-07	1.16E-06	6.42E-06	5.13E-06	1.09E-04	4.395
dense lithic	1.10E-05	0.003	9.51E-10	8.31E-04	2570	0.00	5.19E-04	2.02E-07	3.52E-07	2.12E-06	1.49E-06	1.15E-05	1.334
femic xx	2.00E-03	4.962	7.38E-07	1.53E+00	3280	5.53	2.80E-04	8.54E-08	1.49E-07	9.59E-09	1.65E-08	3.00E-04	0.918
femic xx	1.41E-03	4.005	1.64E-06	1.23E+00	3280	4.47	7.72E-04	2.35E-07	4.10E-07	3.27E-08	5.76E-08	8.45E-04	2.532
femic xx	1.00E-03	3.207	3.75E-06	9.86E-01	3280	3.58	2.20E-03	6.71E-07	1.17E-06	1.16E-07	2.21E-07	2.60E-03	7.216
femic xx	7.07E-04	2.532	9.46E-06	7.79E-01	3280	2.82	7.03E-03	2.14E-06	3.74E-06	4.67E-07	7.21E-07	6.68E-03	23.058
femic xx	5.00E-04	1.962	1.61E-05	6.03E-01	3280	2.19	1.54E-02	4.70E-06	8.19E-06	1.31E-06	1.53E-06	1.09E-02	50.512
femic xx	3.54E-04	1.487	1.43E-05	4.57E-01	3280	1.66	1.81E-02	5.52E-06	9.62E-06	2.02E-06	2.00E-06	1.09E-02	59.368
femic xx	2.50E-04	1.100	7.72E-06	3.39E-01	3280	1.23	1.32E-02	4.02E-06	7.02E-06	1.97E-06	1.77E-06	7.12E-03	43.296
femic xx	1.77E-04	0.792	2.70E-06	2.44E-01	3280	0.88	6.42E-03	1.96E-06	3.41E-06	1.31E-06	1.03E-06	2.99E-03	21.058
femic xx	1.25E-04	0.549	4.03E-07	1.69E-01	3280	0.61	1.38E-03	4.21E-07	7.33E-07	3.95E-07	3.96E-07	7.94E-04	4.526
femic xx	8.84E-05	0.363	1.84E-07	1.12E-01	3280	0.40	9.57E-04	2.92E-07	5.09E-07	3.97E-07	4.44E-07	5.89E-04	3.139
femic xx	6.25E-05	0.224	1.06E-07	6.90E-02	3280	0.25	8.88E-04	2.71E-07	4.72E-07	5.54E-07	6.48E-07	5.32E-04	2.913
femic xx	4.42E-05	0.127	6.09E-08	3.91E-02	3280	0.14	9.02E-04	2.75E-07	4.79E-07	8.69E-07	8.84E-07	4.11E-04	2.959
femic xx	3.13E-05	0.064	2.09E-08	1.97E-02	3280	0.07	6.15E-04	1.88E-07	3.27E-07	9.18E-07	9.51E-07	2.23E-04	2.017
femic xx	2.21E-05	0.028	7.00E-09	8.58E-03	3280	0.03	4.72E-04	1.44E-07	2.51E-07	1.03E-06	9.30E-07	9.48E-05	1.548
femic xx	1.56E-05	0.011	1.45E-09	3.32E-03	3280	0.01	2.53E-04	7.71E-08	1.34E-07	7.03E-07	5.62E-07	2.21E-05	0.830
femic xx	1.10E-05	0.004	1.56E-10	1.20E-03	3280	0.00	7.52E-05	2.29E-08	4.00E-08	2.35E-07	1.65E-07	2.38E-06	0.247

## Summary of results

0.593  $S_r$   
( $\text{kgm}^{-2}\text{s}^{-1}$ )  
1248.9  $\rho_{dep}$   
( $\text{kgm}^{-3}$ )  
4.75E-04  $A_r$

	(ms <sup>-1</sup> )
1.5	d (m)
3159	t (s)

P <sub>navg</sub>	1.950
P <sub>n*</sub>	3.250

Supporting file E. Step-by-step calculation of the sedimentation rate *S*, and deposition time *t* of the tractional layer with sand-wave bedforms

cropping out at Sbirri. In yellow the main parametres. For symbols see the symbol list.  $\gamma_i = \frac{c_{trans_i}}{\left(\left(10.065 * P_{n_i}/P_{n*}\right) + 0.1579\right)}$ ;  $\gamma_{true_i} = \gamma_i * 0.7 + \gamma_{i+1} * 0.3$ .

