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# Search for Higgs boson decays into Z and J/ψ and for Higgs and Z boson decays into J/ψ or Y pairs in pp collisions at $\sqrt{s} = 13$ TeV

The CMS Collaboration\*

CERN, Geneva, Switzerland

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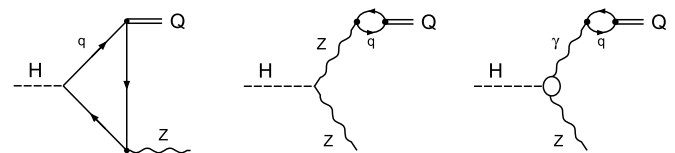
## ABSTRACT

Decays of the Higgs boson into a Z boson and a J/ψ or ψ(2S) meson are searched for in four-lepton final states with the CMS detector at the LHC. A data set of proton-proton collisions corresponding to an integrated luminosity of 138 fb<sup>-1</sup> is used. Using the same data set, decays of the Higgs and Z boson into quarkonium pairs are also searched for. An observation of such decays with this sample would indicate the presence of physics beyond the standard model. No evidence for these decays has been observed and upper limits at the 95% confidence level are placed on the corresponding branching fractions ( $\mathcal{B}$ ). Assuming longitudinal polarization of the Higgs boson decay products, 95% confidence level observed upper limits for  $\mathcal{B}(H \rightarrow ZJ/\psi)$  and  $\mathcal{B}(H \rightarrow Z\psi(2S))$  are  $1.9 \times 10^{-3}$  and  $6.6 \times 10^{-3}$ , respectively.

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## 1. Introduction

A boson with a mass of about 125 GeV was discovered by the ATLAS and CMS Collaborations at the CERN LHC in 2012 [1–3]. Comprehensive studies in various decay channels and production modes followed, and combined measurements show that the properties of the new boson are, so far, consistent with the standard model (SM) predictions for the Higgs boson (H) [4–6]. Rare exclusive decays of the H to mesons provide experimentally clean final states to study deviations of Yukawa couplings from SM predictions that cannot be obtained with inclusive measurements. Several models beyond the SM predict enhanced Yukawa couplings to fermions [7], leading to branching fractions ( $\mathcal{B}$ ) that are enhanced by up to three orders of magnitude. Examples include the Giudice–Lebedev model of quark masses [8], the two Higgs doublet model [9], the single Higgs doublet model with the Froggatt–Nielsen mechanism [10], and Randall–Sundrum models [11,12]. The required sensitivity for observing Yukawa couplings to the second- and first-generation fermions has not yet been reached, although recently the CMS Collaboration published evidence for the Higgs boson decay to a pair of muons [13,14]. The observed upper limit for  $\mathcal{B}(H \rightarrow c\bar{c})$  times the cross section for H production in association with vector bosons as measured by the ATLAS and CMS experiments is found to be 26 and 14 times the SM expectation [15,16], respectively.



**Fig. 1.** Sample Feynman diagrams depicting direct (left) and indirect (middle, right) quark coupling contributions to the  $H \rightarrow ZQ$  decay, where Q represents a quarkonium state and q is a quark, which in the leftmost diagram is either of charm or bottom flavor. The diagrams represent Higgs boson decays into quarkonium pairs when replacing the bottom section with the upper half in each.

The first class of processes that is considered here is the decay of the Higgs boson into a Z boson and a vector meson quarkonium state (Q) [7,17,18]. The relevant SM Feynman diagrams for the decays  $H \rightarrow ZQ$  are shown in Fig. 1. The first diagram in Fig. 1 represents contributing amplitudes at leading order (LO), where the Higgs boson directly couples to a quark-antiquark pair that radiates a Z boson and forms the meson. The last two diagrams depict indirect contributions to the decay amplitude. The last diagram corresponds to one-loop amplitudes as indicated by the blob, as well as tree-level effective vertices [17]. The branching fractions of these processes in the SM are  $\mathcal{B}(H \rightarrow ZJ/\psi) = 2.3 \times 10^{-6}$  [17–19] and  $\mathcal{B}(H \rightarrow Z\psi(2S)) = 1.7 \times 10^{-6}$  [19]. The main source of background events in these final states is from production of a Z boson in association with a genuine [20] or misidentified meson candidate. New physics could affect the direct boson couplings or could enter through loops and alter the interference pattern between the am-

\* E-mail address: [cms-publication-committee-chair@cern.ch](mailto:cms-publication-committee-chair@cern.ch).

plitudes. Any of those possibilities can enhance branching fractions with respect to the SM predictions. For the rare decays  $H \rightarrow Z/\psi$  and  $H \rightarrow Z\psi(2S)$ , the  $H \rightarrow Z\gamma^*$  amplitude contributes significantly. Hence, these rare decays provide complementary information to the decay  $H \rightarrow Z\gamma$ , both in and beyond the SM [17,19]. The decays of the Higgs boson into  $Z/\psi$  and  $Z\eta_c$  have been searched for by the ATLAS Collaboration in hadronic final states, reaching 95% confidence level (CL) upper limits on the branching fraction of the Higgs boson that exceed 100% [21]. Recently, the CMS Collaboration published upper limits on the branching fraction for  $H \rightarrow Z\rho$  and  $H \rightarrow Z\phi$  at 95% CL that are larger than the expected SM branching fraction by more than a factor of 700 [22].

A second related class of processes is the Higgs boson decay into pairs of quarkonia. The Feynman diagrams are variants of the ones depicted in Fig. 1: in each diagram, the on-shell Z boson in the lower part is replaced by a quarkonium decay, similar to the process depicted in the upper part. In the rightmost diagram, both vector bosons could be also gluons, in which case additional soft-gluon exchange occurs. The importance of the measurement of such decays has been pointed out in Refs. [23–26]. Using a phenomenological approach for the direct  $H-q\bar{q}$  coupling, Ref. [23] finds that the dominant quarkonium pair decay mode is  $H \rightarrow YY$  with an estimated branching fraction of  $O(10^{-5})$ . More recently, Ref. [27] predicts values of  $\mathcal{B}(H \rightarrow J/\psi J/\psi) = 1.5 \times 10^{-10}$  and  $\mathcal{B}(H \rightarrow YY) = 2 \times 10^{-9}$  assuming the dominance of indirect amplitudes. Inclusion of the mechanism where the Higgs boson couples directly to charm or bottom quarks, which then hadronize to heavy quarkonia, in the calculation in Ref. [28] leads to an increase by an order of magnitude in the related  $\mathcal{B}(H \rightarrow J/\psi\gamma)$ . The Higgs boson is expected to couple to quarkonium pairs that include radially excited states with comparable strength [29]. Recently, a first search for the decays  $H \rightarrow J/\psi J/\psi$  and  $H \rightarrow Y(nS)Y(mS)$  ( $n, m = 1, 2, 3$ ) was performed by the CMS Collaboration [30].

Related to these two classes of processes is the decay of the Higgs boson into a photon and a vector meson [17,28,31]. The 95% CL upper limits on the branching fractions of the Higgs boson into  $\gamma J/\psi$ ,  $\gamma\rho$ , and  $\gamma\phi$  are found to be two orders of magnitude larger than their expected values in the SM [32–34]. For the  $\gamma\psi(2S)$  and  $\gamma Y(nS)$  decays, the corresponding upper limits are, respectively, three and five orders of magnitude larger than the expected SM branching fractions.

This Letter presents the first search for decays of the 125 GeV Higgs boson into a Z boson and a  $J/\psi$  or  $\psi(2S)$  meson in four-lepton final states. The Z boson is reconstructed from its decays into  $\mu^+\mu^-$  or  $e^+e^-$ , the  $J/\psi$  meson from its decay into  $\mu^+\mu^-$ , and the  $\psi(2S)$  from its inclusive decay into  $J/\psi X$  (feed-down), where X, which is mostly  $\pi\pi$ , is not reconstructed. Furthermore, an update of Higgs boson searches in  $J/\psi J/\psi$  and  $Y(nS)Y(mS)$  decay channels with the full available data sample is presented. New channels are accessed via the inclusive decay of  $\psi(2S)$  into a  $J/\psi$  meson. For the  $Y(nS)$  states the possibilities that they are the result of feed-down transitions from higher Y states before decaying into muon pairs are included. Finally, this Letter also presents the search for decays of the Z boson into quarkonium pairs. The SM prediction for  $\mathcal{B}(Z \rightarrow J/\psi J/\psi)$  calculated in the framework of nonrelativistic quantum chromodynamics (QCD) and leading twist light cone models is of the order of  $10^{-12}$  [35]. The Z decay into a  $J/\psi$  meson and a lepton pair, which is dominated by the electromagnetic fragmentation process, was observed at a rate consistent with SM predictions [36].

The results presented in this Letter are based on proton-proton (pp) collision data recorded in 2016–2018 with the CMS detector at a center-of-mass energy of  $\sqrt{s} = 13$  TeV, amounting to an integrated luminosity of  $138 \text{ fb}^{-1}$  in the  $ZJ/\psi$  channel and  $133 \text{ fb}^{-1}$

in the quarkonium pair channels, where the second number is slightly smaller due to a delayed trigger deployment.

## 2. The CMS detector

The CMS apparatus [37] is a multipurpose, nearly hermetic detector, designed to trigger on and identify electrons, muons, photons, and hadrons. The CMS superconducting solenoid provides an axial magnetic field of 3.8 T. Within the solenoid volume are a silicon pixel and strip tracker, a lead tungstate crystal electromagnetic calorimeter (ECAL), and a brass and scintillator hadron calorimeter, each composed of a barrel and two endcap sections. Forward calorimeters extend the pseudorapidity ( $\eta$ ) coverage provided by the barrel and endcap detectors. Muons are detected in gas-ionization chambers embedded in the steel flux-return yoke outside the solenoid. A more detailed description of the CMS detector, together with a definition of the coordinate system used and the relevant kinematic variables, can be found in Ref. [37].

The silicon tracker measures charged particles within the range  $|\eta| < 2.5$ . It consists of pixel and strip detector modules. An entirely new pixel detector has been installed during a technical stop between the 2016 and 2017 data-taking periods, featuring an all-silicon device with four layers in the barrel and three disks in the endcaps [38], providing four pixel detector measurements within a range  $|\eta| < 3$ .

Muons are measured in the range  $|\eta| < 2.4$ , with detection planes made using three technologies: drift tubes, cathode strip chambers, and resistive-plate chambers. The single-muon trigger efficiency exceeds 90% over the full  $\eta$  range, and the efficiency to reconstruct and identify muons is greater than 96%. Matching muons to tracks measured in the silicon tracker results in relative transverse momentum ( $p_T$ ) resolutions of 1% in the barrel and 3% in the endcaps for muons with  $p_T$  up to 100 GeV, and better than 7% in the barrel for muons with  $p_T$  up to 1 TeV [39].

An electron is reconstructed by combining an energy measurement in the ECAL with a momentum measurement in the tracker. The ECAL consists of 75,848 lead tungstate crystals, which provide coverage of  $|\eta| < 1.48$  in the barrel region and  $1.48 < |\eta| < 3.00$  in the two endcap regions (EE). Preshower detectors consisting of two planes of silicon sensors interleaved with a total of three radiation lengths of lead are located in front of each EE detector. The momentum resolution for electrons with  $p_T \approx 45$  GeV from  $Z \rightarrow e^+e^-$  decays ranges from 2 to 5%. It is generally better in the barrel region than in the endcaps, and also depends on the bremsstrahlung energy emitted by the electron as it traverses the material in front of the ECAL [40].

The candidate vertex with the largest value of summed physics-object  $p_T^2$  is taken to be the primary pp interaction vertex. The physics objects are the jets, clustered using the jet finding algorithm [41,42] with the tracks assigned to candidate vertices as inputs, and the associated missing transverse momentum, taken as the negative vector  $p_T$  sum of those jets. Other collision vertices in the event are considered to have originated from additional inelastic pp collisions in each bunch crossing, referred to as pileup (PU). The average number of PU interactions during the 2016 data-taking period was 23, and increased to 32 during the 2017 and 2018 data-taking periods.

Events of interest are selected using a two-tiered trigger system. The first level, composed of custom hardware processors, uses information from the calorimeters and muon detectors to select events at a rate of around 100 kHz within a fixed latency of about 4  $\mu\text{s}$  [43]. The second level, known as the high-level trigger, consists of a farm of processors running a version of the full event reconstruction software optimized for fast processing, and reduces the event rate to around 1 kHz before data storage [44].

Dedicated triggers were deployed to enhance the selection of events of interest for the present study. Single-lepton (muon or electron) triggers are used to study the  $Z/\psi$  channel. The single muon trigger requires an isolated muon with  $p_T > 27$  GeV. The single electron trigger requires an isolated electron having  $p_T > 27$  (35, 32) GeV during the year 2016 (2017, 2018). The meson-specific triggers are used to study quarkonium channels. They require the presence of at least three muons with  $p_T > 2$  GeV. Two of those must be oppositely charged and originate from a common vertex with a probability greater than 0.5%, as determined by a Kalman vertex filter [45], thus suppressing random combinations of two muons. The  $J/\psi$ - ( $Y$ -) specific trigger requires a dimuon system invariant mass to be between 2.95 and 3.25 (8.5 and 11.4) GeV. The dimuon system  $p_T$  is required to exceed 3.5 GeV for the years 2017 and 2018.

### 3. Simulated samples

Simulated samples of the H and Z boson signals are used to estimate the expected signal yields and model the distribution of signal events in the four-lepton invariant mass. The SM Higgs boson signals are simulated at next-to-leading order (NLO) in perturbative QCD with the POWHEG v2.0 Monte Carlo event generator [46,47], which includes the gluon-gluon fusion and vector boson fusion production processes. The parton distribution function (PDF) set used is NNPDF3.1 [48]. The JHUGen 7.1.4 generator [49,50] is used to decay the Higgs boson into Z bosons and Q mesons. The generator is interfaced with PYTHIA 8.226 [51] for parton showering, hadronization, and underlying event simulation using the CUETP8M1 [52] tune. The total SM Higgs boson production cross section for the calculation of branching fractions is taken from the LHC Higgs cross section working group [7].

The Z bosons are produced and decayed with the PYTHIA 8.226 generator [51] which implements the LO matrix element calculation interfaced with parton showering, hadronization, and underlying event simulation for which the tune CUETP8M1 [52] is used. For the calculation of the branching fraction, the inclusive SM Z boson production cross section is obtained according to Ref. [53], where the prediction includes the next-to-next-to-leading order (NNLO) QCD contributions, and the NLO electroweak corrections from FEWZ3.1 [54] calculated using the NLO PDF set NNPDF3.0. The generated events are then reweighted to match the  $p_T$ -spectrum of the Z boson predicted at NLO [46,47,55].

The generated events are processed through a detailed simulation of the CMS detector based on GEANT4 [56]. Simulated events include additional pp interactions. Events are then reweighted to match the PU profile observed in data.