		Component	$D_i$ (m)	$D_i/D_{sm}$	$\tau_{ri}$ (Pa)	$\tau/\tau_{ri}$	$W_i^*$	$Q_{bi}$ (m <sup>2</sup> s <sup>-1</sup> )	$\phi_i$	$\rho_{si}$ (kgm <sup>-3</sup> )
$u_*$ (ms <sup>-1</sup> )	0.43									
$\rho_s$ average (kgm <sup>-3</sup> )	2440	vesicular glass	5.66E-03	77.534	3.517	0.098	5.645E-08	4.665E-16	0.0013	596
$C_{tot}$ (-)	0.0006	vesicular glass	4.00E-03	54.795	2.787	0.123	1.608E-07	2.657E-15	0.0027	669
$\rho_{mix}$ (kgm <sup>-3</sup> )	1.835	vesicular glass	2.83E-03	38.767	2.210	0.155	4.563E-07	6.034E-15	0.0021	751
$\tau$ (Pa)	0.343	vesicular glass	2.00E-03	27.397	1.752	0.196	1.300E-06	1.718E-14	0.0021	842
$D_{sm}$ (mm)	0.07	vesicular glass	1.41E-03	19.315	1.386	0.247	3.728E-06	1.048E-13	0.0045	945
$Re_*$ (-)	2	vesicular glass	1.00E-03	13.699	1.101	0.311	1.051E-05	2.952E-13	0.0045	1060
$\theta_i$ (-)	0.017	vesicular glass	7.07E-04	9.685	0.872	0.393	2.994E-05	1.237E-12	0.0066	1189
$\tau_{rsm}^*$ (-)	0.0128	vesicular glass	5.00E-04	6.849	0.688	0.499	8.729E-05	4.977E-12	0.0091	1334
$\tau_{rsm}$ (Pa)	0.1906	vesicular glass	3.54E-04	4.849	0.530	0.647	2.827E-04	2.686E-11	0.0152	1497
$S_r$ (kgm <sup>-2</sup> s <sup>-1</sup> )	0.096	vesicular glass	2.50E-04	3.425	0.392	0.876	1.101E-03	1.110E-10	0.0162	1679
$\rho_{dep}$ (kgm <sup>-3</sup> )	1440	vesicular glass	1.77E-04	2.425	0.292	1.176	4.151E-03	5.797E-10	0.0224	1884
$A_r$ (ms <sup>-1</sup> )	6.667E-05	vesicular glass	1.25E-04	1.712	0.233	1.474	3.471E-02	5.192E-09	0.0240	2114
		vesicular glass	8.84E-05	1.211	0.201	1.703	7.723E-02	1.692E-08	0.0351	2372
		vesicular glass	6.25E-05	0.856	0.184	1.864	1.169E-01	3.060E-08	0.0420	2560
		vesicular glass	4.42E-05	0.605	0.173	1.983	1.504E-01	4.321E-08	0.0461	2560
		vesicular glass	3.13E-05	0.429	0.165	2.079	1.799E-01	5.238E-08	0.0467	2560
		vesicular glass	2.21E-05	0.303	0.158	2.166	2.083E-01	4.834E-08	0.0372	2560
		vesicular glass	1.56E-05	0.214	0.152	2.250	2.371E-01	3.003E-08	0.0203	2560
		vesicular glass	1.10E-05	0.151	0.147	2.336	2.676E-01	9.121E-09	0.0055	2560
		salic xx	1.41E-03	19.315	1.386	0.247	3.728E-06	2.465E-14	0.0011	2400
		salic xx	1.00E-03	13.699	1.101	0.311	1.051E-05	8.682E-14	0.0013	2400
		salic xx	7.07E-04	9.685	0.872	0.393	2.994E-05	9.896E-13	0.0053	2400
		salic xx	5.00E-04	6.849	0.688	0.499	8.729E-05	5.698E-12	0.0105	2400
		salic xx	3.54E-04	4.849	0.530	0.647	2.827E-04	2.943E-11	0.0167	2400
		salic xx	2.50E-04	3.425	0.392	0.876	1.101E-03	1.320E-10	0.0192	2400
		salic xx	1.77E-04	2.425	0.292	1.176	4.151E-03	5.626E-10	0.0217	2400
		salic xx	1.25E-04	1.712	0.233	1.474	3.471E-02	4.934E-09	0.0228	2400
		salic xx	8.84E-05	1.211	0.201	1.703	7.723E-02	1.608E-08	0.0334	2400