The acceptance of the final states changes with the angular distribution of leptons in the decay. The distribution of the decay angle  $\theta$ , defined as the angle between the positive lepton momentum in the rest frame of the intermediate particle ( $J/\psi$  meson or Y meson or Z boson) with respect to the direction of this intermediate particle in the rest frame of the parent particle (H or Z boson), is proportional to  $(1 + \lambda_\theta \cos^2 \theta)$ , where  $\lambda_\theta$  is the average polar anisotropy parameter [57]. According to Refs. [27,35,58], longitudinal polarization is expected for the Z boson and assumed for mesons. In this Letter, the nominal results are obtained using a signal acceptance calculated for the longitudinally polarized case ( $\lambda_\theta = -1$ ). Two other extreme scenarios where the intermediate particles are both either fully transversely polarized ( $\lambda_\theta = +1$ ) or unpolarized ( $\lambda_\theta = 0$ ) are also considered. No azimuthal anisotropies are considered.

### 4. Event reconstruction

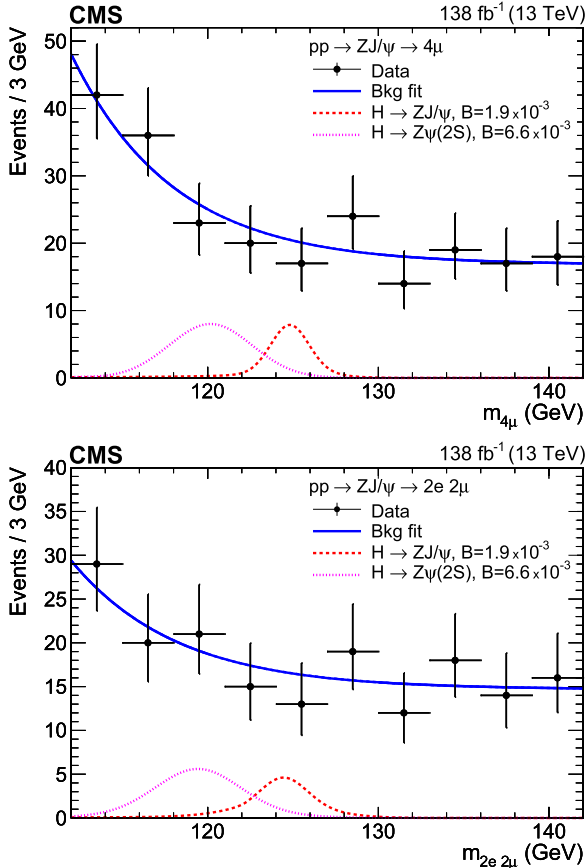
In a first step, events with at least  $2\mu$  and two additional  $\ell$  ( $\ell = e$  or  $\mu$ ) are selected. The  $J/\psi$  and Y candidates are built from  $\mu^+\mu^-$  pairs, and Z candidates from  $\mu^+\mu^-$  or  $e^+e^-$  pairs.

Muons are reconstructed by combining information from the silicon tracker and the muon system [39]. The matching between tracks reconstructed in each of the subsystems proceeds either outside-in, starting from a track in the muon system, or inside-out, starting from a track provided by the silicon tracker. In the latter case, tracks that match track segments in only one or two detector layers of the muon system are also considered in the analysis to collect very low  $p_T$  muons that may not have sufficient energy to penetrate the entire muon system. The muons are selected from the reconstructed muon track candidates that match with at least one segment in any muon detector layer. The number of silicon tracker layers with hits used in the muon track candidate has to be greater than 5 and include at least one pixel detector layer. To suppress muons originating from non-prompt hadron decays, the impact parameter of each muon track, computed with respect to the position of the primary pp interaction vertex, is required to be less than 0.3 (20.0) cm in the transverse plane (along the longitudinal axis). Events with at least four such muons with  $p_T > 3$  GeV and  $|\eta| < 2.4$  and zero sum of charges are accepted. To measure the isolation of the leading (highest  $p_T$ ) muon candidate from other hadronic activity in the event, a cone of size  $\Delta R = \sqrt{(\Delta\eta)^2 + (\Delta\phi)^2} = 0.3$  is constructed around its momentum direction, where  $\phi$  is the azimuthal angle in radians. The  $p_T$  sum of the reconstructed inner-detector tracks originating from the primary pp interaction vertex within the cone has to be less than 50% of the muon's  $p_T$ . The leading muon  $p_T$  is excluded from the sum and the subleading muon  $p_T$  is also excluded if this muon falls within the isolation cone of the leading muon.

The energy of electron candidates is determined from a combination of the track momentum at the main interaction vertex, the corresponding ECAL cluster energy, and the energy sum of all bremsstrahlung photons attached to the track [40]. The dielectron mass resolution for  $Z \rightarrow e^+e^-$  decays when both electrons are in the ECAL barrel is 1.9%, and is 2.9% when they are in the endcaps. To reduce contamination from particles incorrectly identified as electrons, reconstructed electrons are required to pass a multivariate electron identification discriminant. This discriminant, which is described in Ref. [40], combines information about the quality of tracks, the shower shape, kinematic quantities, and hadronic activity in the vicinity of the reconstructed electron. Isolation variables are also included among the discriminant's inputs. Therefore, no additional isolation requirements are applied. The selection based on the multivariate identification discriminant has an electron identification efficiency of 90% while the rate of misidentifying other particles as electrons is 2–5%. Events with at least two oppositely charged electrons with  $p_T > 3$  GeV and  $|\eta| < 2.5$  and two oppositely charged muons are accepted.

Further selection criteria are applied to the different four-lepton final states to achieve the lowest expected upper limit at 95% CL [59–61]. This optimization is performed with data with the signal regions removed [62] and replaced by simulated events, and with the simulated signal shape.

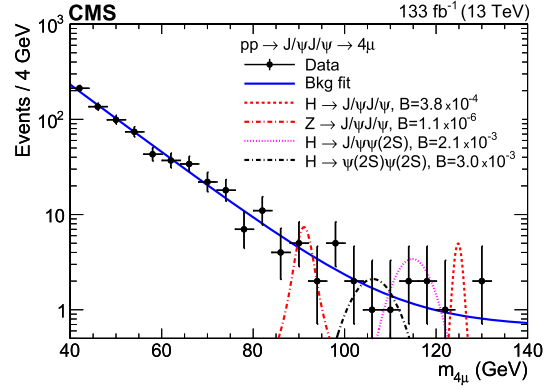
The selection in the  $ZJ/\psi$  channel requires each dilepton resonance to have  $p_T > 5$  GeV, and the candidate invariant mass to lie in the region 80–100 GeV for the Z boson (3.0–3.2 GeV for the  $J/\psi$  meson). The dimuon mass resolution is about 1%. Each dilepton must fit to a common vertex with probability greater than 1%, determined by a vertex fit probability. The four-lepton candidate must have  $p_T > 5$  GeV, and the fit to a common four-lepton vertex must have a probability of greater than 1%. A total of 230 (177) candidate events are found in the  $4\mu$  ( $2e2\mu$ ) invariant mass



**Fig. 2.** Four-lepton invariant mass distributions, (upper) for  $ZJ/\psi \rightarrow 4\mu$  candidates and (lower) for  $ZJ/\psi \rightarrow 2e2\mu$  candidates. The result of the maximum likelihood fit to the background (Bkg) is superimposed (solid blue line). For illustrative purposes, the plots show simulated  $H \rightarrow ZJ/\psi$  (dashed red line) and  $H \rightarrow Z\psi(2S)$  (dotted magenta line) signals normalized to their observed upper limit branching fractions at 95% CL as obtained in this analysis, where  $\psi(2S)$  decays into  $J/\psi + X$ .

window of 112–142 GeV. The lower limit of the four-lepton invariant mass ( $m_{4\ell}$ ) range is chosen to exclude the region close to the  $ZJ/\psi$  threshold where  $m_{4\ell}$  changes rapidly, which is difficult to model. The selection criteria for the decay  $H \rightarrow Z\psi(2S)$ , where  $\psi(2S)$  decays inclusively into  $J/\psi$ , are identical. The respective  $m_{4\ell}$  distributions are shown in Fig. 2. Non-resonant background in the dilepton invariant mass distributions is found to contribute about 20%. These events are an additional source of background which does not concentrate in a specific region of  $m_{4\ell}$ , and thus are accounted for by an empirical parameterization of the background in the fit as described in the next section.

In the case of the  $J/\psi$  pair channel, each dimuon has to be fit to a common vertex with a probability greater than 0.5%. In addition, the  $J/\psi$  candidate must have  $p_T > 3.5$  GeV, matching the trigger requirement, and the invariant masses of the higher and lower- $p_T$   $J/\psi$  candidates have to be within 0.10 and 0.15 GeV, respectively, of the nominal mass of the  $J/\psi$  meson. The mass window of the subleading  $J/\psi$  is wider to allow monitoring the reduction of the sideband population as the selection progresses. In order to suppress contributions from non-prompt hadrons, separately produced  $J/\psi$  mesons and muons from other sources, the four-muon vertex fit probability of  $J/\psi$  pairs must be greater than 5%. Finally, the absolute value of the difference in rapidity ( $|\Delta Y|$ ) between the two  $J/\psi$  candidates has to be less than 3. This criterion marginally affects the signal, while removing about 20% of the background events. After the selection, 720 events are found in data in the 40–140 GeV four-muon invariant mass ( $m_{4\mu}$ ) range. Non-resonant



**Fig. 3.** The four-muon invariant mass distribution of  $J/\psi J/\psi$  candidates (error bars for empty bins are not shown). The result of the maximum likelihood fit to the background (Bkg) is superimposed (solid blue line). For illustrative purposes, the plots show the distributions for simulated signals of the Higgs and Z boson decaying into  $J/\psi J/\psi$  (dashed and dashed-dotted red lines). The signals for the Higgs boson decays  $H \rightarrow \psi(2S)J/\psi$  (dotted magenta line) and  $H \rightarrow \psi(2S)\psi(2S)$  (dashed-dotted black line) are also shown, where  $\psi(2S)$  decays into  $J/\psi$ . Each signal is normalized to their observed 95% CL upper limit branching fraction from this analysis.

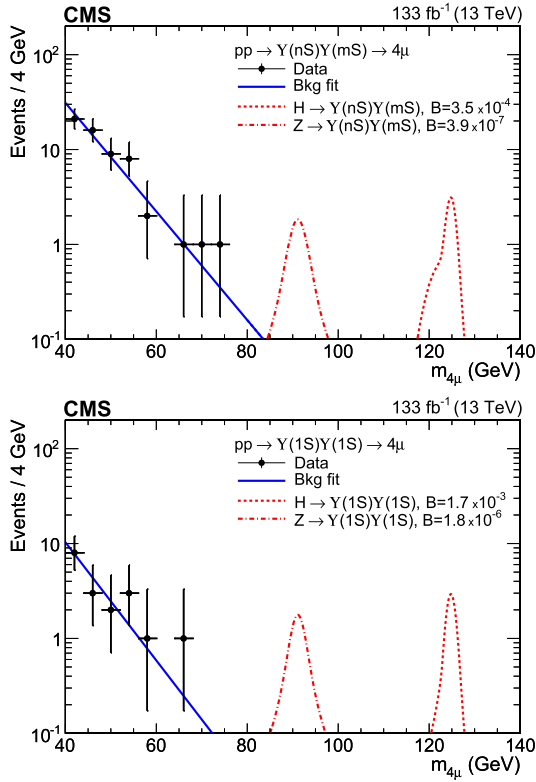
background in the dilepton invariant mass distributions is found to be negligible. Fig. 3 shows the  $m_{4\mu}$  distribution of the  $J/\psi J/\psi$  candidates.

An  $Y$  pair candidate event must have at least four muons each with  $p_T > 4$  GeV. The  $Y(nS)$  ( $n = 1, 2, 3$ ) and  $Y(1S)$  candidates are formed with oppositely charged muon pairs with  $p_T > 5$  GeV, and the dimuon invariant mass within the range 9.0–10.7 GeV and 9.0–9.7 GeV, respectively. In order to suppress random combinations, dimuon and four-muon objects are required to have a vertex fit probability greater than 1%. Between two candidate dimuons, the  $|\Delta Y|$  value has to be less than 2.3 and the azimuthal angle difference has to be greater than 1 radian. The four-muon combination must have  $p_T > 5$  GeV and an absolute rapidity of less than 1.7. After applying the selection criteria in data, 59  $Y(nS)Y(mS)$  (18  $Y(1S)Y(1S)$ ) candidate events are found in the 40–140 GeV  $m_{4\mu}$  range. Fig. 4 upper and lower show the  $m_{4\mu}$  distributions for  $Y(nS)Y(mS)$  and  $Y(1S)Y(1S)$  candidates, respectively.

The differences in the efficiencies between data and simulation for the trigger, offline lepton reconstruction, identification, and isolation requirements are corrected by reweighting the simulated events with data-to-simulation correction factors, which are obtained with the “tag-and-probe” method [63] using  $J/\psi \rightarrow \mu^+\mu^-$  and  $Z \rightarrow \ell^+\ell^-$  events. The total signal efficiency, including kinematic acceptance, trigger, reconstruction, identification, and isolation efficiencies, for the  $ZJ/\psi \rightarrow 4\mu$  ( $ZJ/\psi \rightarrow 2e2\mu$ ) decays with longitudinally polarized  $J/\psi$  and  $Z$ , is found from simulation to be around 30 (24)%. For the Higgs boson decays into  $Y(nS)Y(mS)$  and  $J/\psi J/\psi$ , the corresponding total efficiencies are about 31 and 30%, respectively. For the  $Z$  boson, the corresponding values are about 28 and 32%.

## 5. Signal extraction and systematic uncertainties

Unbinned extended maximum likelihood fits [64] to the  $m_{4\ell}$  distributions are performed. Yields for signals and backgrounds are free parameters in the fit. The background shapes in the  $m_{4\ell}$  distributions are obtained from data and are described by an exponential plus constant function. For the  $ZJ/\psi \rightarrow 4\mu$  channel, the H signal is parameterized with a sum of a Gaussian and a Crystal Ball function [65], and for the  $ZJ/\psi \rightarrow 2e2\mu$  channel with two Crystal Ball functions. Similarly, the H signal in the  $J/\psi$  pair channel is described with a double-Gaussian function, and in the  $Y$  pair channel with a combination of a Gaussian and Crystal Ball functions. In all



**Fig. 4.** The four-muon invariant mass distributions for  $Y(nS)Y(mS)$  (upper) and  $Y(1S)Y(1S)$  (lower) candidates (error bars for empty bins are not shown). The result of the maximum likelihood fit to background (Bkg) is superimposed (solid blue line). For illustrative purposes, the plots show the distributions for simulated Higgs and Z boson signals (dashed and dashed-dotted red lines) normalized to their observed 95% CL upper limit branching fractions from this analysis.

combined functions the mean is a common parameter. The simulated Z boson signal is described with a Voigtian function with the resonance width fixed to the world-average value [66]. The mass resolution and the mean are taken from the fit to the simulation, and they are fixed in the fit to data. The background function from the fit to data in the  $Z/J/\psi$  and  $J/\psi/J/\psi$  channels is superimposed as solid blue lines in Figs. 2 and 3, respectively. In the Y pair sample, no events are observed above the  $m_{4\mu}$  of 80 GeV. The  $m_{4\mu}$  distribution below 80 GeV is well described solely by an exponential function. Fig. 4 shows the observed  $m_{4\mu}$  distribution with the fit superimposed.

Separate fits are performed to the  $m_{4\ell}$  distributions for the different signal hypotheses. Signal shapes for the Higgs boson in decays involving the inclusive transition from  $\psi(2S)$  to  $J/\psi$  meson are modeled with a combination of the same functions as used for the Higgs boson directly decaying into ground state mesons (direct signal). For the fits to the feed-down channels, the background functions are identical and parameters are fixed to the ones from the previous direct signal fits.

Systematic uncertainties originate from imperfect knowledge of the detector and imperfect signal modeling. Systematic uncertainties considered in this analysis are listed below.

- i) The integrated luminosities for the 2016, 2017, and 2018 data-taking years have 1.2–2.5% individual uncertainties [67–69], while the overall uncertainty for the 2016–2018 period is 1.6%.
- ii) The differences between data and simulation for the trigger, offline muon reconstruction, identification, and isolation efficiencies are corrected by reweighting the simulated events with data-to-simulation correction factors, which are obtained with the tag-and-probe method using  $J/\psi$  (or  $Z$ )  $\rightarrow \mu^+\mu^-$

events. The resulting scale factors in muon identification, isolation, and trigger efficiencies are observed to deviate from unity by less than 2(2), 0.5(0.5) and 1(3)% for the  $Z/J/\psi$  (QQ) channel. Analogously, the uncertainty in the electron reconstruction, identification and trigger efficiency is found to be about 2% [40].

- iii) The relative difference in the four-lepton vertex criterion between data and simulation is evaluated with  $Z/J/\psi$  ( $J/\psi$  pair) event samples. It is found to be less than 2(3)% for the  $Z/J/\psi$  (QQ) channel.
- iv) Differences in the lepton momentum scale and resolution in data and simulation are estimated from  $J/\psi$  and Z dilepton signals and extrapolated to the four-lepton signals. The systematic uncertainty is estimated as the relative change in the upper limit when varying the signal mass mean and width by these differences. They are found to be less than 1(3)% in the  $4\mu$  ( $2e2\mu$ ) channel.
- v) The theoretical uncertainties in the production cross section for the H (Z boson) are  $\pm 3.2\%$  ( $\pm 1.7\%$ ) due to the choice of the PDF and the value of the strong coupling constant [7,48,70], and  $^{+4.6}_{-6.7}\%$  ( $\pm 3.5\%$ ) due to the renormalization and factorization scale choice [70–73].
- vi) A common parameterization for each signal model is used for the entire run period. The relative uncertainty in the signal model due to the change in detector conditions in each year is determined to be 1(2)% for the  $Z/J/\psi$  (QQ) channels.
- vii) The background is alternatively parameterized with a second order Chebyshev polynomial or a power law function. The relative uncertainty due to the choice of the background function is found to be negligible.
- viii) The uncertainties in the Z,  $J/\psi$ , and Y branching fractions to lepton pairs, and  $\mathcal{B}(\psi(2S) \rightarrow J/\psi + X)$  are taken from Ref. [66].

## 6. Results

No evidence for the Higgs or Z boson signal is found in any of these channels. The results of this analysis are presented as upper limits on the branching fractions and are set at the 95% CL. Limits are determined with the modified frequentist  $CL_s$  criterion, in which the profile likelihood ratio modified for upper limits is used as the test statistic [59–61]. Systematic uncertainties are incorporated in the likelihood as nuisance parameters. The observed and median expected exclusion limits for the branching fractions at 95% CL for the H and Z boson decays are listed in Table 1. Figs. 2–4 show the distributions for simulated boson signals in different search channels as dashed lines. The signals are normalized to their observed 95% upper limit branching fractions.

The results for  $\mathcal{B}(H \rightarrow Z/J/\psi)$  and  $\mathcal{B}(H \rightarrow Z\psi(2S))$  each are obtained by combining the channels with  $Z \rightarrow e^+e^-$  and  $Z \rightarrow \mu^+\mu^-$ . The values for  $\mathcal{B}(J/\psi \rightarrow \mu^+\mu^-)$ ,  $\mathcal{B}(Z \rightarrow \ell^+\ell^-)$  and  $\mathcal{B}(\psi(2S) \rightarrow J/\psi)$  are taken from Ref. [66]. This analysis does not distinguish between the three Y(nS) states. To calculate their contribution to the corresponding H and Z boson branching fractions, the coupling strength of the bosons to any Y(nS) pairing is assumed to be the same. All Y states can directly decay into muon pairs with the different branching fractions taken from Ref. [66]. In addition, it is assumed that the Y states could be the result of a one step transition  $Y(3S) \rightarrow Y(2S)$ ,  $Y(3S) \rightarrow Y(1S)$ , or  $Y(2S) \rightarrow Y(1S)$  before decaying into muons [66]. Consequently, in the  $Y(1S)Y(1S)$  channel, the feed-down transitions from  $Y(3S)$  and  $Y(2S)$  to  $Y(1S)$  are included. The observed upper limit branching fractions at 95% CL agree with the expected limits. In the case of  $H \rightarrow Z/J/\psi$ , the upper limit branching fraction is about 800 times higher than the SM prediction [17]. In the  $H \rightarrow Y(nS)Y(mS)$  channel, the observed upper limit branching fraction is found to be about one order of

**Table 1**

Exclusion limits at 95% CL on the branching fractions of the H and Z boson decays. The second column lists the observed limits for the case that both intermediate particles are longitudinally polarized ( $\lambda_\theta = -1$ ) as described in the text. The third column shows the median expected limits with the upper and lower bounds in the expected 68% CL intervals. The last two columns list observed upper limits for unpolarized ( $\lambda_\theta = 0$ ) and transversely polarized ( $\lambda_\theta = +1$ ) intermediate particles.

Process	Observed Longitudinal	Expected Longitudinal	Observed	
			Unpolarized	Transverse
<b>Higgs boson channel</b>				
$\mathcal{B}(H \rightarrow ZJ/\psi)$	$1.9 \times 10^{-3}$	$(2.6_{-0.7}^{+1.1}) \times 10^{-3}$	$2.4 \times 10^{-3}$	$2.8 \times 10^{-3}$
$\mathcal{B}(H \rightarrow Z\psi(2S))$	$6.6 \times 10^{-3}$	$(7.1_{-2.0}^{+2.8}) \times 10^{-3}$	$8.3 \times 10^{-3}$	$9.4 \times 10^{-3}$
$\mathcal{B}(H \rightarrow J/\psi J/\psi)$	$3.8 \times 10^{-4}$	$(4.6_{-0.6}^{+2.0}) \times 10^{-4}$	$4.7 \times 10^{-4}$	$5.2 \times 10^{-4}$
$\mathcal{B}(H \rightarrow \psi(2S)J/\psi)$	$2.1 \times 10^{-3}$	$(1.4_{-0.4}^{+0.6}) \times 10^{-3}$	$2.6 \times 10^{-3}$	$2.9 \times 10^{-3}$
$\mathcal{B}(H \rightarrow \psi(2S)\psi(2S))$	$3.0 \times 10^{-3}$	$(3.3_{-0.9}^{+1.5}) \times 10^{-3}$	$3.6 \times 10^{-3}$	$4.7 \times 10^{-3}$
$\mathcal{B}(H \rightarrow Y(nS)Y(mS))$	$3.5 \times 10^{-4}$	$(3.6_{-0.3}^{+0.2}) \times 10^{-4}$	$4.3 \times 10^{-4}$	$4.6 \times 10^{-4}$
$\mathcal{B}(H \rightarrow Y(1S)Y(1S))$	$1.7 \times 10^{-3}$	$(1.7_{-0.1}^{+0.1}) \times 10^{-3}$	$2.0 \times 10^{-3}$	$2.2 \times 10^{-3}$
<b>Z boson channel</b>				
$\mathcal{B}(Z \rightarrow J/\psi J/\psi)$	$11 \times 10^{-7}$	$(9.5_{-2.6}^{+3.8}) \times 10^{-7}$	$14 \times 10^{-7}$	$16 \times 10^{-7}$
$\mathcal{B}(Z \rightarrow Y(nS)Y(mS))$	$3.9 \times 10^{-7}$	$(4.0_{-0.3}^{+0.3}) \times 10^{-7}$	$4.9 \times 10^{-7}$	$5.6 \times 10^{-7}$
$\mathcal{B}(Z \rightarrow Y(1S)Y(1S))$	$1.8 \times 10^{-6}$	$(1.8_{-0.0}^{+0.1}) \times 10^{-6}$	$2.2 \times 10^{-6}$	$2.4 \times 10^{-6}$

magnitude higher than SM predictions that assume dominance of direct quark couplings [23]. The factors are larger for all other channels. Tabulated results are available in the HEPData record for this analysis [74].

## 7. Summary

This Letter presents the first search for decays of the Higgs boson (H) into a Z boson and a  $J/\psi$  meson in four-lepton final states. Data from proton-proton collisions at  $\sqrt{s} = 13$  TeV, corresponding to an integrated luminosity of about  $138 \text{ fb}^{-1}$ , are used. Using the same data, decays of the Higgs and Z boson into quarkonium pairs are also searched for. No excess of a Higgs or Z boson signal above background is found in any of the searched channels and upper limits on branching fractions ( $\mathcal{B}$ ) at the 95% confidence level for various polarization scenarios are set. The Higgs boson decay is also searched for in channels where, before decaying into muon pairs, one or both  $J/\psi$  mesons could be the result of an inclusive  $\psi(2S)$  to  $J/\psi$  transition, and the  $Y(nS)$  ( $n = 1, 2$ ) mesons could be the result of inclusive transitions from  $Y(nS)$  ( $n = 2, 3$ ) mesons. The observed upper limits for the Higgs and Z boson decays for longitudinally polarized mesons are  $\mathcal{B}(H \rightarrow ZJ/\psi) < 1.9 \times 10^{-3}$ ,  $\mathcal{B}(H \rightarrow Z\psi(2S)) < 6.6 \times 10^{-3}$ ,  $\mathcal{B}(H \rightarrow J/\psi J/\psi) < 3.8 \times 10^{-4}$ ,  $\mathcal{B}(H \rightarrow \psi(2S)J/\psi) < 2.1 \times 10^{-3}$ ,  $\mathcal{B}(H \rightarrow \psi(2S)\psi(2S)) < 3.0 \times 10^{-3}$ ,  $\mathcal{B}(H \rightarrow Y(nS)Y(mS)) < 3.5 \times 10^{-4}$ ,  $\mathcal{B}(H \rightarrow Y(1S)Y(1S)) < 1.7 \times 10^{-3}$ ,  $\mathcal{B}(Z \rightarrow J/\psi J/\psi) < 11 \times 10^{-7}$ ,  $\mathcal{B}(Z \rightarrow Y(nS)Y(mS)) < 3.9 \times 10^{-7}$ , and  $\mathcal{B}(Z \rightarrow Y(1S)Y(1S)) < 1.8 \times 10^{-6}$ . The observed upper limit branching fraction for  $H \rightarrow ZJ/\psi$  is about 800 times the value predicted by the standard model [17–19]. For  $H \rightarrow Y(nS)Y(mS)$  the upper limit is about one order of magnitude higher than predicted by earlier standard model calculations [23].

## Declaration of competing interest

The authors declare that they have no known competing financial interests or personal relationships that could have appeared to influence the work reported in this paper.

## Data availability

Release and preservation of data used by the CMS Collaboration as the basis for publications is guided by the CMS policy as stated in “[CMS data preservation, re-use and open access policy](#)”.

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## The CMS Collaboration

### A. Tumasyan<sup>1</sup>

Yerevan Physics Institute, Yerevan, Armenia

W. Adam, J.W. Andrejkovic, T. Bergauer, S. Chatterjee, K. Damanakis, M. Dragicevic, A. Escalante Del Valle, R. Frühwirth<sup>2</sup>, M. Jeitler<sup>2</sup>, N. Krammer, L. Lechner, D. Liko, I. Mikulec, P. Paulitsch, F.M. Pitters, J. Schieck<sup>2</sup>, R. Schöffbeck, D. Schwarz, S. Templ, W. Waltenberger, C.-E. Wulz<sup>2</sup>

Institut für Hochenergiephysik, Vienna, Austria

M.R. Darwish<sup>3</sup>, E.A. De Wolf, T. Janssen, T. Kello<sup>4</sup>, A. Lelek, H. Rejeb Sfar, P. Van Mechelen, S. Van Putte, N. Van Remortel

Universiteit Antwerpen, Antwerpen, Belgium

E.S. Bols, J. D'Hondt, A. De Moor, M. Delcourt, H. El Faham, S. Lowette, S. Moortgat, A. Morton, D. Müller, A.R. Sahasransu, S. Tavernier, W. Van Doninck, D. Vannerom

Vrije Universiteit Brussel, Brussel, Belgium



D. Beghin, B. Clerbaux, G. De Lentdecker, L. Favart, K. Lee, M. Mahdavihorrani, I. Makarenko, A. Malara, S. Paredes, L. P  tr  , A. Popov, N. Postiau, E. Starling, L. Thomas, M. Vanden Bemden, C. Vander Velde, P. Vanlaer

*Universit   Libre de Bruxelles, Bruxelles, Belgium*

D. Dobur, J. Knolle, L. Lambrecht, G. Mestdach, M. Niedziela, C. Rend  n, C. Roskas, A. Samalan, K. Skovpen, M. Tytgat, N. Van Den Bossche, B. Vermassen, L. Wezenbeek

*Ghent University, Ghent, Belgium*

A. Benecke, A. Bethani, G. Bruno, F. Bury, C. Caputo, P. David, C. Delaere, I.S. Donertas, A. Giammanco, K. Jaffel, Sa. Jain, V. Lema  tre, K. Mondal, J. Prisciandaro, A. Taliercio, T.T. Tran, P. Vischia, S. Wertz

*Universit   Catholique de Louvain, Louvain-la-Neuve, Belgium*

G.A. Alves, C. Hensel, A. Moraes, P. Rebello Teles

*Centro Brasileiro de Pesquisas F  sicas, Rio de Janeiro, Brazil*

W.L. Ald   J  nior, M. Alves Gallo Pereira, M. Barroso Ferreira Filho, H. Brandao Malbouisson, W. Carvalho, J. Chinellato<sup>5</sup>, E.M. Da Costa, G.G. Da Silveira<sup>6</sup>, D. De Jesus Damiao, V. Dos Santos Sousa, S. Fonseca De Souza, J. Martins<sup>7</sup>, C. Mora Herrera, K. Mota Amarilo, L. Mundim, H. Nogima, A. Santoro, S.M. Silva Do Amaral, A. Sznajder, M. Thiel, F. Torres Da Silva De Araujo<sup>8</sup>, A. Vilela Pereira

*Universidade do Estado do Rio de Janeiro, Rio de Janeiro, Brazil*

C.A. Bernardes<sup>6</sup>, L. Calligaris, T.R. Fernandez Perez Tomei, E.M. Gregores, P.G. Mercadante, S.F. Novaes, Sandra S. Padula

*Universidade Estadual Paulista, Universidade Federal do ABC, S  o Paulo, Brazil*

A. Aleksandrov, G. Antchev, R. Hadjiiska, P. Iaydjiev, M. Misheva, M. Rodozov, M. Shopova, G. Sultanov

*Institute for Nuclear Research and Nuclear Energy, Bulgarian Academy of Sciences, Sofia, Bulgaria*