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salic xx	6.25E-05	0.856	0.184	1.864	1.169E-01	2.908E-08	0.0399	2400
salic xx	4.42E-05	0.605	0.173	1.983	1.504E-01	4.106E-08	0.0438	2400
salic xx	3.13E-05	0.429	0.165	2.079	1.799E-01	4.977E-08	0.0444	2400
salic xx	2.21E-05	0.303	0.158	2.166	2.083E-01	4.594E-08	0.0354	2400
salic xx	1.56E-05	0.214	0.152	2.250	2.371E-01	2.853E-08	0.0193	2400
salic xx	1.10E-05	0.151	0.147	2.336	2.676E-01	8.667E-09	0.0052	2400
dense lithic	2.83E-03	38.767	2.210	0.155	4.563E-07	5.279E-15	0.0019	2588
dense lithic	2.00E-03	27.397	1.752	0.196	1.300E-06	2.148E-14	0.0027	2588
dense lithic	1.41E-03	19.315	1.386	0.247	3.728E-06	8.011E-14	0.0034	2588
dense lithic	1.00E-03	13.699	1.101	0.311	1.051E-05	2.604E-13	0.0040	2588
dense lithic	7.07E-04	9.685	0.872	0.393	2.994E-05	2.919E-12	0.0156	2588
dense lithic	5.00E-04	6.849	0.688	0.499	8.729E-05	9.305E-12	0.0171	2588
dense lithic	3.54E-04	4.849	0.530	0.647	2.827E-04	3.971E-11	0.0225	2588
dense lithic	2.50E-04	3.425	0.392	0.876	1.101E-03	1.374E-10	0.0200	2588
dense lithic	1.77E-04	2.425	0.292	1.176	4.151E-03	5.797E-10	0.0224	2588
dense lithic	1.25E-04	1.712	0.233	1.474	3.471E-02	4.590E-09	0.0212	2588
dense lithic	8.84E-05	1.211	0.201	1.703	7.723E-02	1.495E-08	0.0310	2588
dense lithic	6.25E-05	0.856	0.184	1.864	1.169E-01	2.705E-08	0.0371	2588
dense lithic	4.42E-05	0.605	0.173	1.983	1.504E-01	3.819E-08	0.0407	2588
dense lithic	3.13E-05	0.429	0.165	2.079	1.799E-01	4.630E-08	0.0413	2588
dense lithic	2.21E-05	0.303	0.158	2.166	2.083E-01	4.273E-08	0.0329	2588
dense lithic	1.56E-05	0.214	0.152	2.250	2.371E-01	2.654E-08	0.0180	2588
dense lithic	1.10E-05	0.151	0.147	2.336	2.676E-01	8.062E-09	0.0048	2588
					$Q_b$ (m <sup>2</sup> s <sup>-1</sup> )	6.665E-07	1	

Supporting file F. Step-by-step calculation of the bedload transportation rate  $Q_b$ , Sedimentation rate  $S_r$  and aggradation rate  $A_r$  of the fine-ash structureless layer at Camaldoli. In yellow the main parametres. For symbols see the symbol list.