A. Dimitrov, T. Ivanov, L. Litov, B. Pavlov, P. Petkov, A. Petrov

*University of Sofia, Sofia, Bulgaria*

T. Cheng, T. Javaid<sup>9</sup>, M. Mittal, L. Yuan

*Beihang University, Beijing, China*

M. Ahmad, G. Bauer<sup>10</sup>, C. Dozen, Z. Hu, Y. Wang, K. Yi<sup>10,11</sup>

*Department of Physics, Tsinghua University, Beijing, China*

E. Chapon, G.M. Chen<sup>9</sup>, H.S. Chen<sup>9</sup>, M. Chen<sup>9</sup>, F. Iemmi, A. Kapoor, H. Liao, Z.-A. Liu<sup>12</sup>, V. Milosevic, F. Monti, R. Sharma, J. Tao, J. Thomas-Wilsker, J. Wang, H. Zhang, J. Zhao

*Institute of High Energy Physics, Beijing, China*

A. Agapitos, Y. An, Y. Ban, C. Chen, A. Levin, Q. Li, X. Lyu, Y. Mao, S.J. Qian, D. Wang, J. Xiao, H. Yang

*State Key Laboratory of Nuclear Physics and Technology, Peking University, Beijing, China*

M. Lu, Z. You

*Sun Yat-Sen University, Guangzhou, China*

X. Gao<sup>4</sup>, D. Leggat, H. Okawa, Y. Zhang

*Institute of Modern Physics and Key Laboratory of Nuclear Physics and Ion-beam Application (MOE) - Fudan University, Shanghai, China*

Z. Lin, M. Xiao

*Zhejiang University, Hangzhou, Zhejiang, China*

C. Avila, A. Cabrera, C. Florez, J. Fraga

*Universidad de Los Andes, Bogota, Colombia*

J. Mejia Guisao, F. Ramirez, J.D. Ruiz Alvarez

*Universidad de Antioquia, Medellin, Colombia*

D. Giljanovic, N. Godinovic, D. Lelas, I. Puljak

*University of Split, Faculty of Electrical Engineering, Mechanical Engineering and Naval Architecture, Split, Croatia*

Z. Antunovic, M. Kovac, T. Sculac

*University of Split, Faculty of Science, Split, Croatia*

V. Brigljevic, B.K. Chitroda, D. Ferencek, D. Majumder, M. Roguljic, A. Starodumov<sup>13</sup>, T. Susa

*Institute Rudjer Boskovic, Zagreb, Croatia*

A. Attikis, K. Christoforou, G. Kole, M. Kolosova, S. Konstantinou, J. Mousa, C. Nicolaou, F. Ptochos, P.A. Razis, H. Rykaczewski, H. Saka

*University of Cyprus, Nicosia, Cyprus*

M. Finger<sup>13</sup>, M. Finger Jr.<sup>13</sup>, A. Kveton

*Charles University, Prague, Czech Republic*

E. Ayala

*Escuela Politecnica Nacional, Quito, Ecuador*

E. Carrera Jarrin

*Universidad San Francisco de Quito, Quito, Ecuador*

A.A. Abdelalim<sup>14,15</sup>, Y. Assran<sup>16,17</sup>

*Academy of Scientific Research and Technology of the Arab Republic of Egypt, Egyptian Network of High Energy Physics, Cairo, Egypt*

M.A. Mahmoud, Y. Mohammed

*Center for High Energy Physics (CHEP-FU), Fayoum University, El-Fayoum, Egypt*

S. Bhowmik, R.K. Dewanjee, K. Ehataht, M. Kadastik, S. Nandan, C. Nielsen, J. Pata, M. Raidal, L. Tani, C. Veelken

*National Institute of Chemical Physics and Biophysics, Tallinn, Estonia*

P. Eerola, H. Kirschenmann, K. Osterberg, M. Voutilainen

*Department of Physics, University of Helsinki, Helsinki, Finland*

S. Bharthuar, E. Brücken, F. Garcia, J. Havukainen, M.S. Kim, R. Kinnunen, T. Lampén, K. Lassila-Perini, S. Lehti, T. Lindén, M. Lotti, L. Martikainen, M. Myllymäki, J. Ott, M.m. Rantanen, H. Siikonen, E. Tuominen, J. Tuominiemi

*Helsinki Institute of Physics, Helsinki, Finland*

P. Luukka, H. Petrow, T. Tuuva

*Lappeenranta-Lahti University of Technology, Lappeenranta, Finland*

C. Amendola, M. Besancon, F. Couderc, M. Dejardin, D. Denegri, J.L. Faure, F. Ferri, S. Ganjour, P. Gras, G. Hamel de Monchenault, P. Jarry, B. Lenzi, V. Lohezic, J. Malcles, J. Rander, A. Rosowsky, M.Ö. Sahin, A. Savoy-Navarro<sup>18</sup>, P. Simkina, M. Titov

IRFU, CEA, Université Paris-Saclay, Gif-sur-Yvette, France

S. Ahuja, C. Baldenegro Barrera, F. Beaudette, M. Bonanomi, A. Buchot Perraguin, P. Busson, A. Cappati, C. Charlot, O. Davignon, B. Diab, G. Falmagne, B.A. Fontana Santos Alves, S. Ghosh, R. Granier de Cassagnac, A. Hakimi, B. Harikrishnan, J. Motta, M. Nguyen, C. Ochando, P. Paganini, L. Portales, J. Rembser, R. Salerno, U. Sarkar, J.B. Sauvan, Y. Sirois, A. Tarabini, E. Vernazza, A. Zabi, A. Zghiche

Laboratoire Leprince-Ringuet, CNRS/IN2P3, Ecole Polytechnique, Institut Polytechnique de Paris, Palaiseau, France

J.-L. Agram<sup>19</sup>, J. Andrea, D. Apparú, D. Bloch, G. Bourgatte, J.-M. Brom, E.C. Chabert, C. Collard, D. Darej, J.-C. Fontaine<sup>19</sup>, U. Goerlach, C. Grimault, A.-C. Le Bihan, E. Nibigira, P. Van Hove

Université de Strasbourg, CNRS, IPHC UMR 7178, Strasbourg, France

S. Beauceron, C. Bernet, G. Boudoul, C. Camen, A. Carle, N. Chanon, D. Contardo, P. Depasse, H. El Mamouni, J. Fay, S. Gascon, M. Gouzevitch, B. Ille, I.B. Laktineh, M. Lethuillier, L. Mirabito, S. Perries, K. Shchablo, V. Sordini, L. Torterotot, M. Vander Donckt, S. Viret

Institut de Physique des 2 Infinis de Lyon (IP2I), Villeurbanne, France

A. Khvedelidze<sup>13</sup>, I. Lomidze, Z. Tsamalaidze<sup>13</sup>

Georgian Technical University, Tbilisi, Georgia

V. Botta, L. Feld, K. Klein, M. Lipinski, D. Meuser, A. Pauls, N. Röwert, J. Schulz, M. Teroerde

RWTH Aachen University, I. Physikalisches Institut, Aachen, Germany

A. Dodonova, N. Eich, D. Eliseev, M. Erdmann, P. Fackeldey, B. Fischer, T. Hebbeker, K. Hoepfner, F. Ivone, M.y. Lee, L. Mastrolorenzo, M. Merschmeyer, A. Meyer, S. Mondal, S. Mukherjee, D. Noll, A. Novak, A. Pozdnyakov, Y. Rath, H. Reithler, A. Schmidt, S.C. Schuler, A. Sharma, L. Vigilante, S. Wiedenbeck, S. Zaleski

RWTH Aachen University, III. Physikalisches Institut A, Aachen, Germany

C. Dziwok, G. Flügge, W. Haj Ahmad<sup>20</sup>, O. Hlushchenko, T. Kress, A. Nowack, O. Pooth, D. Roy, A. Stahl<sup>21</sup>, T. Ziemons, A. Zötz

RWTH Aachen University, III. Physikalisches Institut B, Aachen, Germany

H. Aarup Petersen, M. Aldaya Martin, P. Asmuss, S. Baxter, M. Bayatmakou, O. Behnke, A. Bermúdez Martínez, S. Bhattacharya, A.A. Bin Anuar, F. Blekman<sup>22</sup>, K. Borras<sup>23</sup>, D. Brunner, A. Campbell, A. Cardini, C. Cheng, F. Colombina, S. Consuegra Rodríguez, G. Correia Silva, M. De Silva, L. Didukh, G. Eckerlin, D. Eckstein, L.I. Estevez Banos, O. Filatov, E. Gallo<sup>22</sup>, A. Geiser, A. Giraldi, G. Greau, A. Grohsjean, V. Guglielmi, M. Guthoff, A. Jafari<sup>24</sup>, N.Z. Jomhari, B. Kaech, A. Kasem<sup>23</sup>, M. Kasemann, H. Kaveh, C. Kleinwort, R. Kogler, D. Krücker, W. Lange, K. Lipka, W. Lohmann<sup>25</sup>, R. Mankel, I.-A. Melzer-Pellmann, M. Mendizabal Morentin, J. Metwally, A.B. Meyer, M. Meyer, G. Milella, M. Mormile, A. Mussgiller, A. Nürnberg, Y. Otari, D. Pérez Adán, D. Pitzl, A. Raspereza, B. Ribeiro Lopes, J. Rübenach, A. Saggio, A. Saibel, M. Savitskyi, M. Scham<sup>26,23</sup>, V. Scheurer, S. Schnake<sup>23</sup>, P. Schütze, C. Schwanenberger<sup>22</sup>, M. Shchedrolosiev, R.E. Sosa Ricardo, D. Stafford, N. Tonon<sup>†</sup>, M. Van De Klundert, F. Vazzoler, R. Walsh, D. Walter, Q. Wang, Y. Wen, K. Wichmann, L. Wiens<sup>23</sup>, C. Wissing, S. Wuchterl, A. Zimmermann Castro Santos

Deutsches Elektronen-Synchrotron, Hamburg, Germany

R. Aggleton, S. Albrecht, S. Bein, L. Benato, P. Connor, K. De Leo, M. Eich, K. El Morabit, F. Feindt, A. Fröhlich, C. Garbers, E. Garutti, M. Hajheidari, J. Haller, A. Hinzmann, G. Kasieczka, R. Klanner, T. Kramer, V. Kutzner, J. Lange, T. Lange, A. Lobanov, C. Matthies, A. Mehta, L. Moureaux, A. Nigamova, K.J. Pena Rodriguez, M. Rieger, O. Rieger, P. Schleper, M. Schröder, J. Schwandt, H. Stadie, G. Steinbrück, A. Tews

*University of Hamburg, Hamburg, Germany*

J. Bechtel, S. Brommer, M. Burkart, E. Butz, R. Caspart, T. Chwalek, W. De Boer<sup>†</sup>, A. Dierlamm, A. Droll, N. Faltermann, M. Giffels, J.O. Gosewisch, A. Gottmann, F. Hartmann<sup>21</sup>, C. Heidecker, M. Horzela, U. Husemann, P. Keicher, R. Koppenhöfer, S. Maier, S. Mitra, Th. Müller, M. Neukum, G. Quast, K. Rabbertz, J. Rauser, D. Savoie, M. Schnepf, D. Seith, I. Shvetsov, H.J. Simonis, R. Ulrich, J. van der Linden, R.F. Von Cube, M. Wassmer, M. Weber, S. Wieland, R. Wolf, S. Wozniewski, S. Wunsch

*Karlsruher Institut fuer Technologie, Karlsruhe, Germany*

G. Anagnostou, P. Assiouras, G. Daskalakis, A. Kyriakis, A. Stakia

*Institute of Nuclear and Particle Physics (INPP), NCSR Demokritos, Aghia Paraskevi, Greece*

M. Diamantopoulou, D. Karasavvas, P. Kontaxakis, C.K. Koraka, A. Manousakis-Katsikakis, A. Panagiotou, I. Papavergou, N. Saoulidou, K. Theofilatos, E. Tziaferi, K. Vellidis, E. Vourliotis

*National and Kapodistrian University of Athens, Athens, Greece*

G. Bakas, K. Kousouris, I. Papakrivopoulos, G. Tsipolitis, A. Zacharopoulou

*National Technical University of Athens, Athens, Greece*

K. Adamidis, I. Bestintzanos, I. Evangelou, C. Foudas, P. Gianneios, P. Katsoulis, P. Kokkas, N. Manthos, I. Papadopoulos, J. Strologas

*University of Ioánnina, Ioánnina, Greece*

M. Csanád, K. Farkas, M.M.A. Gadallah<sup>27</sup>, S. Lökös<sup>28</sup>, P. Major, K. Mandal, G. Pásztor, A.J. Rádl, O. Surányi, G.I. Veres

*MTA-ELTE Lendület CMS Particle and Nuclear Physics Group, Eötvös Loránd University, Budapest, Hungary*

M. Bartók<sup>29</sup>, G. Bencze, C. Hajdu, D. Horvath<sup>30,31</sup>, F. Sikler, V. Veszpremi

*Wigner Research Centre for Physics, Budapest, Hungary*

S. Czellar, D. Fasanella, F. Fienga, J. Karancsi<sup>29</sup>, J. Molnar, Z. Szillasi, D. Teyssier

*Institute of Nuclear Research ATOMKI, Debrecen, Hungary*

P. Raics, Z.L. Trocsanyi<sup>32</sup>, B. Ujvari<sup>33</sup>

*Institute of Physics, University of Debrecen, Debrecen, Hungary*

T. Csorgo<sup>34</sup>, F. Nemes<sup>34</sup>, T. Novak

*Karoly Robert Campus, MATE Institute of Technology, Gyongyos, Hungary*

J. Babbar, S. Bansal, S.B. Beri, V. Bhatnagar, G. Chaudhary, S. Chauhan, N. Dhingra<sup>35</sup>, R. Gupta, A. Kaur, H. Kaur, M. Kaur, P. Kumari, M. Meena, K. Sandeep, J.B. Singh<sup>36</sup>, A.K. Viridi

*Panjab University, Chandigarh, India*

A. Ahmed, A. Bhardwaj, B.C. Choudhary, M. Gola, S. Keshri, A. Kumar, M. Naimuddin, P. Priyanka, K. Ranjan, S. Saumya, A. Shah

*University of Delhi, Delhi, India*

R. Bhattacharya, S. Bhattacharya, D. Bhowmik, S. Dutta, S. Dutta, B. Gomber<sup>37</sup>, M. Maity<sup>38</sup>, P. Palit, P.K. Rout, G. Saha, B. Sahu, S. Sarkar, M. Sharan

*Saha Institute of Nuclear Physics, HBNI, Kolkata, India*

P.K. Behera, S.C. Behera, P. Kalbhor, J.R. Komaragiri<sup>39</sup>, D. Kumar<sup>39</sup>, A. Muhammad, L. Panwar<sup>39</sup>, R. Pradhan, P.R. Pujahari, A. Sharma, A.K. Sikdar, P.C. Tiwari<sup>39</sup>

*Indian Institute of Technology Madras, Madras, India*

K. Naskar<sup>40</sup>

*Bhabha Atomic Research Centre, Mumbai, India*

T. Aziz, S. Dugad, M. Kumar, G.B. Mohanty, P. Suryadevara

*Tata Institute of Fundamental Research-A, Mumbai, India*

S. Banerjee, R. Chudasama, M. Guchait, S. Karmakar, S. Kumar, G. Majumder, K. Mazumdar, S. Mukherjee

*Tata Institute of Fundamental Research-B, Mumbai, India*

S. Bahinipati<sup>41</sup>, C. Kar, P. Mal, T. Mishra, V.K. Muraleedharan Nair Bindhu<sup>42</sup>, A. Nayak<sup>42</sup>, P. Saha, N. Sur, S.K. Swain, D. Vats<sup>42</sup>

*National Institute of Science Education and Research, An OCC of Homi Bhabha National Institute, Bhubaneswar, Odisha, India*

A. Alpana, S. Dube, B. Kansal, A. Laha, S. Pandey, A. Rastogi, S. Sharma

*Indian Institute of Science Education and Research (IISER), Pune, India*

H. Bakhshiansohi<sup>43</sup>, E. Khazaie, M. Zeinali<sup>44</sup>

*Isfahan University of Technology, Isfahan, Iran*

S. Chenarani<sup>45</sup>, S.M. Etesami, M. Khakzad, M. Mohammadi Najafabadi

*Institute for Research in Fundamental Sciences (IPM), Tehran, Iran*

M. Grunewald

*University College Dublin, Dublin, Ireland*

M. Abbrescia<sup>a,b</sup>, R. Aly<sup>a,c,14</sup>, C. Aruta<sup>a,b</sup>, A. Colaleo<sup>a</sup>, D. Creanza<sup>a,c</sup>, N. De Filippis<sup>a,c</sup>, M. De Palma<sup>a,b</sup>, A. Di Florio<sup>a,b</sup>, W. Elmetenawee<sup>a,b</sup>, F. Errico<sup>a,b</sup>, L. Fiore<sup>a</sup>, G. Iaselli<sup>a,c</sup>, M. Ince<sup>a,b</sup>, S. Lezki<sup>a,b</sup>, G. Maggi<sup>a,c</sup>, M. Maggi<sup>a</sup>, I. Margjeka<sup>a,b</sup>, V. Mastrapasqua<sup>a,b</sup>, S. My<sup>a,b</sup>, S. Nuzzo<sup>a,b</sup>, A. Pellecchia<sup>a,b</sup>, A. Pompili<sup>a,b</sup>, G. Pugliese<sup>a,c</sup>, D. Ramos<sup>a</sup>, A. Ranieri<sup>a</sup>, G. Selvaggi<sup>a,b</sup>, L. Silvestris<sup>a</sup>, F.M. Simone<sup>a,b</sup>, Ü. Sözbilir<sup>a</sup>, R. Venditti<sup>a</sup>, P. Verwilligen<sup>a</sup>

<sup>a</sup> INFN Sezione di Bari, Bari, Italy

<sup>b</sup> Università di Bari, Bari, Italy

<sup>c</sup> Politecnico di Bari, Bari, Italy

G. Abbiendi<sup>a</sup>, C. Battilana<sup>a,b</sup>, D. Bonacorsi<sup>a,b</sup>, L. Borbonovi<sup>a</sup>, L. Brigliadori<sup>a</sup>, R. Campanini<sup>a,b</sup>, P. Capiluppi<sup>a,b</sup>, A. Castro<sup>a,b</sup>, F.R. Cavallo<sup>a</sup>, C. Ciocca<sup>a</sup>, M. Cuffiani<sup>a,b</sup>, G.M. Dallavalle<sup>a</sup>, T. Diotallevi<sup>a,b</sup>, F. Fabbri<sup>a</sup>, A. Fanfani<sup>a,b</sup>, P. Giacomelli<sup>a</sup>, L. Giommi<sup>a,b</sup>, C. Grandi<sup>a</sup>, L. Guiducci<sup>a,b</sup>, S. Lo Meo<sup>a,46</sup>, L. Lunerti<sup>a,b</sup>, S. Marcellini<sup>a</sup>, G. Masetti<sup>a</sup>, F.L. Navarria<sup>a,b</sup>, A. Perrotta<sup>a</sup>, F. Primavera<sup>a,b</sup>, A.M. Rossi<sup>a,b</sup>, T. Rovelli<sup>a,b</sup>, G.P. Siroli<sup>a,b</sup>

<sup>a</sup> INFN Sezione di Bologna, Bologna, Italy

<sup>b</sup> Università di Bologna, Bologna, Italy

S. Albergo<sup>a,b,47</sup>, S. Costa<sup>a,b,47</sup>, A. Di Mattia<sup>a</sup>, R. Potenza<sup>a,b</sup>, A. Tricomi<sup>a,b,47</sup>, C. Tuve<sup>a,b</sup>

<sup>a</sup> INFN Sezione di Catania, Catania, Italy<sup>b</sup> Università di Catania, Catania, Italy

G. Barbagli<sup>a</sup>, B. Camaiani<sup>a,b</sup>, A. Cassese<sup>a</sup>, R. Ceccarelli<sup>a,b</sup>, V. Ciulli<sup>a,b</sup>, C. Civinini<sup>a</sup>, R. D'Alessandro<sup>a,b</sup>,  
E. Focardi<sup>a,b</sup>, G. Latino<sup>a,b</sup>, P. Lenzi<sup>a,b</sup>, M. Lizzo<sup>a,b</sup>, M. Meschini<sup>a</sup>, S. Paoletti<sup>a</sup>, R. Seidita<sup>a,b</sup>,  
G. Sguazzoni<sup>a</sup>, L. Viliani<sup>a</sup>

<sup>a</sup> INFN Sezione di Firenze, Firenze, Italy<sup>b</sup> Università di Firenze, Firenze, Italy

L. Benussi, S. Bianco, D. Piccolo

INFN Laboratori Nazionali di Frascati, Frascati, Italy

M. Bozzo<sup>a,b</sup>, F. Ferro<sup>a</sup>, R. Mulargia<sup>a</sup>, E. Robutti<sup>a</sup>, S. Tosi<sup>a,b</sup>

<sup>a</sup> INFN Sezione di Genova, Genova, Italy<sup>b</sup> Università di Genova, Genova, Italy

A. Benaglia<sup>a</sup>, G. Boldrini<sup>a</sup>, F. Brivio<sup>a,b</sup>, F. Ceteorelli<sup>a,b</sup>, F. De Guio<sup>a,b</sup>, M.E. Dinardo<sup>a,b</sup>, P. Dini<sup>a</sup>, S. Gennai<sup>a</sup>,  
A. Ghezzi<sup>a,b</sup>, P. Govoni<sup>a,b</sup>, L. Guzzi<sup>a,b</sup>, M.T. Lucchini<sup>a,b</sup>, M. Malberti<sup>a</sup>, S. Malvezzi<sup>a</sup>, A. Massironi<sup>a</sup>,  
D. Menasce<sup>a</sup>, L. Moroni<sup>a</sup>, M. Paganoni<sup>a,b</sup>, D. Pedrini<sup>a</sup>, B.S. Pinolini<sup>a</sup>, S. Ragazzi<sup>a,b</sup>, N. Redaelli<sup>a</sup>,  
T. Tabarelli de Fatis<sup>a,b</sup>, D. Zuolo<sup>a,b</sup>

<sup>a</sup> INFN Sezione di Milano-Bicocca, Milano, Italy<sup>b</sup> Università di Milano-Bicocca, Milano, Italy

S. Buontempo<sup>a</sup>, F. Carnevali<sup>a,b</sup>, N. Cavallo<sup>a,c</sup>, A. De Iorio<sup>a,b</sup>, F. Fabozzi<sup>a,c</sup>, A.O.M. Iorio<sup>a,b</sup>, L. Lista<sup>a,b,48</sup>,  
S. Meola<sup>a,d,21</sup>, P. Paolucci<sup>a,21</sup>, B. Rossi<sup>a</sup>, C. Sciacca<sup>a,b</sup>

<sup>a</sup> INFN Sezione di Napoli, Napoli, Italy<sup>b</sup> Università di Napoli "Federico II", Napoli, Italy<sup>c</sup> Università della Basilicata, Potenza, Italy<sup>d</sup> Università G. Marconi, Roma, Italy

P. Azzi<sup>a</sup>, N. Bacchetta<sup>a</sup>, D. Bisello<sup>a,b</sup>, P. Bortignon<sup>a</sup>, A. Bragagnolo<sup>a,b</sup>, R. Carlin<sup>a,b</sup>, P. Checchia<sup>a</sup>,  
T. Dorigo<sup>a</sup>, U. Dosselli<sup>a</sup>, F. Gasparini<sup>a,b</sup>, U. Gasparini<sup>a,b</sup>, G. Grosso<sup>a</sup>, L. Layer<sup>a,49</sup>, E. Lusiani<sup>a</sup>,  
M. Margoni<sup>a,b</sup>, F. Marini<sup>a</sup>, A.T. Meneguzzo<sup>a,b</sup>, J. Pazzini<sup>a,b</sup>, P. Ronchese<sup>a,b</sup>, R. Rossin<sup>a,b</sup>, F. Simonetto<sup>a,b</sup>,  
G. Strong<sup>a</sup>, M. Tosi<sup>a,b</sup>, H. Yarar<sup>a,b</sup>, M. Zanetti<sup>a,b</sup>, P. Zotto<sup>a,b</sup>, A. Zucchetta<sup>a,b</sup>, G. Zumerle<sup>a,b</sup>

<sup>a</sup> INFN Sezione di Padova, Padova, Italy<sup>b</sup> Università di Padova, Padova, Italy<sup>c</sup> Università di Trento, Trento, Italy

C. Aimè<sup>a,b</sup>, A. Braghieri<sup>a</sup>, S. Calzaferri<sup>a,b</sup>, D. Fiorina<sup>a,b</sup>, P. Montagna<sup>a,b</sup>, S.P. Ratti<sup>a,b</sup>, V. Re<sup>a</sup>,  
C. Riccardi<sup>a,b</sup>, P. Salvini<sup>a</sup>, I. Vai<sup>a</sup>, P. Vitulo<sup>a,b</sup>

<sup>a</sup> INFN Sezione di Pavia, Pavia, Italy<sup>b</sup> Università di Pavia, Pavia, Italy

P. Asenov<sup>a,50</sup>, G.M. Bilei<sup>a</sup>, D. Ciangottini<sup>a,b</sup>, L. Fanò<sup>a,b</sup>, M. Magherini<sup>a,b</sup>, G. Mantovani<sup>a,b</sup>, V. Mariani<sup>a,b</sup>,  
M. Menichelli<sup>a</sup>, F. Moscatelli<sup>a,50</sup>, A. Piccinelli<sup>a,b</sup>, M. Presilla<sup>a,b</sup>, A. Rossi<sup>a,b</sup>, A. Santocchia<sup>a,b</sup>, D. Spiga<sup>a</sup>,  
T. Tedeschi<sup>a,b</sup>

<sup>a</sup> INFN Sezione di Perugia, Perugia, Italy<sup>b</sup> Università di Perugia, Perugia, Italy

P. Azzurri<sup>a</sup>, G. Bagliesi<sup>a</sup>, V. Bertacchi<sup>a,c</sup>, L. Bianchini<sup>a,b</sup>, T. Boccali<sup>a</sup>, E. Bossini<sup>a,b</sup>, D. Bruschini<sup>a,c</sup>,  
R. Castaldi<sup>a</sup>, M.A. Ciocci<sup>a,b</sup>, V. D'Amante<sup>a,d</sup>, R. Dell'Orso<sup>a</sup>, M.R. Di Domenico<sup>a,d</sup>, S. Donato<sup>a</sup>, A. Giassi<sup>a</sup>,  
F. Ligabue<sup>a,c</sup>, E. Manca<sup>a,c</sup>, G. Mandorli<sup>a,c</sup>, D. Matos Figueiredo<sup>a</sup>, A. Messineo<sup>a,b</sup>, M. Musich<sup>a,b</sup>, F. Palla<sup>a</sup>,  
S. Parolia<sup>a,b</sup>, G. Ramirez-Sanchez<sup>a,c</sup>, A. Rizzi<sup>a,b</sup>, G. Rolandi<sup>a,c</sup>, S. Roy Chowdhury<sup>a,c</sup>, A. Scribano<sup>a</sup>,  
N. Shafiei<sup>a,b</sup>, P. Spagnolo<sup>a</sup>, R. Tenchini<sup>a</sup>, G. Tonelli<sup>a,b</sup>, N. Turini<sup>a,d</sup>, A. Venturi<sup>a</sup>, P.G. Verdini<sup>a</sup>

<sup>a</sup> INFN Sezione di Pisa, Pisa, Italy<sup>b</sup> Università di Pisa, Pisa, Italy<sup>c</sup> Scuola Normale Superiore di Pisa, Pisa, Italy<sup>d</sup> Università di Siena, Siena, Italy

P. Barria <sup>a</sup>, M. Campana <sup>a,b</sup>, F. Cavallari <sup>a</sup>, D. Del Re <sup>a,b</sup>, E. Di Marco <sup>a</sup>, M. Diemoz <sup>a</sup>, E. Longo <sup>a,b</sup>,  
 P. Meridiani <sup>a</sup>, G. Organtini <sup>a,b</sup>, F. Pandolfi <sup>a</sup>, R. Paramatti <sup>a,b</sup>, C. Quaranta <sup>a,b</sup>, S. Rahatlou <sup>a,b</sup>, C. Rovelli <sup>a</sup>,  
 F. Santanastasio <sup>a,b</sup>, L. Soffi <sup>a</sup>, R. Tramontano <sup>a,b</sup>

<sup>a</sup> INFN Sezione di Roma, Roma, Italy<sup>b</sup> Sapienza Università di Roma, Roma, Italy

N. Amapane <sup>a,b</sup>, R. Arcidiacono <sup>a,c</sup>, S. Argiro <sup>a,b</sup>, M. Arneodo <sup>a,c</sup>, N. Bartosik <sup>a</sup>, R. Bellan <sup>a,b</sup>, A. Bellora <sup>a,b</sup>,  
 J. Berenguer Antequera <sup>a,b</sup>, C. Biino <sup>a</sup>, N. Cartiglia <sup>a</sup>, M. Costa <sup>a,b</sup>, R. Covarelli <sup>a,b</sup>, N. Demaria <sup>a</sup>,  
 M. Grippo <sup>a,b</sup>, B. Kiani <sup>a,b</sup>, F. Legger <sup>a</sup>, C. Mariotti <sup>a</sup>, S. Maselli <sup>a</sup>, A. Mecca <sup>a,b</sup>, E. Migliore <sup>a,b</sup>, E. Monteil <sup>a,b</sup>,  
 M. Monteno <sup>a</sup>, M.M. Obertino <sup>a,b</sup>, G. Ortona <sup>a</sup>, L. Pacher <sup>a,b</sup>, N. Pastrone <sup>a</sup>, M. Pelliccioni <sup>a</sup>, M. Ruspa <sup>a,c</sup>,  
 K. Shchelina <sup>a</sup>, F. Siviero <sup>a,b</sup>, V. Sola <sup>a</sup>, A. Solano <sup>a,b</sup>, D. Soldi <sup>a,b</sup>, A. Staiano <sup>a</sup>, M. Tornago <sup>a,b</sup>, D. Trocino <sup>a</sup>,  
 G. Umoret <sup>a,b</sup>, A. Vagnerini <sup>a,b</sup>