Input data from Pyflow		Component	$D_i$ (m)	$P_{ni}$	$P_{ni} * C_{transi}$	$P_{ni}/P_n^*$	$\rho_{si}$ (kgm <sup>-3</sup> )	$w_i$ (ms <sup>-1</sup> )	$\phi_i$	Vol. fraction (m <sup>3</sup> kg <sup>-1</sup> )	$C_{transi}$	$\gamma_i$	$\gamma_{truei}$	$S_{ri}$ (kgm <sup>-2</sup> s <sup>-1</sup> )	$\rho_{si} * \phi_i$ (kgm <sup>-3</sup> )
		vesicular glass	2.83E-03	4.484	9.51E-05	2.14	1508	3.015	7.62E-03	5.05E-06	2.12E-05	9.68E-07	1.58E-06	7.16E-03	11.49
		vesicular glass	2.00E-03	3.831	1.45E-04	1.83	1615	2.576	1.46E-02	9.04E-06	3.79E-05	2.02E-06	2.65E-06	1.10E-02	23.58
$u^*$ (ms <sup>-1</sup> )	2.00	vesicular glass	1.41E-03	3.194	2.05E-04	1.52	1723	2.148	2.64E-02	1.53E-05	6.43E-05	4.11E-06	5.14E-06	1.90E-02	45.48
$C_{tot}$	0.00175	vesicular glass	1.00E-03	2.613	2.53E-04	1.25	1830	1.757	4.22E-02	2.31E-05	9.68E-05	7.54E-06	7.99E-06	2.57E-02	77.23
$P_{nsusp}$	0.861	vesicular glass	7.07E-04	2.102	1.97E-04	1.00	1937	1.414	4.32E-02	2.23E-05	9.36E-05	9.04E-06	1.00E-05	2.74E-02	83.70
		vesicular glass	5.00E-04	1.664	1.69E-04	0.79	2045	1.119	4.94E-02	2.42E-05	1.01E-04	1.23E-05	1.26E-05	2.89E-02	101.02
		vesicular glass	3.54E-04	1.293	1.11E-04	0.62	2152	0.870	4.39E-02	2.04E-05	8.56E-05	1.33E-05	1.29E-05	2.41E-02	94.49
		vesicular glass	2.50E-04	1.292	9.88E-05	0.62	2260	0.869	4.12E-02	1.82E-05	7.65E-05	1.19E-05	1.19E-05	2.34E-02	93.10
		vesicular glass	1.77E-04	0.996	5.91E-05	0.48	2367	0.670	3.35E-02	1.42E-05	5.94E-05	1.19E-05	1.27E-05	2.01E-02	79.30
		vesicular glass	1.25E-04	0.677	3.39E-05	0.32	2475	0.456	2.95E-02	1.19E-05	5.00E-05	1.45E-05	1.67E-05	1.88E-02	73.00
		vesicular glass	8.84E-05	0.433	2.13E-05	0.21	2582	0.291	3.03E-02	1.17E-05	4.92E-05	2.18E-05	2.09E-05	1.57E-02	78.24
		vesicular glass	6.25E-05	0.257	6.73E-06	0.12	2690	0.173	1.68E-02	6.25E-06	2.62E-05	1.87E-05	2.05E-05	9.55E-03	45.18