<sup>a</sup> INFN Sezione di Torino, Torino, Italy<sup>b</sup> Università di Torino, Torino, Italy<sup>c</sup> Università del Piemonte Orientale, Novara, Italy

S. Belforte <sup>a</sup>, V. Candelise <sup>a,b</sup>, M. Casarsa <sup>a</sup>, F. Cossutti <sup>a</sup>, A. Da Rold <sup>a,b</sup>, G. Della Ricca <sup>a,b</sup>, G. Sorrentino <sup>a,b</sup>

<sup>a</sup> INFN Sezione di Trieste, Trieste, Italy<sup>b</sup> Università di Trieste, Trieste, Italy

S. Dogra, C. Huh, B. Kim, D.H. Kim, G.N. Kim, J. Kim, J. Lee, S.W. Lee, C.S. Moon, Y.D. Oh, S.I. Pak,  
 S. Sekmen, Y.C. Yang

*Kyungpook National University, Daegu, Korea*

H. Kim, D.H. Moon

*Chonnam National University, Institute for Universe and Elementary Particles, Kwangju, Korea*

E. Asilar, B. Francois, T.J. Kim, J. Park

*Hanyang University, Seoul, Korea*

S. Cho, S. Choi, B. Hong, K. Lee, K.S. Lee, J. Lim, J. Park, S.K. Park, J. Yoo

*Korea University, Seoul, Korea*

J. Goh, A. Gurtu

*Kyung Hee University, Department of Physics, Seoul, Korea*

H.S. Kim, Y. Kim, G.B. Yu

*Sejong University, Seoul, Korea*

J. Almond, J.H. Bhyun, J. Choi, S. Jeon, J. Kim, J.S. Kim, S. Ko, H. Kwon, H. Lee, S. Lee, B.H. Oh, M. Oh,  
 S.B. Oh, H. Seo, U.K. Yang, I. Yoon

*Seoul National University, Seoul, Korea*

W. Jang, D.Y. Kang, Y. Kang, S. Kim, B. Ko, J.S.H. Lee, Y. Lee, J.A. Merlin, I.C. Park, Y. Roh, M.S. Ryu,  
 D. Song, I.J. Watson, S. Yang

*University of Seoul, Seoul, Korea*

S. Ha, H.D. Yoo

Yonsei University, Department of Physics, Seoul, Korea

**M. Choi, H. Lee, Y. Lee, I. Yu**

Sungkyunkwan University, Suwon, Korea

**T. Beyrouthy, Y. Maghrbi**

College of Engineering and Technology, American University of the Middle East (AUM), Dasman, Kuwait

**K. Dreimanis, V. Veckalns**

Riga Technical University, Riga, Latvia

**M. Ambrozas, A. Carvalho Antunes De Oliveira, A. Juodagalvis, A. Rinkevicius, G. Tamulaitis**

Vilnius University, Vilnius, Lithuania

**N. Bin Norjoharuddeen, S.Y. Hoh<sup>51</sup>, Z. Zolkapli**

National Centre for Particle Physics, Universiti Malaya, Kuala Lumpur, Malaysia

**J.F. Benitez, A. Castaneda Hernandez, H.A. Encinas Acosta, L.G. Gallegos Maríñez, M. León Coello, J.A. Murillo Quijada, A. Sehrawat, L. Valencia Palomo**

Universidad de Sonora (UNISON), Hermosillo, Mexico

**G. Ayala, H. Castilla-Valdez, E. De La Cruz-Burelo, I. Heredia-De La Cruz<sup>52</sup>, R. Lopez-Fernandez, C.A. Mondragon Herrera, D.A. Perez Navarro, R. Reyes-Almanza, A. Sánchez Hernández**

Centro de Investigacion y de Estudios Avanzados del IPN, Mexico City, Mexico

**C. Oropeza Barrera, F. Vazquez Valencia**

Universidad Iberoamericana, Mexico City, Mexico

**I. Pedraza, H.A. Salazar Ibarguen, C. Uribe Estrada**

Benemerita Universidad Autonoma de Puebla, Puebla, Mexico

**I. Bubanja, J. Mijuskovic<sup>53</sup>, N. Raicevic**

University of Montenegro, Podgorica, Montenegro

**D. Krofcheck**

University of Auckland, Auckland, New Zealand

**P.H. Butler**

University of Canterbury, Christchurch, New Zealand

**A. Ahmad, M.I. Asghar, A. Awais, M.I.M. Awan, M. Gul, H.R. Hoorani, W.A. Khan, M.A. Shah, M. Shoaib, M. Waqas**

National Centre for Physics, Quaid-I-Azam University, Islamabad, Pakistan

**V. Avati, L. Grzanka, M. Malawski**

AGH University of Science and Technology Faculty of Computer Science, Electronics and Telecommunications, Krakow, Poland

**H. Bialkowska, M. Bluj, B. Boimska, M. Górski, M. Kazana, M. Szeleper, P. Zalewski**

National Centre for Nuclear Research, Swierk, Poland

**K. Bunkowski, K. Doroba, A. Kalinowski, M. Konecki, J. Krolikowski**

Institute of Experimental Physics, Faculty of Physics, University of Warsaw, Warsaw, Poland



M. Araujo, P. Bargassa, D. Bastos, A. Boletti, P. Faccioli, M. Gallinaro, J. Hollar, N. Leonardo, T. Niknejad, M. Pisano, J. Seixas, O. Toldaiev, J. Varela

*Laboratório de Instrumentação e Física Experimental de Partículas, Lisboa, Portugal*

P. Adzic<sup>54</sup>, M. Dordevic, P. Milenovic, J. Milosevic

*VINCA Institute of Nuclear Sciences, University of Belgrade, Belgrade, Serbia*

M. Aguilar-Benitez, J. Alcaraz Maestre, A. Álvarez Fernández, I. Bachiller, M. Barrio Luna, Cristina F. Bedoya, C.A. Carrillo Montoya, M. Cepeda, M. Cerrada, N. Colino, B. De La Cruz, A. Delgado Peris, J.P. Fernández Ramos, J. Flix, M.C. Fouz, O. Gonzalez Lopez, S. Goy Lopez, J.M. Hernandez, M.I. Josa, J. León Holgado, D. Moran, Á. Navarro Tobar, C. Perez Dengra, A. Pérez-Calero Yzquierdo, J. Puerta Pelayo, I. Redondo, L. Romero, S. Sánchez Navas, L. Urda Gómez, C. Willmott

*Centro de Investigaciones Energéticas Medioambientales y Tecnológicas (CIEMAT), Madrid, Spain*

J.F. de Trocóniz

*Universidad Autónoma de Madrid, Madrid, Spain*

B. Alvarez Gonzalez, J. Cuevas, J. Fernandez Menendez, S. Folgueras, I. Gonzalez Caballero, J.R. González Fernández, E. Palencia Cortezon, C. Ramón Álvarez, V. Rodríguez Bouza, A. Soto Rodríguez, A. Trapote, N. Trevisani, C. Vico Villalba

*Universidad de Oviedo, Instituto Universitario de Ciencias y Tecnologías Espaciales de Asturias (ICTEA), Oviedo, Spain*

J.A. Brochero Cifuentes, I.J. Cabrillo, A. Calderon, J. Duarte Campderros, M. Fernandez, C. Fernandez Madrazo, P.J. Fernández Manteca, A. García Alonso, G. Gomez, C. Martinez Rivero, P. Martinez Ruiz del Arbol, F. Matorras, P. Matorras Cuevas, J. Piedra Gomez, C. Prieels, A. Ruiz-Jimeno, L. Scodellaro, I. Vila, J.M. Vizán Garcia

*Instituto de Física de Cantabria (IFCA), CSIC-Universidad de Cantabria, Santander, Spain*

M.K. Jayananda, B. Kailasapathy<sup>55</sup>, D.U.J. Sonnadara, D.D.C. Wickramarathna

*University of Colombo, Colombo, Sri Lanka*

W.G.D. Dharmaratna, K. Liyanage, N. Perera, N. Wickramage

*University of Ruhuna, Department of Physics, Matara, Sri Lanka*

T.K. Aarrestad, D. Abbaneo, J. Alimena, E. Auffray, G. Auzinger, J. Baechler, P. Baillon<sup>†</sup>, D. Barney, J. Bendavid, M. Bianco, B. Bilin, A. Bocci, C. Caillol, T. Camporesi, M. Capeans Garrido, G. Cerminara, N. Chernyavskaya, S.S. Chhibra, S. Choudhury, M. Cipriani, L. Cristella, D. d'Enterria, A. Dabrowski, A. David, A. De Roeck, M.M. Defranchis, M. Deile, M. Dobson, M. Dünser, N. Dupont, A. Elliott-Peisert, F. Fallavollita<sup>56</sup>, A. Florent, L. Forthomme, G. Franzoni, W. Funk, S. Ghosh, S. Giani, D. Gigi, K. Gill, F. Glege, L. Gouskos, E. Govorkova, M. Haranko, J. Hegeman, V. Innocente, T. James, P. Janot, J. Kaspar, J. Kieseler, M. Komm, N. Kratochwil, C. Lange, S. Laurila, P. Lecoq, A. Lintuluoto, C. Lourenço, B. Maier, L. Malgeri, S. Mallios, M. Mannelli, A.C. Marini, F. Meijers, S. Mersi, E. Meschi, F. Moortgat, M. Mulders, S. Orfanelli, L. Orsini, F. Pantaleo, E. Perez, M. Peruzzi, A. Petrilli, G. Petrucciani, A. Pfeiffer, M. Pierini, D. Piparo, M. Pitt, H. Qu, T. Quast, D. Rabaday, A. Racz, G. Reales Gutiérrez, M. Rovere, H. Sakulin, J. Salfeld-Nebgen, S. Scarfi, M. Selvaggi, A. Sharma, P. Silva, W. Snoeys, P. Sphicas<sup>57</sup>, A.G. Stahl Leitner, S. Summers, K. Tatar, V.R. Tavolaro, D. Treille, P. Tropea, A. Tsiros, J. Wanczyk<sup>58</sup>, K.A. Wozniak, W.D. Zeuner

*CERN, European Organization for Nuclear Research, Geneva, Switzerland*

L. Caminada<sup>59</sup>, A. Ebrahimi, W. Erdmann, R. Horisberger, Q. Ingram, H.C. Kaestli, D. Kotlinski, M. Missiroli<sup>59</sup>, L. Noehte<sup>59</sup>, T. Rohe

Paul Scherrer Institut, Villigen, Switzerland

K. Androso<sup>58</sup>, M. Backhaus, P. Berger, A. Calandri, A. De Cosa, G. Dissertori, M. Dittmar, M. Donegà, C. Dorfer, F. Eble, K. Gedia, F. Glessgen, T.A. Gómez Espinosa, C. Grab, D. Hits, W. Lusterhmann, A.-M. Lyon, R.A. Manzoni, L. Marchese, C. Martin Perez, A. Mascellani<sup>58</sup>, M.T. Meinhard, F. Nessi-Tedaldi, J. Niedziela, F. Pauss, V. Perovic, S. Pigazzini, M.G. Ratti, M. Reichmann, C. Reissel, T. Reitenspiess, B. Ristic, D. Ruini, D.A. Sanz Becerra, J. Steggemann<sup>58</sup>, D. Valsecchi<sup>21</sup>, R. Wallny

ETH Zurich - Institute for Particle Physics and Astrophysics (IPA), Zurich, Switzerland

C. Amsler<sup>60</sup>, P. Bäertschi, C. Botta, D. Brzhechko, M.F. Canelli, K. Cormier, A. De Wit, R. Del Burgo, J.K. Heikkilä, M. Huwiler, W. Jin, A. Jofrehei, B. Kilminster, S. Leontsinis, S.P. Liechti, A. Macchiolo, P. Meiring, V.M. Mikuni, U. Molinatti, I. Neutelings, A. Reimers, P. Robmann, S. Sanchez Cruz, K. Schweiger, M. Senger, Y. Takahashi

Universität Zürich, Zurich, Switzerland

C. Adloff<sup>61</sup>, C.M. Kuo, W. Lin, A. Roy, T. Sarkar<sup>38</sup>, S.S. Yu

National Central University, Chung-Li, Taiwan

L. Ceard, Y. Chao, K.F. Chen, P.H. Chen, P.s. Chen, H. Cheng, W.-S. Hou, Y.y. Li, R.-S. Lu, E. Paganis, A. Psallidas, A. Steen, H.y. Wu, E. Yazgan, P.r. Yu

National Taiwan University (NTU), Taipei, Taiwan

B. Asavapibhop, C. Asawatangtrakuldee, N. Srimanobhas

Chulalongkorn University, Faculty of Science, Department of Physics, Bangkok, Thailand

F. Boran, S. Damarsecin<sup>62</sup>, Z.S. Demiroglu, F. Dolek, I. Dumanoglu<sup>63</sup>, E. Eskut, Y. Guler<sup>64</sup>, E. Gurpinar Guler<sup>64</sup>, C. Isik, O. Kara, A. Kayis Topaksu, U. Kiminsu, G. Onengut, K. Ozdemir<sup>65</sup>, A. Polatoz, A.E. Simsek, B. Tali<sup>66</sup>, U.G. Tok, S. Turkcapar, I.S. Zorbakir

Çukurova University, Physics Department, Science and Art Faculty, Adana, Turkey

G. Karapinar, K. Ocalan<sup>67</sup>, M. Yalvac<sup>68</sup>

Middle East Technical University, Physics Department, Ankara, Turkey

B. Akgun, I.O. Atakisi, E. Gülmez, M. Kaya<sup>69</sup>, O. Kaya<sup>70</sup>, Ö. Özçelik, S. Tekten<sup>71</sup>, E.A. Yetkin<sup>72</sup>

Bogazici University, Istanbul, Turkey

A. Cakir, K. Cankocak<sup>63</sup>, Y. Komurcu, S. Sen<sup>73</sup>

Istanbul Technical University, Istanbul, Turkey

O. Aydilek, S. Cerci<sup>66</sup>, B. Haciosahinoglu, I. Hos<sup>74</sup>, B. Isildak<sup>75</sup>, B. Kaynak, S. Ozkorucuklu, H. Sert, C. Simsek, D. Sunar Cerci<sup>66</sup>, C. Zorbilmez

Istanbul University, Istanbul, Turkey

B. Grynyov

Institute for Scintillation Materials of National Academy of Science of Ukraine, Kharkiv, Ukraine

L. Levchuk

National Science Centre, Kharkiv Institute of Physics and Technology, Kharkiv, Ukraine

D. Anthony, E. Bhal, S. Bologna, J.J. Brooke, A. Bundock, E. Clement, D. Cussans, H. Flacher, M. Glowacki, J. Goldstein, G.P. Heath, H.F. Heath, L. Kreczko, B. Krikler, S. Paramesvaran, S. Seif El Nasr-Storey, V.J. Smith, N. Stylianou<sup>76</sup>, K. Walkingshaw Pass, R. White

*University of Bristol, Bristol, United Kingdom*

K.W. Bell, A. Belyaev<sup>77</sup>, C. Brew, R.M. Brown, D.J.A. Cockerill, C. Cooke, K.V. Ellis, K. Harder, S. Harper, M.-L. Holmberg<sup>78</sup>, J. Linacre, K. Manolopoulos, D.M. Newbold, E. Olaiya, D. Petyt, T. Reis, T. Schuh, C.H. Shepherd-Themistocleous, I.R. Tomalin, T. Williams

*Rutherford Appleton Laboratory, Didcot, United Kingdom*

R. Bainbridge, P. Bloch, S. Bonomally, J. Borg, S. Breeze, O. Buchmuller, V. Cepaitis, G.S. Chahal<sup>79</sup>, D. Colling, P. Dauncey, G. Davies, M. Della Negra, S. Fayer, G. Fedi, G. Hall, M.H. Hassanshahi, G. Iles, J. Langford, L. Lyons, A.-M. Magnan, S. Malik, A. Martelli, D.G. Monk, J. Nash<sup>80</sup>, M. Pesaresi, B.C. Radburn-Smith, D.M. Raymond, A. Richards, A. Rose, E. Scott, C. Seez, A. Shtipliyski, A. Tapper, K. Uchida, T. Virdee<sup>21</sup>, M. Vojinovic, N. Wardle, S.N. Webb, D. Winterbottom

*Imperial College, London, United Kingdom*

K. Coldham, J.E. Cole, A. Khan, P. Kyberd, I.D. Reid, L. Teodorescu, S. Zahid

*Brunel University, Uxbridge, United Kingdom*

S. Abdullin, A. Brinkerhoff, B. Caraway, J. Dittmann, K. Hatakeyama, A.R. Kanuganti, B. McMaster, M. Saunders, S. Sawant, C. Sutantawibul, J. Wilson

*Baylor University, Waco, TX, USA*

R. Bartek, A. Dominguez, R. Uniyal, A.M. Vargas Hernandez

*Catholic University of America, Washington, DC, USA*

A. Buccilli, S.I. Cooper, D. Di Croce, S.V. Gleyzer, C. Henderson, C.U. Perez, P. Rumerio<sup>81</sup>, C. West

*The University of Alabama, Tuscaloosa, AL, USA*

A. Akpinar, A. Albert, D. Arcaro, C. Cosby, Z. Demiragli, C. Erice, E. Fontanesi, D. Gastler, S. May, J. Rohlf, K. Salyer, D. Sperka, D. Spitzbart, I. Suarez, A. Tsatsos, S. Yuan

*Boston University, Boston, MA, USA*

G. Benelli, B. Burkle, X. Coubez<sup>23</sup>, D. Cutts, M. Hadley, U. Heintz, J.M. Hogan<sup>82</sup>, T. Kwon, G. Landsberg, K.T. Lau, D. Li, M. Lukasik, J. Luo, M. Narain, N. Pervan, S. Sagir<sup>83</sup>, F. Simpson, E. Usai, W.Y. Wong, X. Yan, D. Yu, W. Zhang

*Brown University, Providence, RI, USA*

J. Bonilla, C. Brainerd, R. Breedon, M. Calderon De La Barca Sanchez, M. Chertok, J. Conway, P.T. Cox, R. Erbacher, G. Haza, F. Jensen, O. Kukral, G. Mocellin, M. Mulhearn, D. Pellett, B. Regnery, D. Taylor, Y. Yao, F. Zhang

*University of California, Davis, Davis, CA, USA*

M. Bachtis, R. Cousins, A. Datta, D. Hamilton, J. Hauser, M. Ignatenko, M.A. Iqbal, T. Lam, W.A. Nash, S. Regnard, D. Saltzberg, B. Stone, V. Valuev

*University of California, Los Angeles, CA, USA*

Y. Chen, R. Clare, J.W. Gary, M. Gordon, G. Hanson, G. Karapostoli, O.R. Long, N. Manganeli, W. Si, S. Wimpenny, Y. Zhang

*University of California, Riverside, Riverside, CA, USA*

J.G. Branson, P. Chang, S. Cittolin, S. Cooperstein, D. Diaz, J. Duarte, R. Gerosa, L. Giannini, J. Guiang, R. Kansal, V. Krutelyov, R. Lee, J. Letts, M. Masciovecchio, F. Mokhtar, M. Pieri, B.V. Sathia Narayanan, V. Sharma, M. Tadel, F. Würthwein, Y. Xiang, A. Yagil

*University of California, San Diego, La Jolla, CA, USA*

N. Amin, C. Campagnari, M. Citron, G. Collura, A. Dorsett, V. Dutta, J. Incandela, M. Kilpatrick, J. Kim, B. Marsh, H. Mei, M. Oshiro, M. Quinnan, J. Richman, U. Sarica, F. Setti, J. Sheplock, P. Siddireddy, D. Stuart, S. Wang

*University of California, Santa Barbara - Department of Physics, Santa Barbara, CA, USA*

A. Bornheim, O. Cerri, I. Dutta, J.M. Lawhorn, N. Lu, J. Mao, H.B. Newman, T.Q. Nguyen, M. Spiropulu, J.R. Vlimant, C. Wang, S. Xie, Z. Zhang, R.Y. Zhu

*California Institute of Technology, Pasadena, CA, USA*

J. Alison, S. An, M.B. Andrews, P. Bryant, T. Ferguson, A. Harilal, C. Liu, T. Mudholkar, S. Murthy, M. Paulini, A. Roberts, A. Sanchez, W. Terrill

*Carnegie Mellon University, Pittsburgh, PA, USA*

J.P. Cumalat, W.T. Ford, A. Hassani, G. Karathanasis, E. MacDonald, R. Patel, A. Perloff, C. Savard, N. Schonbeck, K. Stenson, K.A. Ulmer, S.R. Wagner, N. Zipper

*University of Colorado Boulder, Boulder, CO, USA*

J. Alexander, S. Bright-Thonney, X. Chen, Y. Cheng, D.J. Cranshaw, J. Fan, X. Fan, D. Gadkari, S. Hogan, J. Monroy, J.R. Patterson, D. Quach, J. Reichert, M. Reid, A. Ryd, J. Thom, P. Wittich, R. Zou

*Cornell University, Ithaca, NY, USA*

M. Albrow, M. Alyari, G. Apollinari, A. Apresyan, A. Apyan, L.A.T. Bauerdick, D. Berry, J. Berryhill, P.C. Bhat, K. Burkett, J.N. Butler, A. Canepa, G.B. Cerati, H.W.K. Cheung, F. Chlebana, K.F. Di Petrillo, J. Dickinson, V.D. Elvira, Y. Feng, J. Freeman, A. Gandrakota, Z. Gecse, L. Gray, D. Green, S. Grünendahl, O. Gutsche, R.M. Harris, R. Heller, T.C. Herwig, J. Hirschauer, B. Jayatilaka, S. Jindariani, M. Johnson, U. Joshi, T. Klijnsma, B. Klima, K.H.M. Kwok, S. Lammel, D. Lincoln, R. Lipton, T. Liu, C. Madrid, K. Maeshima, C. Mantilla, D. Mason, P. McBride, P. Merkel, S. Mrenna, S. Nahn, J. Ngadiuba, V. Papadimitriou, N. Pastika, K. Pedro, C. Pena<sup>84</sup>, F. Ravera, A. Reinsvold Hall<sup>85</sup>, L. Ristori, E. Sexton-Kennedy, N. Smith, A. Soha, L. Spiegel, J. Strait, L. Taylor, S. Tkaczyk, N.V. Tran, L. Uplegger, E.W. Vaandering, H.A. Weber, I. Zoi

*Fermi National Accelerator Laboratory, Batavia, IL, USA*

P. Avery, D. Bourilkov, L. Cadamuro, V. Cherepanov, R.D. Field, D. Guerrero, M. Kim, E. Koenig, J. Konigsberg, A. Korytov, K.H. Lo, K. Matchev, N. Menendez, G. Mitselmakher, A. Muthirakalayi Madhu, N. Rawal, D. Rosenzweig, S. Rosenzweig, K. Shi, J. Wang, Z. Wu

*University of Florida, Gainesville, FL, USA*

T. Adams, A. Askew, R. Habibullah, V. Hagopian, K.F. Johnson, R. Khurana, T. Kolberg, G. Martinez, H. Prosper, C. Schiber, O. Viazlo, R. Yohay, J. Zhang

*Florida State University, Tallahassee, FL, USA*

M.M. Baarmand, S. Butalla, T. Elkafrawy<sup>86</sup>, M. Hohlmann, R. Kumar Verma, D. Noonan, M. Rahmani, F. Yumiceva

*Florida Institute of Technology, Melbourne, FL, USA*

M.R. Adams, H. Becerril Gonzalez, R. Cavanaugh, S. Dittmer, O. Evdokimov, C.E. Gerber, D.J. Hofman, D.S. Lemos, A.H. Merrit, C. Mills, G. Oh, T. Roy, S. Rudrabhatla, M.B. Tonjes, N. Varelas, J. Viinikainen, X. Wang, Z. Ye

*University of Illinois at Chicago (UIC), Chicago, IL, USA*

M. Alhousseini, K. Dilsiz<sup>87</sup>, L. Emediato, R.P. Gandrajula, O.K. Köseyan, J.-P. Merlo, A. Mestvirishvili<sup>88</sup>, J. Nachtman, H. Ogul<sup>89</sup>, Y. Onel, A. Penzo, C. Snyder, E. Tiras<sup>90</sup>

*The University of Iowa, Iowa City, IA, USA*

O. Amram, B. Blumenfeld, L. Corcodilos, J. Davis, A.V. Gritsan, S. Kyriacou, P. Maksimovic, J. Roskes, M. Swartz, T.Á. Vámi

*Johns Hopkins University, Baltimore, MD, USA*

A. Abreu, J. Anguiano, P. Baringer, A. Bean, Z. Flowers, T. Isidori, S. Khalil, J. King, G. Krintiras, A. Kropivnitskaya, M. Lazarovits, C. Le Mahieu, C. Lindsey, J. Marquez, N. Minafra, M. Murray, M. Nickel, C. Rogan, C. Royon, R. Salvatico, S. Sanders, E. Schmitz, C. Smith, Q. Wang, Z. Warner, J. Williams, G. Wilson

*The University of Kansas, Lawrence, KS, USA*

S. Duric, A. Ivanov, K. Kaadze, D. Kim, Y. Maravin, T. Mitchell, A. Modak, K. Nam

*Kansas State University, Manhattan, KS, USA*

F. Rebassoo, D. Wright

*Lawrence Livermore National Laboratory, Livermore, CA, USA*

E. Adams, A. Baden, O. Baron, A. Belloni, S.C. Eno, N.J. Hadley, S. Jabeen, R.G. Kellogg, T. Koeth, Y. Lai, S. Lascio, A.C. Mignerey, S. Nabili, C. Palmer, M. Seidel, A. Skuja, L. Wang, K. Wong

*University of Maryland, College Park, MD, USA*

D. Abercrombie, G. Andreassi, R. Bi, W. Busza, I.A. Cali, Y. Chen, M. D'Alfonso, J. Eysermans, C. Freer, G. Gomez-Ceballos, M. Goncharov, P. Harris, M. Hu, M. Klute, D. Kovalskyi, J. Krupa, Y.-J. Lee, K. Long, C. Mironov, C. Paus, D. Rankin, C. Roland, G. Roland, Z. Shi, G.S.F. Stephans, J. Wang, Z. Wang, B. Wyslouch

*Massachusetts Institute of Technology, Cambridge, MA, USA*

R.M. Chatterjee, A. Evans, J. Hiltbrand, Sh. Jain, B.M. Joshi, M. Krohn, Y. Kubota, J. Mans, M. Revering, R. Rusack, R. Saradhy, N. Schroeder, N. Strobbe, M.A. Wadud

*University of Minnesota, Minneapolis, MN, USA*

K. Bloom, M. Bryson, S. Chauhan, D.R. Claes, C. Fangmeier, L. Finco, F. Golf, C. Joo, I. Kravchenko, I. Reed, J.E. Siado, G.R. Snow<sup>†</sup>, W. Tabb, A. Wightman, F. Yan, A.G. Zecchinelli

*University of Nebraska-Lincoln, Lincoln, NE, USA*

G. Agarwal, H. Bandyopadhyay, L. Hay, I. Iashvili, A. Kharchilava, C. McLean, D. Nguyen, J. Pekkanen, S. Rappoccio, A. Williams

*State University of New York at Buffalo, Buffalo, NY, USA*

G. Alverson, E. Barberis, Y. Haddad, Y. Han, A. Hortiangtham, A. Krishna, J. Li, J. Lidrych, G. Madigan, B. Marzocchi, D.M. Morse, V. Nguyen, T. Orimoto, A. Parker, L. Skinnari, A. Tishelman-Charny, T. Wamorkar, B. Wang, A. Wisecarver, D. Wood

*Northeastern University, Boston, MA, USA*

S. Bhattacharya, J. Bueghly, Z. Chen, A. Gilbert, T. Gunter, K.A. Hahn, Y. Liu, N. Odell, M.H. Schmitt, M. Velasco

*Northwestern University, Evanston, IL, USA*

R. Band, R. Bucci, M. Cremonesi, A. Das, R. Goldouzian, M. Hildreth, K. Hurtado Anampa, C. Jessop, K. Lannon, J. Lawrence, N. Loukas, L. Lutton, J. Mariano, N. Marinelli, I. Mcalister, T. McCauley, C. Mcgrady, K. Mohrman, C. Moore, Y. Musienko<sup>13</sup>, R. Ruchti, A. Townsend, M. Wayne, M. Zarucki, L. Zygala

*University of Notre Dame, Notre Dame, IN, USA*

B. Bylsma, L.S. Durkin, B. Francis, C. Hill, A. Lesauvage, M. Nunez Ornelas, K. Wei, B.L. Winer, B.R. Yates

*The Ohio State University, Columbus, OH, USA*

F.M. Addesa, B. Bonham, P. Das, G. Dezoort, P. Elmer, A. Frankenthal, B. Greenberg, N. Haubrich, S. Higginbotham, A. Kalogeropoulos, G. Kopp, S. Kwan, D. Lange, D. Marlow, K. Mei, I. Ojalvo, J. Olsen, D. Stickland, C. Tully

*Princeton University, Princeton, NJ, USA*

S. Malik, S. Norberg

*University of Puerto Rico, Mayaguez, PR, USA*

A.S. Bakshi, V.E. Barnes, R. Chawla, S. Das, L. Gutay, M. Jones, A.W. Jung, D. Kondratyev, A.M. Koshy, M. Liu, G. Negro, N. Neumeister, G. Paspalaki, S. Piperov, A. Purohit, J.F. Schulte, M. Stojanovic, J. Thieman, F. Wang, R. Xiao, W. Xie

*Purdue University, West Lafayette, IN, USA*

J. Dolen, N. Parashar

*Purdue University Northwest, Hammond, IN, USA*

D. Acosta, A. Baty, T. Carnahan, M. Decaro, S. Dildick, K.M. Ecklund, S. Freed, P. Gardner, F.J.M. Geurts, A. Kumar, W. Li, B.P. Padley, R. Redjimi, J. Rotter, W. Shi, S. Yang, E. Yigitbasi, L. Zhang<sup>91</sup>, Y. Zhang, X. Zuo