		vesicular glass	4.42E-05	0.140	2.92E-06	0.07	2797	0.094	1.39E-02	4.97E-06	2.09E-05	2.49E-05	2.55E-05	6.73E-03	38.88
		vesicular glass	3.13E-05	0.066	8.55E-07	0.03	2760	0.045	8.51E-03	3.08E-06	1.29E-05	2.70E-05	2.87E-05	3.53E-03	23.49
		vesicular glass	2.21E-05	0.028	2.74E-07	0.01	2760	0.019	6.35E-03	2.30E-06	9.65E-06	3.27E-05	3.11E-05	1.64E-03	17.53
		vesicular glass	1.56E-05	0.011	6.03E-08	0.01	2760	0.007	3.78E-03	1.37E-06	5.75E-06	2.75E-05	2.30E-05	4.50E-04	10.43
		vesicular glass	1.10E-05	0.003	6.39E-09	0.00	2760	0.002	1.40E-03	5.07E-07	2.13E-06	1.23E-05	8.64E-06	4.77E-05	3.86
		femic xx	2.00E-03	6.344	1.10E-05	3.03	3280	4.267	1.35E-03	4.12E-07	1.73E-06	5.58E-08	1.16E-07	1.62E-03	4.43
		femic xx	1.41E-03	5.206	2.64E-05	2.48	3280	3.501	3.96E-03	1.21E-06	5.07E-06	1.99E-07	2.20E-07	2.53E-03	12.99
		femic xx	1.00E-03	4.166	2.29E-05	1.99	3280	2.802	4.29E-03	1.31E-06	5.49E-06	2.70E-07	3.57E-07	3.28E-03	14.07
		femic xx	7.07E-04	3.270	2.93E-05	1.56	3280	2.199	7.01E-03	2.14E-06	8.97E-06	5.60E-07	7.04E-07	5.08E-03	22.99
		femic xx	5.00E-04	2.527	3.26E-05	1.21	3280	1.699	1.01E-02	3.08E-06	1.29E-05	1.04E-06	1.18E-06	6.57E-03	33.13
		femic xx	3.54E-04	1.920	2.73E-05	0.92	3280	1.291	1.11E-02	3.38E-06	1.42E-05	1.50E-06	1.69E-06	7.17E-03	36.41
		femic xx	2.50E-04	1.428	2.17E-05	0.68	3280	0.961	1.19E-02	3.63E-06	1.52E-05	2.15E-06	2.16E-06	6.80E-03	39.03
		femic xx	1.77E-04	1.322	1.89E-05	0.63	3280	0.889	1.12E-02	3.41E-06	1.43E-05	2.18E-06	2.09E-06	6.11E-03	36.74
		femic xx	1.25E-04	0.898	7.67E-06	0.43	3280	0.604	6.68E-03	2.04E-06	8.55E-06	1.89E-06	2.22E-06	4.39E-03	21.91
		femic xx	8.84E-05	0.577	5.06E-06	0.28	3280	0.388	6.86E-03	2.09E-06	8.78E-06	2.97E-06	2.87E-06	3.66E-03	22.50
		femic xx	6.25E-05	0.347	1.69E-06	0.17	3280	0.233	3.81E-03	1.16E-06	4.87E-06	2.65E-06	2.96E-06	2.26E-03	12.50

Sedimentology

femic xx	4.42E-05	0.193	7.78E-07	0.09	3280	0.130	3.15E-03	9.60E-07	4.03E-06	3.68E-06	3.73E-06	1.59E-03	10.33
femic xx	3.13E-05	0.100	2.46E-07	0.05	3280	0.067	1.93E-03	5.88E-07	2.47E-06	3.85E-06	4.12E-06	9.06E-04	6.33
femic xx	2.21E-05	0.048	8.89E-08	0.02	3280	0.032	1.45E-03	4.42E-07	1.86E-06	4.75E-06	4.60E-06	4.86E-04	4.76
femic xx	1.56E-05	0.021	2.34E-08	0.01	3280	0.014	8.64E-04	2.63E-07	1.11E-06	4.24E-06	3.56E-06	1.67E-04	2.83
femic xx	1.10E-05	0.008	3.28E-09	0.00	3280	0.006	3.05E-04	9.30E-08	3.90E-07	1.96E-06	1.38E-06	2.57E-05	1.00
dense lithic	5.66E-03	6.527	5.41E-05	3.11	2570	4.390	5.08E-03	1.98E-06	8.29E-06	2.61E-07	3.45E-07	3.90E-03	13.06
dense lithic	4.00E-03	5.730	4.52E-05	2.73	2570	3.853	4.83E-03	1.88E-06	7.89E-06	2.82E-07	3.90E-07	3.86E-03	12.41
dense lithic	2.83E-03	4.904	7.53E-05	2.34	2570	3.298	9.40E-03	3.66E-06	1.53E-05	6.41E-07	1.11E-06	9.45E-03	24.16
dense lithic	2.00E-03	4.078	1.80E-04	1.95	2570	2.743	2.71E-02	1.05E-05	4.42E-05	2.22E-06	3.01E-06	2.12E-02	69.65
dense lithic	1.41E-03	3.307	2.61E-04	1.58	2570	2.224	4.83E-02	1.88E-05	7.89E-05	4.87E-06	5.60E-06	3.20E-02	124.13
dense lithic	1.00E-03	2.630	2.48E-04	1.26	2570	1.769	5.78E-02	2.25E-05	9.44E-05	7.31E-06	7.82E-06	3.55E-02	148.55
dense lithic	7.07E-04	2.059	1.88E-04	0.98	2570	1.385	5.60E-02	2.18E-05	9.14E-05	9.01E-06	9.57E-06	3.41E-02	143.92
dense lithic	5.00E-04	1.587	1.35E-04	0.76	2570	1.067	5.23E-02	2.04E-05	8.54E-05	1.09E-05	1.08E-05	2.96E-02	134.41
dense lithic	3.54E-04	1.200	7.62E-05	0.57	2570	0.807	3.89E-02	1.51E-05	6.35E-05	1.06E-05	9.98E-06	2.07E-02	99.97
dense lithic	2.50E-04	1.277	6.86E-05	0.61	2570	0.859	3.29E-02	1.28E-05	5.37E-05	8.46E-06	8.43E-06	1.86E-02	84.55
dense lithic	1.77E-04	0.974	3.97E-05	0.46	2570	0.655	2.50E-02	9.73E-06	4.08E-05	8.36E-06	9.09E-06	1.53E-02	64.25
dense lithic	1.25E-04	0.628	2.17E-05	0.30	2570	0.423	2.12E-02	8.25E-06	3.46E-05	1.08E-05	1.30E-05	1.41E-02	54.48
dense lithic	8.84E-05	0.375	1.34E-05	0.18	2570	0.253	2.18E-02	8.48E-06	3.56E-05	1.80E-05	1.78E-05	1.15E-02	56.03
dense lithic	6.25E-05	0.204	4.02E-06	0.10	2570	0.137	1.21E-02	4.71E-06	1.98E-05	1.72E-05	1.98E-05	6.96E-03	31.10
dense lithic	4.42E-05	0.098	1.61E-06	0.05	2570	0.066	1.00E-02	3.89E-06	1.63E-05	2.57E-05	2.64E-05	4.49E-03	25.70
dense lithic	3.13E-05	0.041	4.13E-07	0.02	2570	0.028	6.12E-03	2.38E-06	9.99E-06	2.79E-05	2.94E-05	2.10E-03	15.73
dense lithic	2.21E-05	0.014	1.06E-07	0.01	2570	0.010	4.57E-03	1.78E-06	7.46E-06	3.29E-05	3.07E-05	7.50E-04	11.74
dense lithic	1.56E-05	0.003	1.38E-08	0.00	2570	0.002	2.72E-03	1.06E-06	4.44E-06	2.57E-05	2.02E-05	1.09E-04	6.99
dense lithic	1.10E-05	0.013	2.06E-08	0.01	2570	0.009	9.91E-04	3.86E-07	1.62E-06	7.37E-06	5.16E-06	1.13E-04	2.55
salic xx	2.00E-03	5.666	2.01E-06	2.70	2400	3.811	2.03E-04	8.46E-08	3.55E-07	1.28E-08	4.42E-08	4.05E-04	0.49
salic xx	1.41E-03	4.616	1.09E-05	2.20	2400	3.104	1.35E-03	5.63E-07	2.36E-06	1.05E-07	1.35E-07	1.01E-03	3.24
salic xx	1.00E-03	3.672	1.37E-05	1.75	2400	2.469	2.13E-03	8.88E-07	3.72E-06	2.07E-07	2.77E-07	1.64E-03	5.11
salic xx	7.07E-04	2.868	1.77E-05	1.37	2400	1.929	3.53E-03	1.47E-06	6.17E-06	4.39E-07	5.27E-07	2.44E-03	8.47
salic xx	5.00E-04	2.205	1.75E-05	1.05	2400	1.483	4.55E-03	1.90E-06	7.96E-06	7.33E-07	1.01E-06	3.61E-03	10.92
salic xx	3.54E-04	1.665	2.29E-05	0.79	2400	1.120	7.87E-03	3.28E-06	1.38E-05	1.67E-06	1.83E-06	4.93E-03	18.89