*Rice University, Houston, TX, USA*

A. Bodek, P. de Barbaro, R. Demina, J.L. Dulemba, C. Fallon, T. Ferbel, M. Galanti, A. Garcia-Bellido, O. Hindrichs, A. Khukhunaishvili, E. Ranken, R. Taus, G.P. Van Onsem

*University of Rochester, Rochester, NY, USA*

K. Goulios

*The Rockefeller University, New York, NY, USA*

B. Chiarito, J.P. Chou, Y. Gershtein, E. Halkiadakis, A. Hart, M. Heindl, O. Karacheban<sup>25</sup>, I. Laflotte, A. Lath, R. Montalvo, K. Nash, M. Osherson, S. Salur, S. Schnetzer, S. Somalwar, R. Stone, S.A. Thayil, S. Thomas, H. Wang

*Rutgers, The State University of New Jersey, Piscataway, NJ, USA*

H. Acharya, A.G. Delannoy, S. Fiorendi, T. Holmes, S. Spanier

*University of Tennessee, Knoxville, TN, USA*

O. Bouhali<sup>92</sup>, M. Dalchenko, A. Delgado, R. Eusebi, J. Gilmore, T. Huang, T. Kamon<sup>93</sup>, H. Kim, S. Luo, S. Malhotra, R. Mueller, D. Overton, D. Rathjens, A. Safonov

*Texas A&M University, College Station, TX, USA*

N. Akchurin, J. Damgov, V. Hegde, K. Lamichhane, S.W. Lee, T. Mengke, S. Muthumuni, T. Peltola, I. Volobouev, Z. Wang, A. Whitbeck

Texas Tech University, Lubbock, TX, USA

E. Appelt, S. Greene, A. Gurrola, W. Johns, A. Melo, F. Romeo, P. Sheldon, S. Tuo, J. Velkovska

Vanderbilt University, Nashville, TN, USA

M.W. Arenton, B. Cardwell, B. Cox, G. Cummings, J. Hakala, R. Hirosky, M. Joyce, A. Ledovskoy, A. Li, C. Neu, C.E. Perez Lara, B. Tannenwald, S. White

University of Virginia, Charlottesville, VA, USA

N. Poudyal

Wayne State University, Detroit, MI, USA

S. Banerjee, K. Black, T. Bose, S. Dasu, I. De Bruyn, P. Everaerts, C. Galloni, H. He, M. Herndon, A. Herve, U. Hussain, A. Lanaro, A. Loeliger, R. Loveless, J. Madhusudanan Sreekala, A. Mallampalli, A. Mohammadi, D. Pinna, A. Savin, V. Shang, V. Sharma, W.H. Smith, D. Teague, S. Trembath-Reichert, W. Vetens

University of Wisconsin - Madison, Madison, WI, USA

S. Afanasiev, V. Andreev, Yu. Andreev, T. Aushev, M. Azarkin, A. Babaev, A. Belyaev, V. Blinov<sup>94</sup>, E. Boos, V. Borshch, D. Budkouski, V. Bunichev, O. Bychkova, V. Chekhovsky, R. Chistov<sup>94</sup>, M. Danilov<sup>94</sup>, A. Dermenev, T. Dimova<sup>94</sup>, I. Dremin, M. Dubinin<sup>84</sup>, L. Dudko, V. Epshteyn<sup>95</sup>, G. Gavrillov, V. Gavrillov, S. Gninenko, V. Golovtsov, N. Golubev, I. Golutvin, I. Gorbunov, V. Ivanchenko, Y. Ivanov, V. Kachanov, L. Kardapoltsev<sup>94</sup>, V. Karjavine, A. Karneyeu, V. Kim<sup>94</sup>, M. Kirakosyan, D. Kirpichnikov, M. Kirsanov, V. Klyukhin, O. Kodolova<sup>96</sup>, D. Konstantinov, V. Korenkov, A. Kozyrev<sup>94</sup>, N. Krasnikov, E. Kuznetsova<sup>97</sup>, A. Lanev, A. Litomin, N. Lychkovskaya, V. Makarenko, A. Malakhov, V. Matveev<sup>94</sup>, V. Murzin, A. Nikitenko<sup>98</sup>, S. Obraztsov, V. Okhotnikov, V. Oreshkin, A. Oskin, I. Ovtin<sup>94</sup>, V. Palichik, P. Parygin<sup>99</sup>, A. Pashenkov, V. Perelygin, M. Perfilov, S. Petrushanko, G. Pivovarov, S. Polikarpov<sup>94</sup>, V. Popov, O. Radchenko<sup>94</sup>, M. Savina, V. Savrin, V. Shalaev, S. Shmatov, S. Shulha, Y. Skovpen<sup>94</sup>, S. Slabospitskii, I. Smirnov, V. Smirnov, A. Snigirev, D. Sosnov, A. Stepenov, V. Sulimov, E. Tcherniaev, A. Terkulov, O. Teryaev, M. Toms<sup>100</sup>, A. Toropin, L. Uvarov, A. Uzunian, E. Vlasov<sup>101</sup>, S. Volkov, A. Vorobyev, N. Voytishin, B.S. Yuldashev<sup>102</sup>, A. Zarubin, I. Zhizhin, A. Zhokin

Authors affiliated with an institute or an international laboratory covered by a cooperation agreement with CERN

<sup>†</sup> Deceased.

<sup>1</sup> Also at Yerevan State University, Yerevan, Armenia.

<sup>2</sup> Also at TU Wien, Vienna, Austria.

<sup>3</sup> Also at Institute of Basic and Applied Sciences, Faculty of Engineering, Arab Academy for Science, Technology and Maritime Transport, Alexandria, Egypt.

<sup>4</sup> Also at Université Libre de Bruxelles, Bruxelles, Belgium.

<sup>5</sup> Also at Universidade Estadual de Campinas, Campinas, Brazil.

<sup>6</sup> Also at Federal University of Rio Grande do Sul, Porto Alegre, Brazil.

<sup>7</sup> Also at UFMS, Nova Andradina, Brazil.

<sup>8</sup> Also at The University of the State of Amazonas, Manaus, Brazil.

<sup>9</sup> Also at University of Chinese Academy of Sciences, Beijing, China.

<sup>10</sup> Also at Nanjing Normal University Department of Physics, Nanjing, China.

<sup>11</sup> Now at The University of Iowa, Iowa City, Iowa, USA.

<sup>12</sup> Also at University of Chinese Academy of Sciences, Beijing, China.

<sup>13</sup> Also at an institute or an international laboratory covered by a cooperation agreement with CERN.

<sup>14</sup> Also at Helwan University, Cairo, Egypt.

<sup>15</sup> Now at Zewail City of Science and Technology, Zewail, Egypt.

<sup>16</sup> Also at Suez University, Suez, Egypt.

<sup>17</sup> Now at British University in Egypt, Cairo, Egypt.

<sup>18</sup> Also at Purdue University, West Lafayette, Indiana, USA.

<sup>19</sup> Also at Université de Haute Alsace, Mulhouse, France.

- <sup>20</sup> Also at Erzincan Binali Yildirim University, Erzincan, Turkey.
- <sup>21</sup> Also at CERN, European Organization for Nuclear Research, Geneva, Switzerland.
- <sup>22</sup> Also at University of Hamburg, Hamburg, Germany.
- <sup>23</sup> Also at RWTH Aachen University, III. Physikalisches Institut A, Aachen, Germany.
- <sup>24</sup> Also at Isfahan University of Technology, Isfahan, Iran.
- <sup>25</sup> Also at Brandenburg University of Technology, Cottbus, Germany.
- <sup>26</sup> Also at Forschungszentrum Jülich, Juelich, Germany.
- <sup>27</sup> Also at Physics Department, Faculty of Science, Assiut University, Assiut, Egypt.
- <sup>28</sup> Also at Karoly Robert Campus, MATE Institute of Technology, Gyongyos, Hungary.
- <sup>29</sup> Also at Institute of Physics, University of Debrecen, Debrecen, Hungary.
- <sup>30</sup> Also at Institute of Nuclear Research ATOMKI, Debrecen, Hungary.
- <sup>31</sup> Now at Universitatea Babeş-Bolyai - Facultatea de Fizica, Cluj-Napoca, Romania.
- <sup>32</sup> Also at MTA-ELTE Lendület CMS Particle and Nuclear Physics Group, Eötvös Loránd University, Budapest, Hungary.
- <sup>33</sup> Also at Faculty of Informatics, University of Debrecen, Debrecen, Hungary.
- <sup>34</sup> Also at Wigner Research Centre for Physics, Budapest, Hungary.
- <sup>35</sup> Also at Punjab Agricultural University, Ludhiana, India.
- <sup>36</sup> Also at UPES - University of Petroleum and Energy Studies, Dehradun, India.
- <sup>37</sup> Also at University of Hyderabad, Hyderabad, India.
- <sup>38</sup> Also at University of Visva-Bharati, Santiniketan, India.
- <sup>39</sup> Also at Indian Institute of Science (IISc), Bangalore, India.
- <sup>40</sup> Also at Indian Institute of Technology (IIT), Mumbai, India.
- <sup>41</sup> Also at IIT Bhubaneswar, Bhubaneswar, India.
- <sup>42</sup> Also at Institute of Physics, Bhubaneswar, India.
- <sup>43</sup> Also at Deutsches Elektronen-Synchrotron, Hamburg, Germany.
- <sup>44</sup> Also at Sharif University of Technology, Tehran, Iran.
- <sup>45</sup> Also at Department of Physics, University of Science and Technology of Mazandaran, Behshahr, Iran.
- <sup>46</sup> Also at Italian National Agency for New Technologies, Energy and Sustainable Economic Development, Bologna, Italy.
- <sup>47</sup> Also at Centro Siciliano di Fisica Nucleare e di Struttura Della Materia, Catania, Italy.
- <sup>48</sup> Also at Scuola Superiore Meridionale, Università di Napoli 'Federico II', Napoli, Italy.
- <sup>49</sup> Also at Università di Napoli 'Federico II', Napoli, Italy.
- <sup>50</sup> Also at Consiglio Nazionale delle Ricerche - Istituto Officina dei Materiali, Perugia, Italy.
- <sup>51</sup> Also at Department of Applied Physics, Faculty of Science and Technology, Universiti Kebangsaan Malaysia, Bangi, Malaysia.
- <sup>52</sup> Also at Consejo Nacional de Ciencia y Tecnología, Mexico City, Mexico.
- <sup>53</sup> Also at IRFU, CEA, Université Paris-Saclay, Gif-sur-Yvette, France.
- <sup>54</sup> Also at Faculty of Physics, University of Belgrade, Belgrade, Serbia.
- <sup>55</sup> Also at Trincomalee Campus, Eastern University, Sri Lanka, Nilaveli, Sri Lanka.
- <sup>56</sup> Also at INFN Sezione di Pavia, Università di Pavia, Pavia, Italy.
- <sup>57</sup> Also at National and Kapodistrian University of Athens, Athens, Greece.
- <sup>58</sup> Also at Ecole Polytechnique Fédérale Lausanne, Lausanne, Switzerland.
- <sup>59</sup> Also at Universität Zürich, Zurich, Switzerland.
- <sup>60</sup> Also at Stefan Meyer Institute for Subatomic Physics, Vienna, Austria.
- <sup>61</sup> Also at Laboratoire d'Annecy-le-Vieux de Physique des Particules, IN2P3-CNRS, Annecy-le-Vieux, France.
- <sup>62</sup> Also at Şırnak University, Şırnak, Turkey.
- <sup>63</sup> Also at Near East University, Research Center of Experimental Health Science, Mersin, Turkey.
- <sup>64</sup> Also at Konya Technical University, Konya, Turkey.
- <sup>65</sup> Also at Izmir Bakircay University, Izmir, Turkey.
- <sup>66</sup> Also at Adiyaman University, Adiyaman, Turkey.
- <sup>67</sup> Also at Necmettin Erbakan University, Konya, Turkey.
- <sup>68</sup> Also at Bozok Universitetesi Rektörlüğü, Yozgat, Turkey.
- <sup>69</sup> Also at Marmara University, Istanbul, Turkey.
- <sup>70</sup> Also at Milli Savunma University, Istanbul, Turkey.
- <sup>71</sup> Also at Kafkas University, Kars, Turkey.
- <sup>72</sup> Also at Istanbul Bilgi University, Istanbul, Turkey.
- <sup>73</sup> Also at Hacettepe University, Ankara, Turkey.
- <sup>74</sup> Also at Istanbul University - Cerrahpasa, Faculty of Engineering, Istanbul, Turkey.
- <sup>75</sup> Also at Yildiz Technical University, Istanbul, Turkey.
- <sup>76</sup> Also at Vrije Universiteit Brussel, Brussel, Belgium.
- <sup>77</sup> Also at School of Physics and Astronomy, University of Southampton, Southampton, United Kingdom.
- <sup>78</sup> Also at University of Bristol, Bristol, United Kingdom.
- <sup>79</sup> Also at IPPP Durham University, Durham, United Kingdom.
- <sup>80</sup> Also at Monash University, Faculty of Science, Clayton, Australia.
- <sup>81</sup> Also at Università di Torino, Torino, Italy.
- <sup>82</sup> Also at Bethel University, St. Paul, Minnesota, USA.



- <sup>83</sup> Also at Karamanoğlu Mehmetbey University, Karaman, Turkey.
- <sup>84</sup> Also at California Institute of Technology, Pasadena, California, USA.
- <sup>85</sup> Also at United States Naval Academy, Annapolis, Maryland, USA.
- <sup>86</sup> Also at Ain Shams University, Cairo, Egypt.
- <sup>87</sup> Also at Bingol University, Bingol, Turkey.
- <sup>88</sup> Also at Georgian Technical University, Tbilisi, Georgia.
- <sup>89</sup> Also at Sinop University, Sinop, Turkey.
- <sup>90</sup> Also at Erciyes University, Kayseri, Turkey.
- <sup>91</sup> Also at Institute of Modern Physics and Key Laboratory of Nuclear Physics and Ion-beam Application (MOE) - Fudan University, Shanghai, China.
- <sup>92</sup> Also at Texas A&M University at Qatar, Doha, Qatar.
- <sup>93</sup> Also at Kyungpook National University, Daegu, Korea.
- <sup>94</sup> Also at another institute or international laboratory covered by a cooperation agreement with CERN.
- <sup>95</sup> Now at Istanbul University, Istanbul, Turkey.
- <sup>96</sup> Also at Yerevan Physics Institute, Yerevan, Armenia.
- <sup>97</sup> Now at University of Florida, Gainesville, Florida, USA.
- <sup>98</sup> Also at Imperial College, London, United Kingdom.
- <sup>99</sup> Now at University of Rochester, Rochester, New York, USA.
- <sup>100</sup> Now at Baylor University, Waco, Texas, USA.
- <sup>101</sup> Now at INFN Sezione di Torino, Università di Torino, Torino, Italy; Università del Piemonte Orientale, Novara, Italy.
- <sup>102</sup> Also at Institute of Nuclear Physics of the Uzbekistan Academy of Sciences, Tashkent, Uzbekistan.