Summary of results	
0.589	$S_r$ ( $\text{kgm}^{-2}\text{s}^{-1}$ )
1476	$\rho_{dep}$ ( $\text{kgm}^{-3}$ )
0.00040	$A_r$ ( $\text{ms}^{-1}$ )
0.15	d (m)
375.7	t (s)

salic xx	2.50E-04	1.277	1.79E-05	0.61	2400	0.859	8.03E-03	3.35E-06	1.40E-05	2.21E-06	2.46E-06	5.06E-03	19.27
salic xx	1.77E-04	1.087	1.79E-05	0.52	2400	0.731	9.40E-03	3.92E-06	1.64E-05	3.03E-06	3.16E-06	5.54E-03	22.56
salic xx	1.25E-04	0.731	9.38E-06	0.35	2400	0.492	7.34E-03	3.06E-06	1.28E-05	3.46E-06	4.07E-06	4.80E-03	17.62
salic xx	8.84E-05	0.465	6.13E-06	0.22	2400	0.312	7.54E-03	3.14E-06	1.32E-05	5.47E-06	5.29E-06	3.97E-03	18.10
salic xx	6.25E-05	0.276	2.01E-06	0.13	2400	0.186	4.17E-03	1.74E-06	7.29E-06	4.87E-06	5.43E-06	2.42E-03	10.01
salic xx	4.42E-05	0.153	9.21E-07	0.07	2400	0.103	3.45E-03	1.44E-06	6.03E-06	6.72E-06	6.75E-06	1.66E-03	8.28
salic xx	3.13E-05	0.079	2.91E-07	0.04	2400	0.053	2.11E-03	8.79E-07	3.69E-06	6.83E-06	7.18E-06	9.14E-04	5.06
salic xx	2.21E-05	0.038	1.05E-07	0.02	2400	0.026	1.57E-03	6.54E-07	2.75E-06	8.01E-06	7.65E-06	4.70E-04	3.77
salic xx	1.56E-05	0.017	2.83E-08	0.01	2400	0.012	9.40E-04	3.92E-07	1.64E-06	6.81E-06	5.74E-06	1.58E-04	2.26
salic xx	1.10E-05	0.007	4.42E-09	0.00	2400	0.005	3.56E-04	1.48E-07	6.22E-07	3.24E-06	2.27E-06	2.56E-05	0.85

$$P_{\text{navg}} = 1.80\text{E}+00$$

$$P_n^* = 2.10\text{E}+00$$

Supporting file G. Step-by-step calculation of the sedimentation rate  $S_r$  of the tractive layer cropping under the massive bed at the “split location”. In yellow the main parametres. For symbols see the symbol list.  $\gamma_i = \frac{c_{trans_i}}{\left(\left(10.065 * \frac{P_{n_i}}{P_n^*}\right) + 0.1579\right)}$ ;  $\gamma_{true_i} = \gamma_i * 0.7 + \gamma_{i+1} * 0.3$ .